

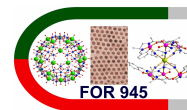
Quasi exact few-body quantum magnetism

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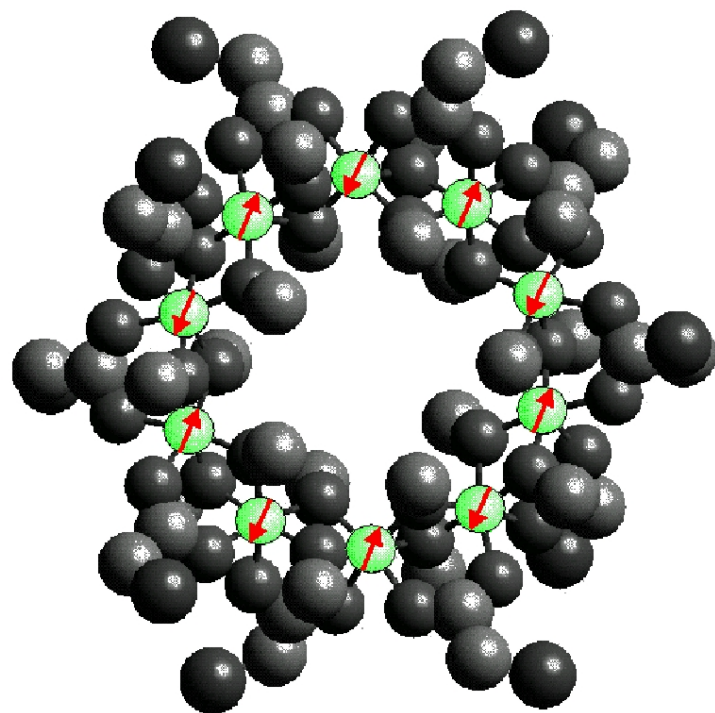
<http://obelix.physik.uni-bielefeld.de/~schnack/>

SuperMUC Status and Results Workshop
LRZ in Garching, April 26-27, 2016



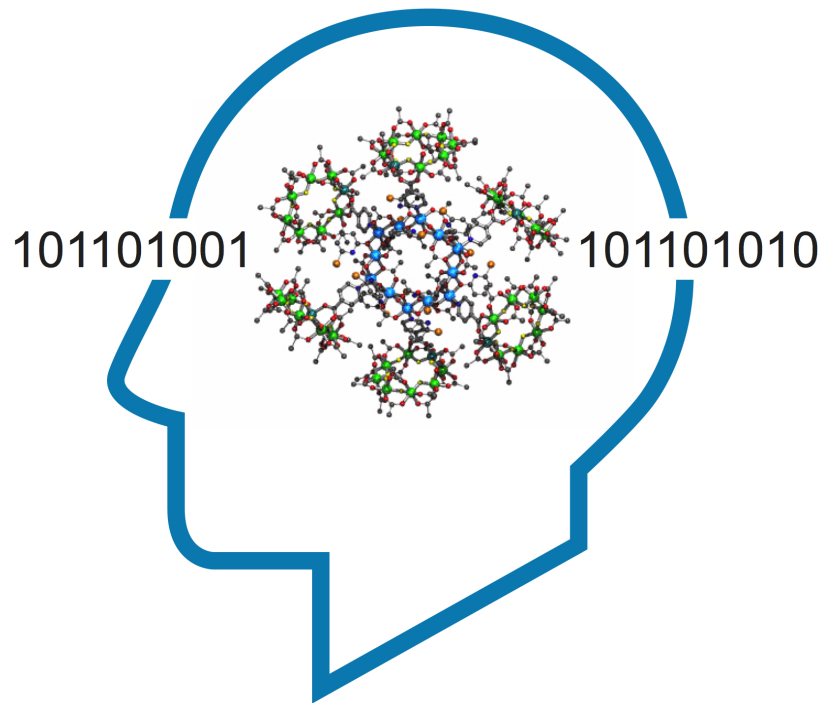
Typical problems in quantum magnetism

You have got a molecule!



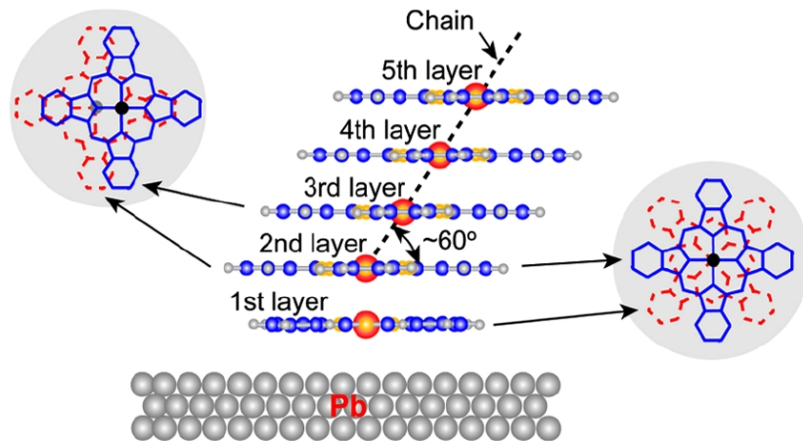
Congratulations!

You want to build a quantum computer!



Very smart!

You want to deposit your molecule!



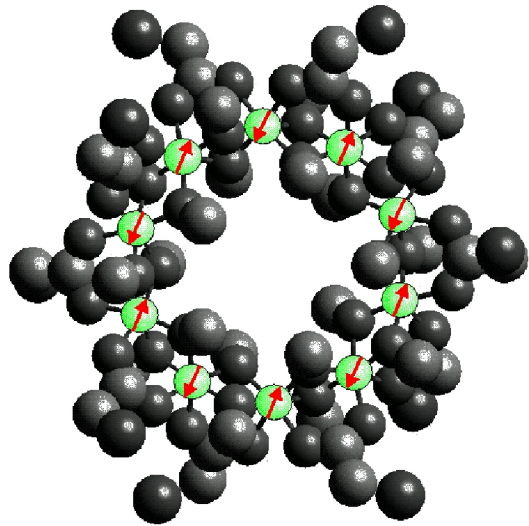
Next generation magnetic storage!

