Modeling of isotropic and anisotropic magnetic molecules

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> Seminar Aachen, October 30th, 2009









Many thanks to my collaborators worldwide

- T. Englisch, T. Glaser, M. Höck, S. Leiding, A. Müller, R. Schnalle, Chr. Schröder, J. Ummethum (Bielefeld)
- K. Bärwinkel, H.-J. Schmidt, M. Neumann (Osnabrück);
- M. Luban, D. Vaknin (Ames Lab, USA); P. Kögerler (Aachen, Jülich, Ames) J. Musfeld (U. of Tennessee, USA); N. Dalal (Florida State, USA); R.E.P. Winpenny, E.J.L. McInnes (Man U, UK); L. Cronin (Glasgow, UK); J. van Slageren (Nottingham); H. Nojiri (Sendai, Japan); A. Postnikov (Metz, France)
- J. Richter, J. Schulenburg (Magdeburg); S. Blügel (FZ Jülich); A. Honecker (Göttingen); U. Kortz (Bremen); A. Tennant, B. Lake (HMI Berlin); B. Büchner, V. Kataev, R. Klingeler, H.-H. Klauß (Dresden); P. Chaudhuri (Mühlheim); J. Wosnitza (Dresden-Rossendorf)

Contents for you today



- 1. Introduction
- 2. Isotropic molecules: ITO & point group
- 3. Anisotropic molecules

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Introduction

Introduction

The beauty of magnetic molecules I



 Mn_{12}

- Inorganic or organic macro molecules, where paramagnetic ions such as Iron (Fe), Chromium (Cr), Copper (Cu), Nickel (Ni), Vanadium (V), Manganese (Mn), or rare earth ions are embedded in a host matrix;
- Pure organic magnetic molecules: magnetic coupling between high spin units (e.g. free radicals);
- Speculative applications: magnetic storage devices, magnets in biological systems, lightinduced nano switches, displays, catalysts, transparent magnets, qubits for quantum computers.

The beauty of magnetic molecules II



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- Dimers (Fe₂), tetrahedra (Cr₄), cubes (Cr₈);
- Rings, especially iron and chromium rings
- Complex structures (Mn₁₂) drosophila of molecular magnetism;
- "Soccer balls", more precisely icosidodecahedra (Fe₃₀) and other macro molecules;
- Chain like and planar structures of interlinked magnetic molecules, e.g. triangular Cu chain:

J. Schnack, H. Nojiri, P. Kögerler, G. J. T. Cooper, L. Cronin, Phys. Rev. B 70, 174420 (2004); Sato, Sakai, Läuchli, Mila, ...

The beauty of magnetic molecules III



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- Single Molecule Magnets (SMM): magnetic molecules with large ground state moment; e.g. S = 10 for Mn₁₂ or Fe₈
- Anisotropy barrier dominates behavior (as in your hard drive);
- Single molecule is a magnet and shows metastable magnetization and hysteresis; but also magnetization tunneling.
- Today's major efforts: improve stability of magnetization; rational design; investigate on surfaces.

Model Hamiltonian (spin only)

$$\begin{split} H &= \sum_{i,j} \vec{\underline{s}}(i) \cdot \mathbf{J}_{ij} \cdot \vec{\underline{s}}(j) + \sum_{i,j} \vec{D}_{ij} \cdot \left[\vec{\underline{s}}(i) \times \vec{\underline{s}}(j) \right] + \mu_B \vec{B} \sum_{i}^{N} \mathbf{g}_i \vec{\underline{s}}(i) \\ & \mathsf{Exchange/Anisotropy} \quad \mathsf{Dzyaloshinskii-Moriya} \quad \mathsf{Zeeman} \end{split}$$

Isotropic Hamiltonian

$$\begin{array}{lll} H &=& -\sum_{i,j} \,\, J_{ij} \, \vec{\underline{s}}(i) \cdot \vec{\underline{s}}(j) & + & g \, \mu_B \, B \, \sum_i^N \,\, \underline{\underline{s}}_z(i) \\ & &$$

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Thank God, we have computers



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Isotropic magnetic molecules



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Group theory for highly symmetric molecules:

- $H = -\sum_{i,j} J_{ij} \, \vec{s_i} \cdot \vec{s_j} + g \mu_B \vec{S} \cdot \vec{B}$;
- $\left[\underset{\sim}{H}, \overset{\overrightarrow{S}^2}{\underset{\sim}{\Sigma}} \right] = 0$, $\left[\underset{\sim}{H}, \overset{S}{\underset{\sim}{\Sigma}} z \right] = 0$;
- Irreducible Tensor Operator (ITO) approach, MAGPACK (1);
- Additional point group symmetries (2).

D. Gatteschi and L. Pardi, Gazz. Chim. Ital. **123**, 231 (1993); J. J. Borras-Almenar, J. M. Clemente-Juan, E. Coronado, and B. S. Tsukerblat, Inorg. Chem. **38**, 6081 (1999).
O. Waldmann, Phys. Rev. B **61**, 6138 (2000); V. E. Sinitsyn, I. G. Bostrem, and A. S. Ovchinnikov, J. Phys. A-Math. Theor. **40**, 645 (2007); R. Schnalle and J. Schnack, Phys. Rev. B **79**, 104419 (2009).

Reminder ITO

$$H_{\sim}^{H}\text{Heisenberg} = \sqrt{3} \sum_{i,j} J_{ij} T_{\sim}^{(0)}(\{k_i\}, \{\overline{k}_i\} | k_i = k_j = 1)$$

Irreducible Tensor Operator approach

- Express spin operators and functions thereof as ITOs;
- Use vector coupling basis $| \alpha S M \rangle$ and recursive recoupling;
- Numerical implementation e.g. MAGPACK.