You have got an idea about the modeling!

$$\tilde{H} = -2 \sum_{i < j} J_{ij} \vec{\tilde{s}}(i) \cdot \vec{\tilde{s}}(j) + g \mu_B B \sum_i^N s_z(i)$$

Heisenberg

Zeeman



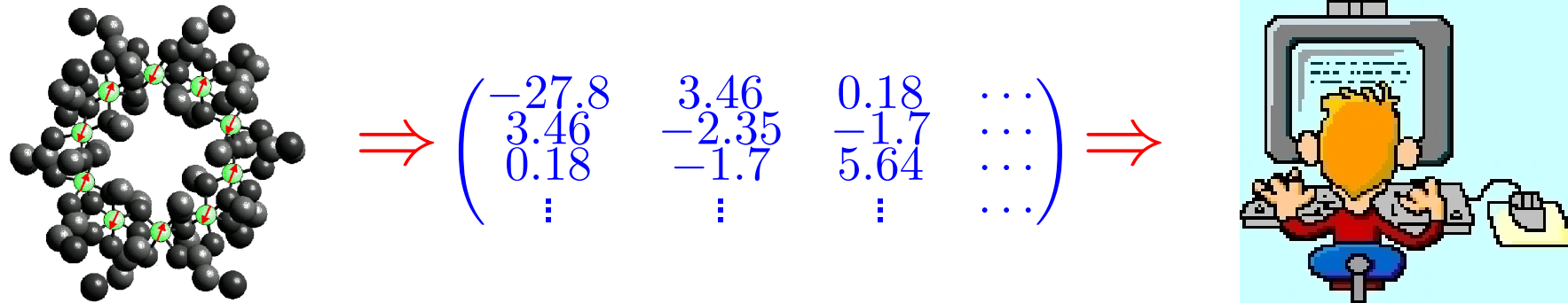
You have to solve the Schrödinger equation!

$$\underline{H} | \phi_n \rangle = E_n | \phi_n \rangle$$

Eigenvalues E_n and eigenvectors $| \phi_n \rangle$

- needed for spectroscopy (EPR, INS, NMR);
- needed for thermodynamic functions (magnetization, susceptibility, heat capacity);
- needed for time evolution (pulsed EPR, simulate quantum computing, thermalization).

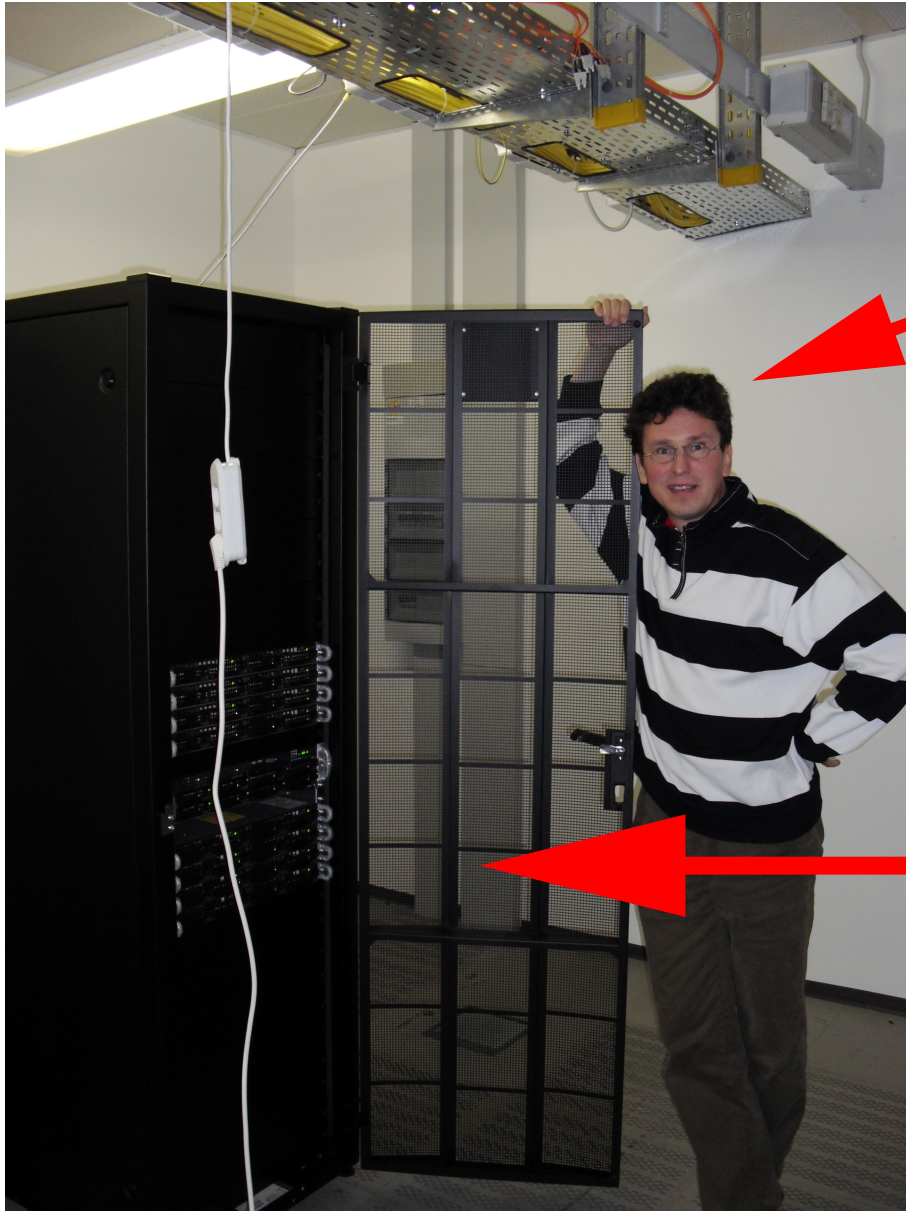
In the end it's always a big matrix!



$$\text{Fe}_{10}^{\text{III}}: N = 10, s = 5/2$$

Dimension=**60,466,176**. Maybe **too** big?

Thank God, we have computers



“Espresso-doped multi-core”

128 cores, 384 GB RAM

... but that's not enough!



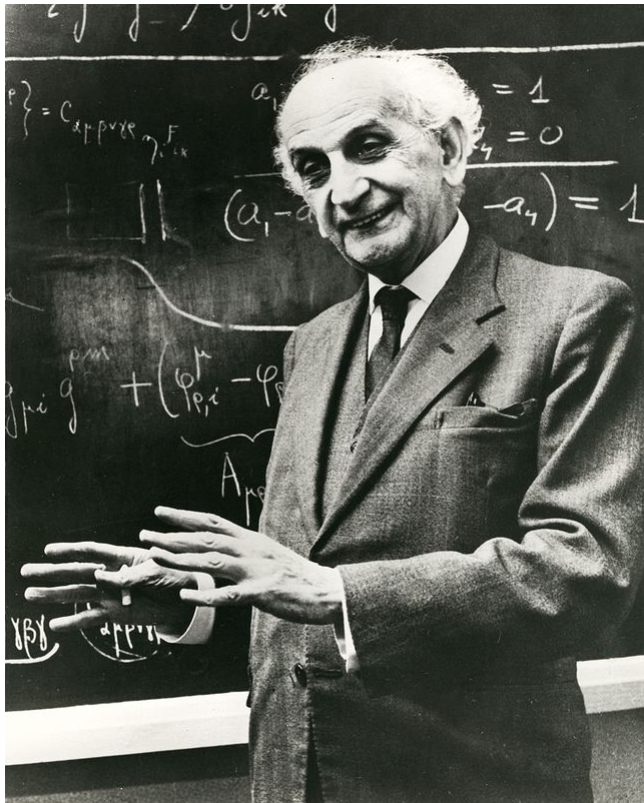
SuperMUC @ LRZ!

If matrix diagonalization does not work:

Finite-temperature Lanczos Method

(Good for dimensions up to 10^{10} .)

Lanczos – a Krylov space method



- Idea: exact diagonalization in reduced basis sets.
- But which set to choose???
- Idea: generate the basis set with the operator you want to diagonalize:
 $\{ |\phi\rangle, \tilde{H}|\phi\rangle, \tilde{H}^2|\phi\rangle, \tilde{H}^3|\phi\rangle, \dots \}$
- But which starting vector to choose???
- Idea: almost any will do!
- Cornelius Lanczos (Lánczos Kornél, 1893-1974)

(1) C. Lanczos, J. Res. Nat. Bur. Stand. **45**, 255 (1950).

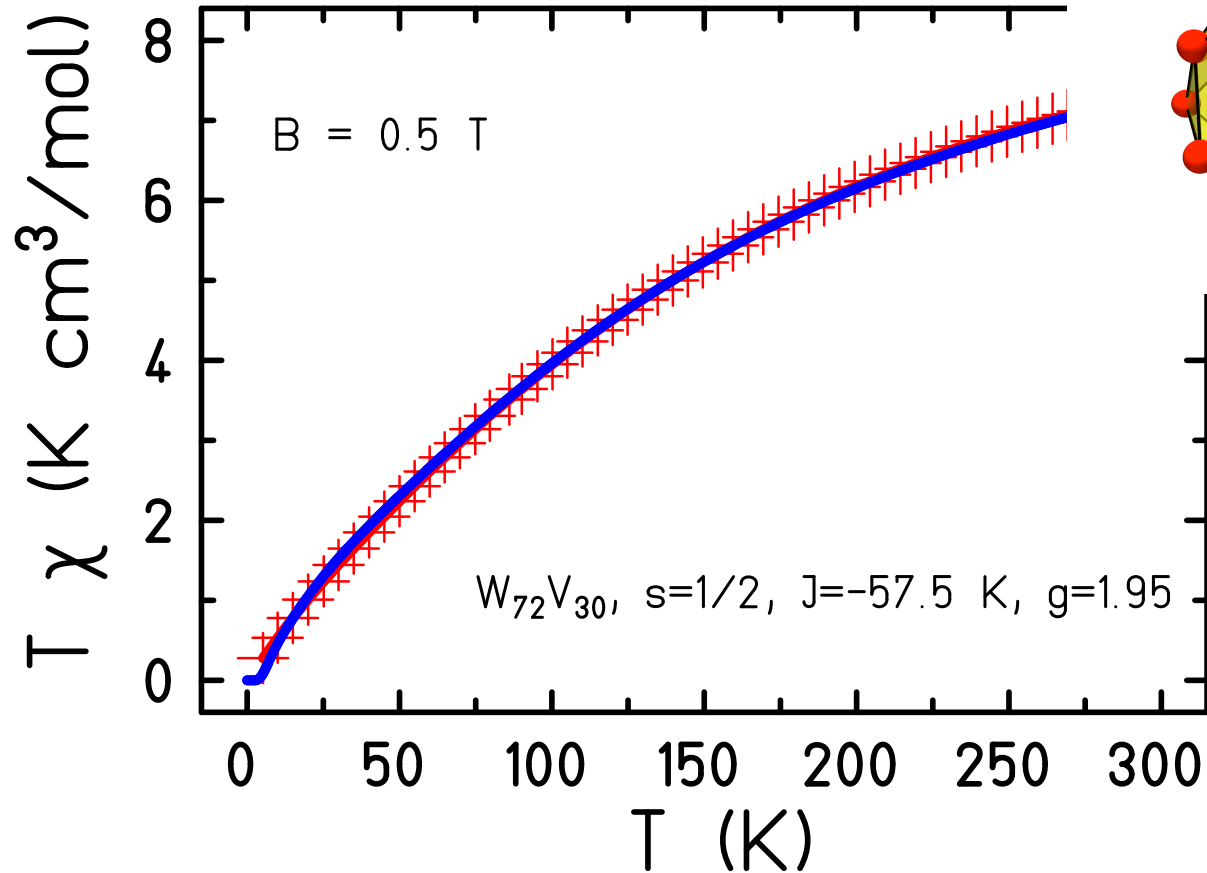
Finite-temperature Lanczos Method I

$$\begin{aligned}
 Z(T, B) &= \sum_{\nu} \langle \nu | \exp \left\{ -\beta \tilde{H} \right\} | \nu \rangle \\
 \langle \nu | \exp \left\{ -\beta \tilde{H} \right\} | \nu \rangle &\approx \sum_n \langle \nu | n(\nu) \rangle \exp \left\{ -\beta \epsilon_n \right\} \langle n(\nu) | \nu \rangle \\
 Z(T, B) &\approx \frac{\dim(\mathcal{H})}{R} \sum_{\nu=1}^R \sum_{n=1}^{N_L} \exp \left\{ -\beta \epsilon_n \right\} |\langle n(\nu) | \nu \rangle|^2
 \end{aligned}$$

- $|n(\nu)\rangle$ n-th Lanczos eigenvector starting from $|\nu\rangle$
- Partition function replaced by a small sum: $R = 1 \dots 10, N_L \approx 100$.

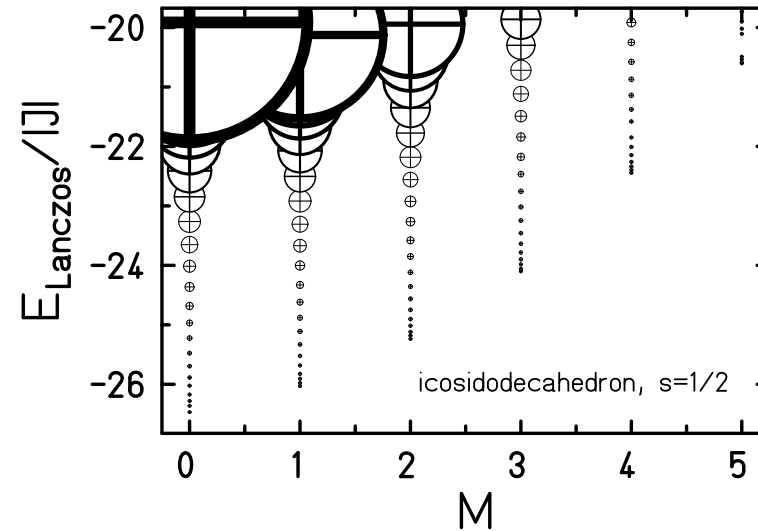
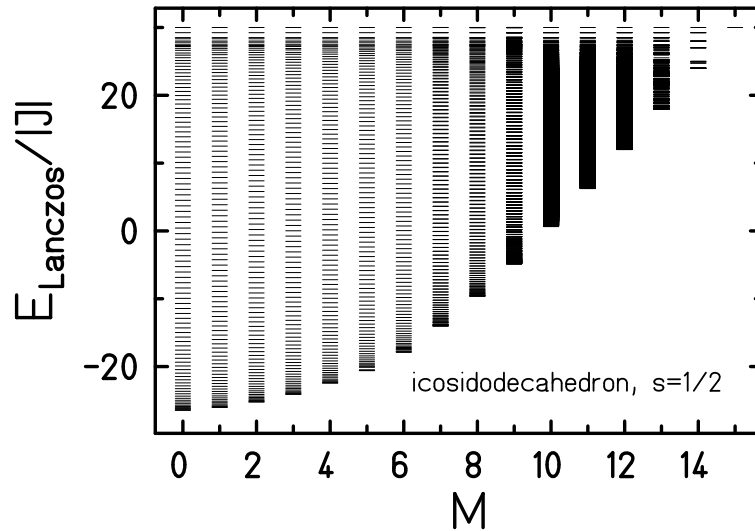
J. Jaklic and P. Prelovsek, Phys. Rev. B **49**, 5065 (1994).

Icosidodecahedron $s = 1/2$



Exp. data: A. M. Todea, A. Merca, H. Bögge, T. Glaser, L. Engelhardt, R. Prozorov, M. Luban, A. Müller, Chem. Commun., 3351 (2009).

Icosidodecahedron $s = 1/2$

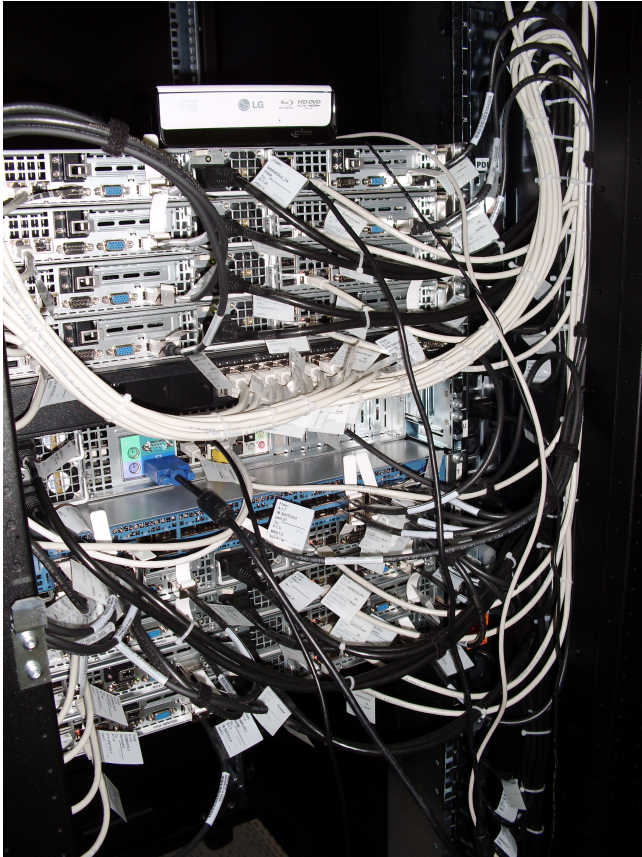


- The true spectrum will be much denser. This is miraculously compensated for by the weights.

$$Z(T, B) \approx \frac{\dim(\mathcal{H})}{R} \sum_{\nu=1}^R \sum_{n=1}^{N_L} \exp \{-\beta \epsilon_n\} |\langle n(\nu, \Gamma) | \nu, \Gamma \rangle|^2$$

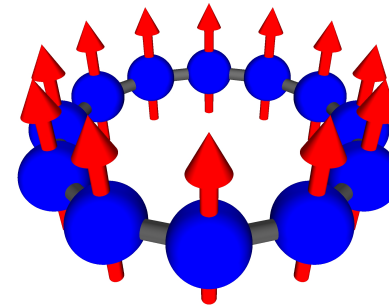
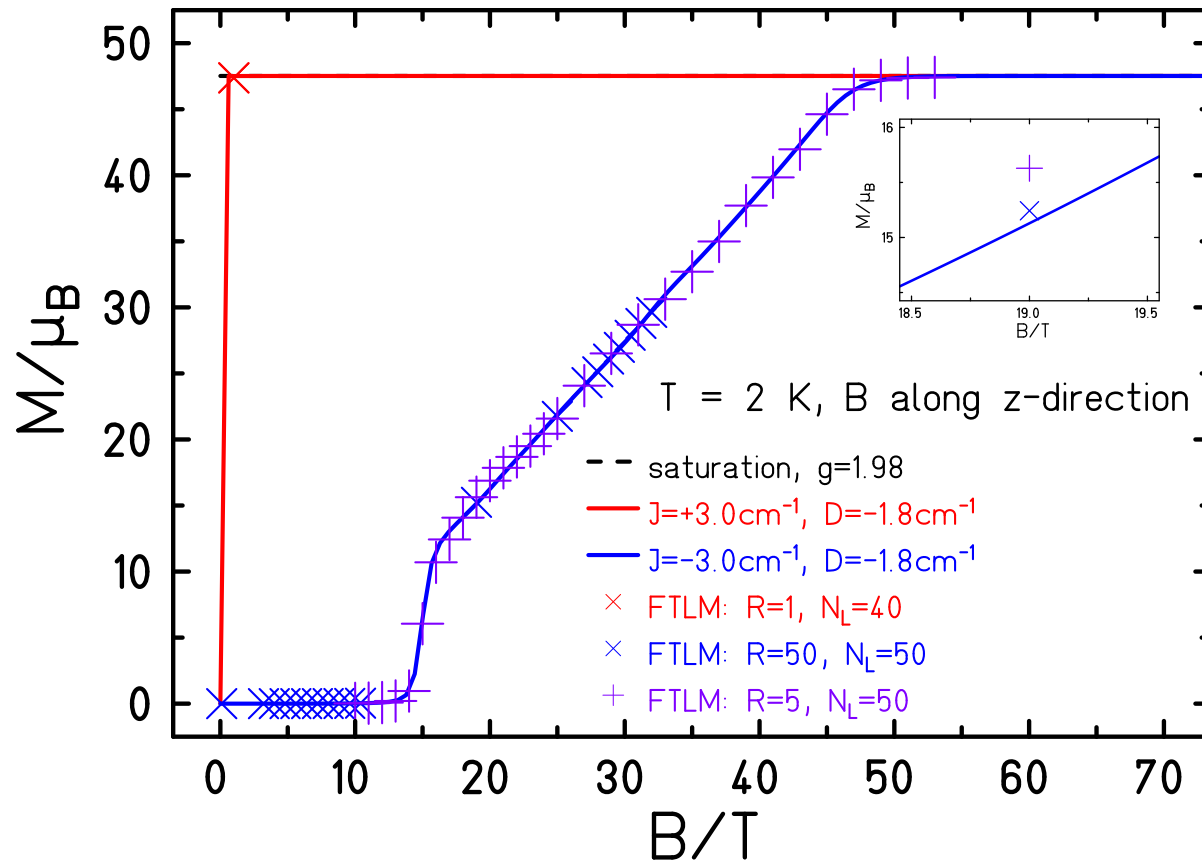
J. Schnack, O. Wendland, Eur. Phys. J. B **78** (2010) 535-541

Finite-temperature Lanczos Method II



- FTLM is a trace estimator.
- FTLM consists of matrix vector multiplications.
- Ideal for openMP and MPI!
- Vector spaces with dimensions of up to 10^{10} can be used.
- Use of symmetries improves the accuracy.

A fictitious $\text{Mn}_{12}^{\text{III}}$ – M_z vs B_z



$s = 2$

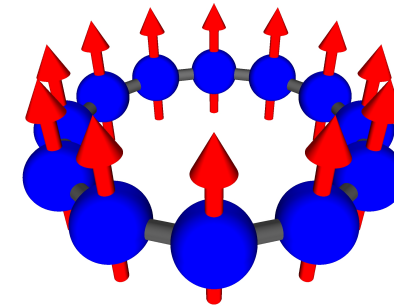
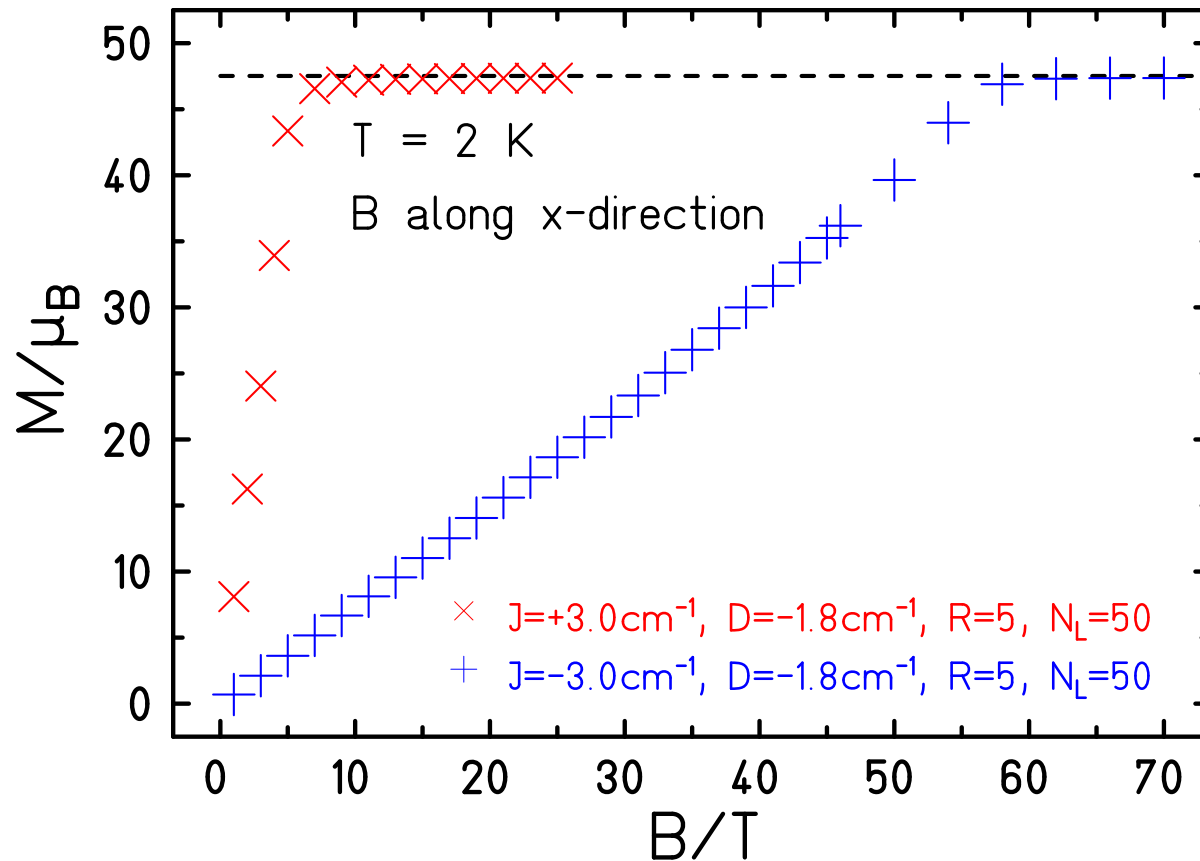
$\dim(\mathcal{H}) = 244, 140, 625$

collinear easy axes

A few days compared to *impossible!*

O. Hanebaum, J. Schnack, Eur. Phys. J. B **87**, 194 (2014)

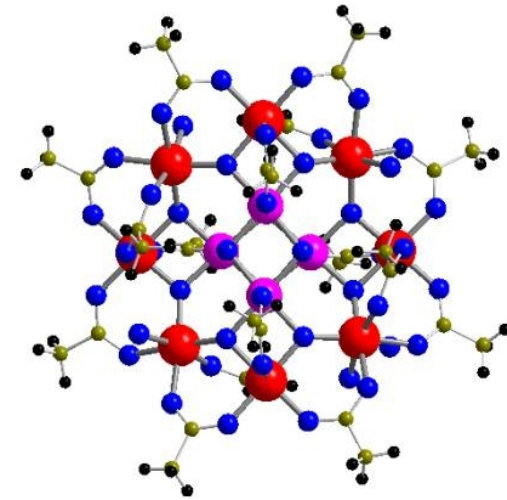
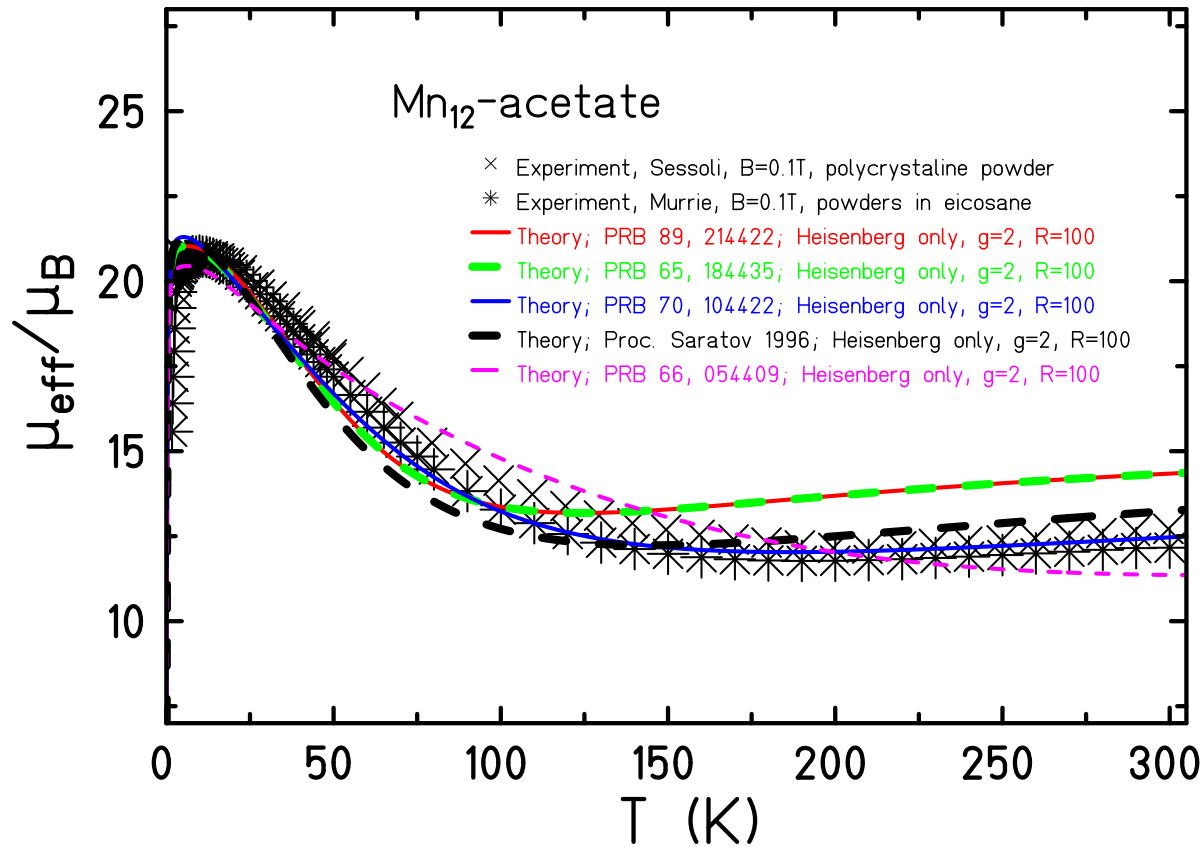
A fictitious $\text{Mn}_{12}^{\text{III}}$ – M_x vs B_x



No other method can deliver these curves!

O. Hanebaum, J. Schnack, Eur. Phys. J. B **87**, 194 (2014)

Effective magnetic moment of



No other method can deliver these curves!

O. Hanebaum, J. Schnack, Eur. Phys. J. B **87**, 194 (2014)

The magnetocaloric effect

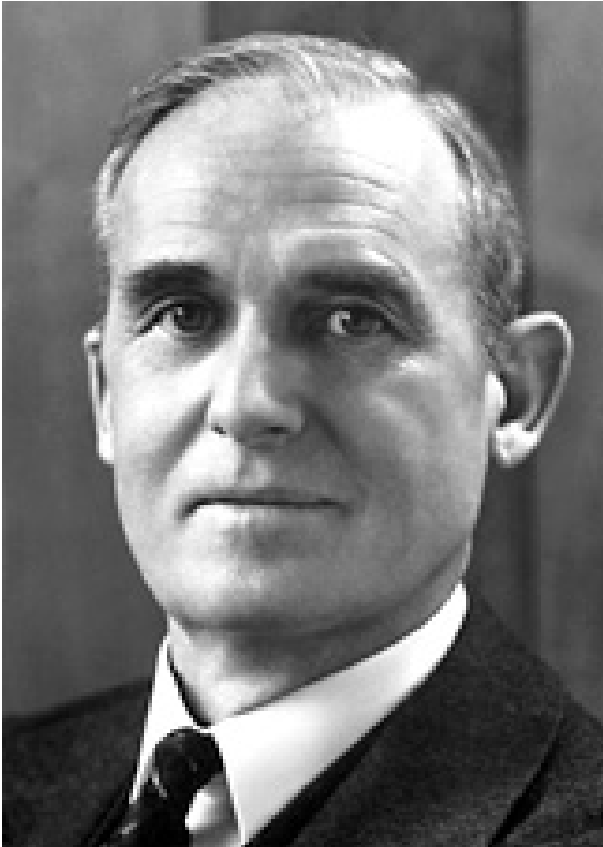
Magnetocaloric effect – Basics



- Heating or cooling in a varying magnetic field. Predicted, discussed, discovered by Thomson, Warburg, Weiss, and Piccard (1).
- Typical rates: 0.5 . . . 2 K/T.
- Giant magnetocaloric effect: 3 . . . 4 K/T e.g. in $\text{Gd}_5(\text{Si}_x\text{Ge}_{1-x})_4$ alloys ($x \leq 0.5$).
- **Scientific goal I: room temperature applications.**
- **Scientific goal II: sub-Kelvin cooling.**

(1) A. Smith, Eur. Phys. J. H **38**, 507 (2013).

Sub-Kelvin cooling: Nobel prize 1949



The Nobel Prize in Chemistry 1949 was awarded to William F. Giaouque *for his contributions in the field of chemical thermodynamics, particularly concerning the behaviour of substances at extremely low temperatures.*

Sub-Kelvin cooling: Nobel prize 1949

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LETTERS TO THE EDITOR

Attainment of Temperatures Below 1° Absolute by Demagnetization of $\text{Gd}_2(\text{SO}_4)_3 \cdot 8\text{H}_2\text{O}$

We have recently carried out some preliminary experiments on the adiabatic demagnetization of $\text{Gd}_2(\text{SO}_4)_3 \cdot 8\text{H}_2\text{O}$ at the temperatures of liquid helium. As previously predicted by one of us, a large fractional lowering of the absolute temperature was obtained.

An iron-free solenoid producing a field of about 8000 gauss was used for all the measurements. The amount of $\text{Gd}_2(\text{SO}_4)_3 \cdot 8\text{H}_2\text{O}$ was 61 g. The observations were checked by many repetitions of the cooling. The temperatures were measured by means of the inductance of a coil surrounding the gadolinium sulfate. The coil was immersed in liquid helium and isolated from the gadolinium by means of an evacuated space. The thermometer was in excellent agreement with the temperature of liquid helium as indicated by its vapor pressure down to 1.5°K.

On March 19, starting at a temperature of about 3.4°K, the material cooled to 0.53°K. On April 8, starting at about 2°, a temperature of 0.34°K was reached. On April 9, starting at about 1.5°, a temperature of 0.25°K was attained.

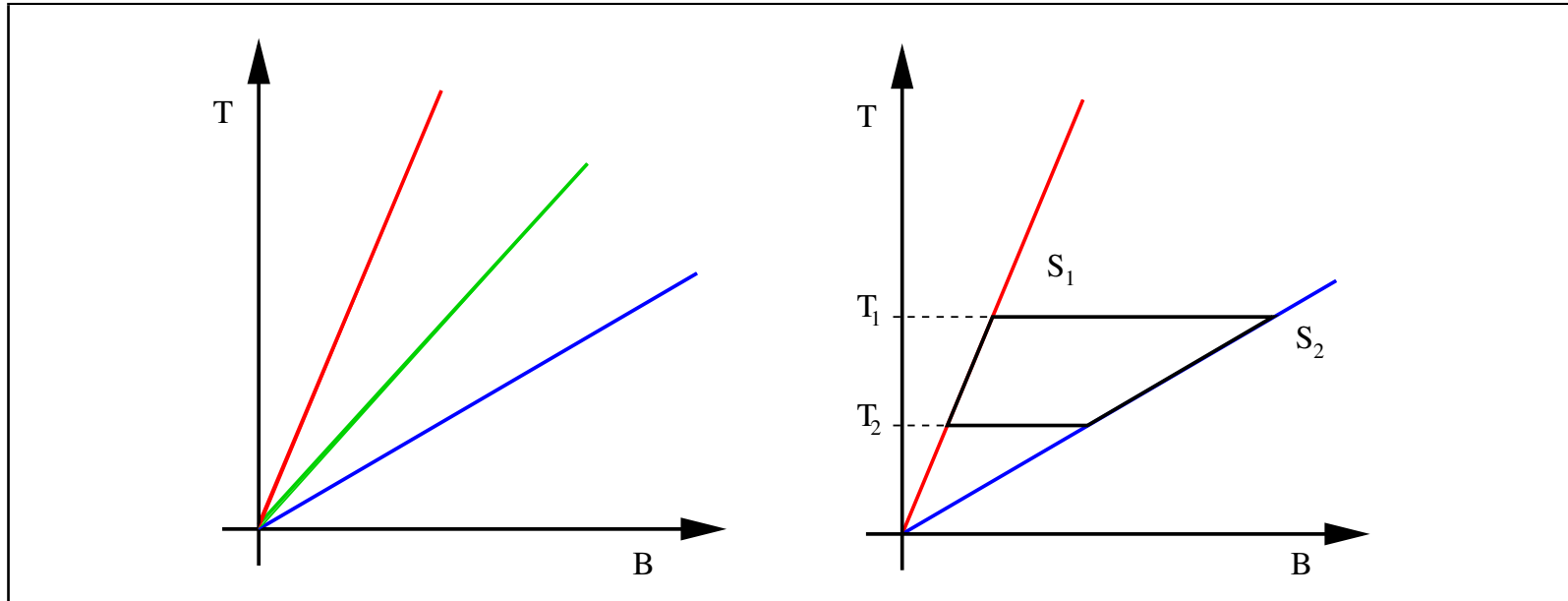
It is apparent that it will be possible to obtain much lower temperatures, especially when successive demagnetizations are utilized.

W. F. GIAUQUE
D. P. MACDOUGALL

Department of Chemistry,
University of California,
Berkeley, California,
April 12, 1933.

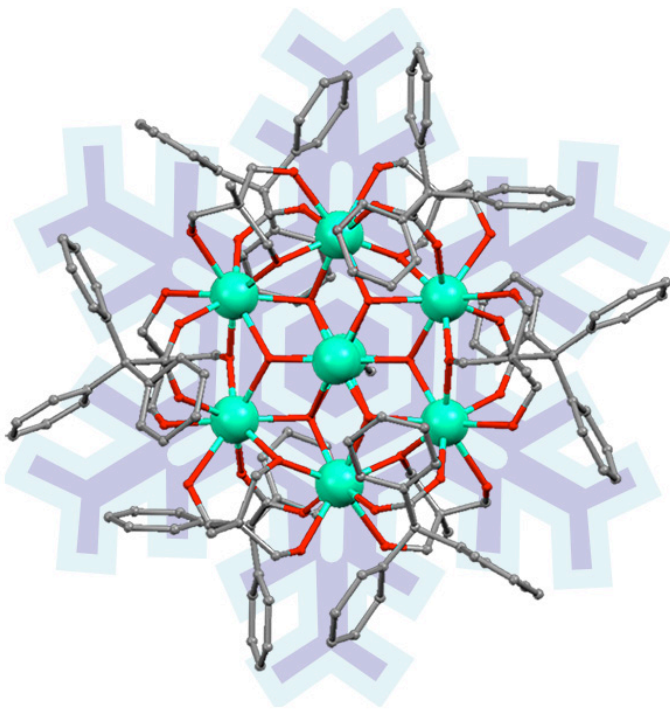
W. F. Giauque and D. MacDougall, *Phys. Rev.* **43**, 768 (1933).

Magnetocaloric effect – Paramagnets



- Ideal paramagnet: $S(T, B) = f(B/T)$, i.e. $S = const \Rightarrow T \propto B$.
- At low T pronounced effects of dipolar interaction prevent further effective cooling.

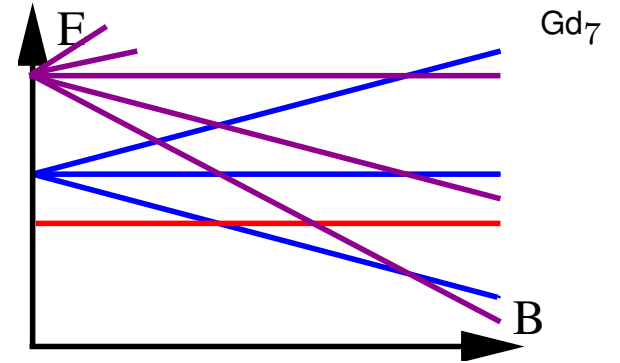
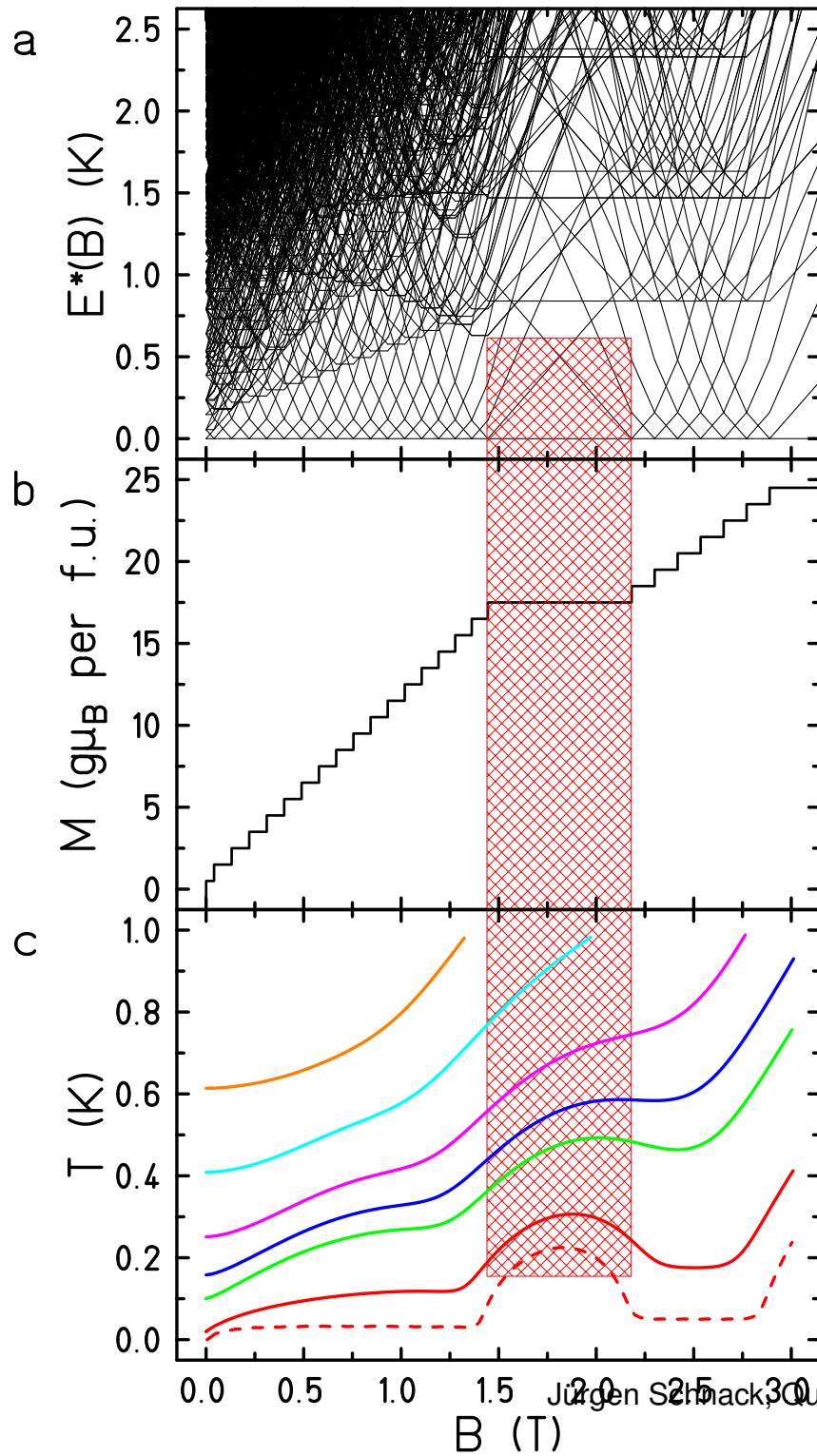
Gd₇ – Basics



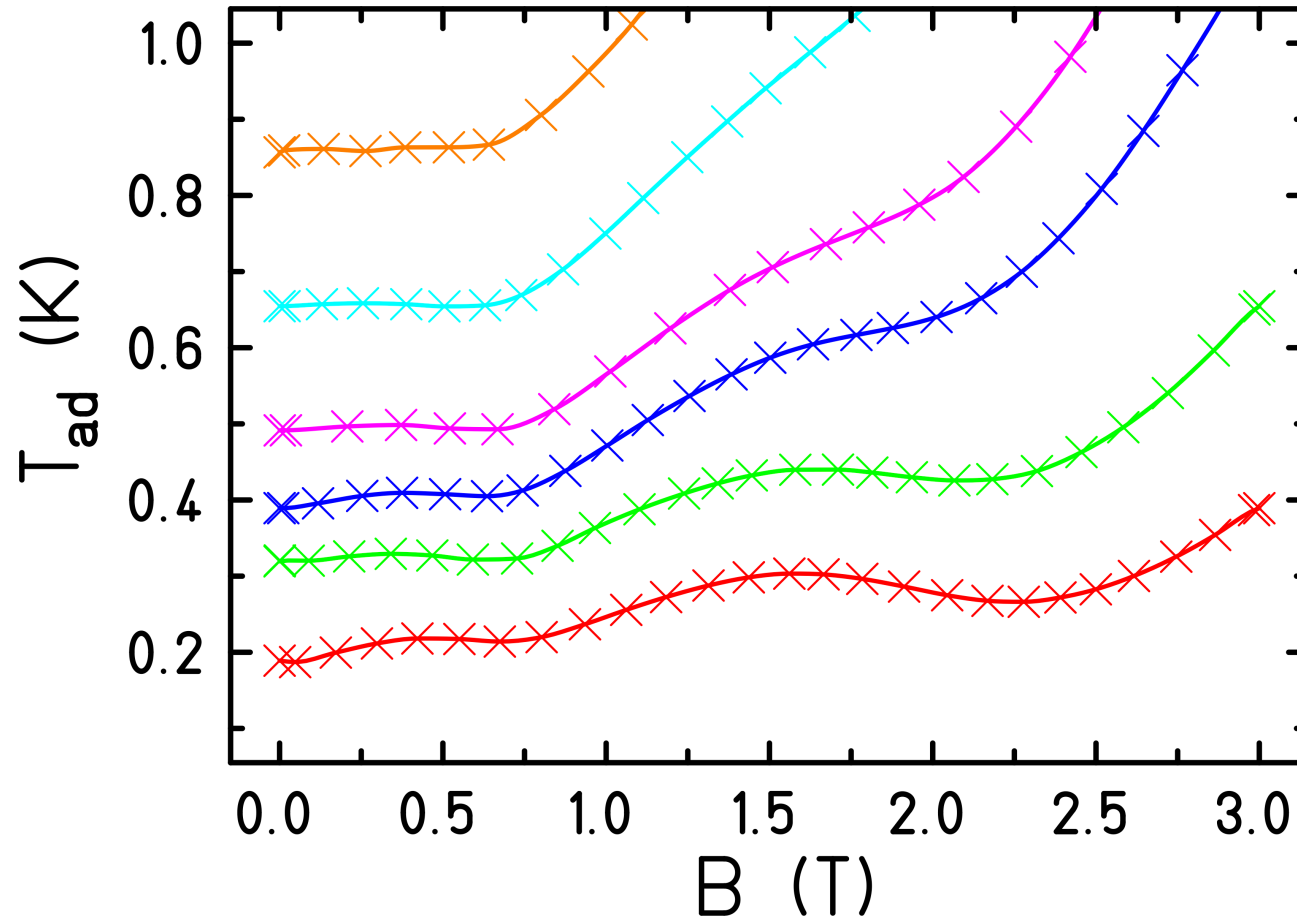
- Often magnetocaloric observables not directly measured, but inferred from Maxwell's relations.
- First real cooling experiment with a molecule.
- $$\underline{H} = -2 \sum_{i < j} J_{ij} \vec{\tilde{s}}_i \cdot \vec{\tilde{s}}_j + g \mu_B B \sum_i^N \tilde{s}_i^z$$

 $J_1 = -0.090(5) \text{ K}, J_2 = -0.080(5) \text{ K}$
 and $g = 2.02$.
- **Very good agreement down to the lowest temperatures.**

J. W. Sharples, D. Collison, E. J. L. McInnes, J. Schnack, E. Palacios, M. Evangelisti, Nat. Commun. **5**, 5321 (2014).



Gd₇ – Experimental cooling



J. W. Sharples, D. Collison, E. J. L. McInnes, J. Schnack, E. Palacios, M. Evangelisti, Nat. Commun. **5**, 5321 (2014).

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Thank you very much for your
attention.

The end.

Molecular Magnetism Web

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Highlights. Tutorials. Who is who. Conferences.