(1) Gatteschi, Tsukerblat, Coronado, Waldmann, ...(2) R. Schnalle, Ph.D. thesis, Osnabrück University (2009)

Advanced ITO & Point Groups II

$$\mathcal{P}^{(n)} | \alpha S M \rangle = \left(\frac{l_n}{h} \sum_R \left(\chi^{(n)}(R) \right)^* \tilde{G}(R) \right) | \alpha S M \rangle$$

Point Group Symmetry

- Projection on irreducible representations (Wigner);
- Basis function generating machine;
- Orthonormalization necessary.
- (1) O. Waldmann, Phys. Rev. B **61**, 6138 (2000).
- (2) R. Schnalle, Ph.D. thesis, Osnabrück University (2009)

Advanced ITO & Point Groups III

$$\mathcal{G}(R) \mid \alpha \, S \, M \,\rangle_a = \sum_{\alpha'} \mid \alpha' \, S \, M \,\rangle_a \, {}_a \langle \, \alpha' \, S \, M \mid \alpha \, S \, M \,\rangle_b$$

Serious problem: Recoupling

- So far: only point groups that are compatible with the coupling scheme are used (1);
- Problem: otherwise complicated basis transformation between different coupling schemes;
- Solution: implementation of graph-theoretical results to evaluate recoupling coefficients $_a \langle \alpha' S M | \alpha S M \rangle_b$ (2).
- (1) O. Waldmann, Phys. Rev. B **61**, 6138 (2000).
- (2) R. Schnalle, Ph.D. thesis, Osnabrück University (2009)

Advanced ITO & Point Groups IV



Cuboctahedron, s = 3/2, Hilbert space dimension 16,777,216; symmetry O_h (1). Evaluation of recoupling coefficients very time consuming. (1,2)

(1) J. Schnack and R. Schnalle, Polyhedron **28**, 1620 (2009);

(2) R. Schnalle and J. Schnack, Phys. Rev. B **79**, 104419 (2009).

Summary – Isotropic molecules

- Combined use of SU(2) and point group symmetries possible.
- Numerical effort for recoupling coefficients enormous.
- Representation with smaller matrices.
- Further insight: spectroscopic labeling, selection rules for transitions.
- Works also for frustrated molecules, where QMC does not converge.



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Anisotropic magnetic molecules

Getting eigenvalues

$$H_{\widetilde{B}}(\vec{B}) = -\sum_{i,j} J_{ij} \vec{s}(i) \cdot \vec{s}(j) + \sum_{i} d_{i} \left(\vec{e}_{i} \cdot \vec{s}(i)\right)^{2} + \mu_{B} \vec{B} \cdot \sum_{i} \mathbf{g}_{i} \cdot \vec{s}(i)$$

- $\left[\frac{H}{\approx}, \vec{S}^2\right] \neq 0$, $\left[\frac{H}{\approx}, S_z\right] \neq 0$; \Rightarrow MAGPACK does not work!
- You have to diagonalize $H(\vec{B})$ for every field (direction and strength)!
- If you are lucky, point group symmetries still exist. Use them!
- (1) J. Schnack, Condens. Matter Phys. 12, 323 (2009);

How to obtain the magnetization?

Numerical differentiation

• For each field \vec{B} you evaluate the energy eigenvalues TWICE: $M_{\nu}(\vec{B}) = -\left(E_{\nu}(\vec{B}(1+\varepsilon)) - E_{\nu}(\vec{B})\right)/(\varepsilon B)$

 Numerical differentiation is a serious mathematical subject – good accuracy requires fine-tuning.

Using eigenvectors of $H(\vec{B})$

- Evaluate the energy eigenvectors: greater numerical effort, for INS anyway necessary;
- For each \vec{B} , evaluate and store E_{ν} and $\vec{M}_{\nu}(\vec{B}) = \mu_B \langle \nu | \sum_i^N \mathbf{g}_i \cdot \vec{s}(i) | \nu \rangle$.
- Accurate, but time consuming (eigenvectors!).

Orientational average for powder samples

- If you have a single crystal, doze off for the rest of my talk.
- Average over *x*-, *y*-, and *z*-direction: poor;
- Average over random directions: large fluctuations;
- Use Lebedev-Laikov grids: The parameters ensure that angular integration of polynomials $x^k \cdot y^l \cdot z^m$, where $k + l + m \leq 131$ can be performed with a relative accuracy of $2 \cdot 10^{-14}$.
 - I am using LLG with 50 (25) orientations.

V. I. Lebedev and D. N. Laikov, Dokl. Akad. Nauk 366, 741 (1999); and link to program on www.molmag.de



$Mn_3Cr I$

Mn₃Cr:

- Assume C_3 symmetry;
- Two couplings: J_1 to central Cr, J_2 between Mn;
- Model Mn anisotropy by local axis $\vec{e}(\vartheta, \phi)$. Due to C_3 symmetry $\vartheta_{Mn1} = \vartheta_{Mn2} = \vartheta_{Mn3}$. Only relative $\phi = 120^\circ$ determined.
- Model Cr anisotropy by local axis $\vec{e}(\vartheta, \phi)$. Due to C_3 symmetry $\vartheta_{Cr} = 0, \phi_{Cr} = 0$.
- Mn: s=5/2, g=2.0; Cr: s=3/2, g=1.95



Mn₃Cr II – Results



Result: $J_1 = -0.29 \text{ cm}^{-1}$, $J_2 = -0.07 \text{ cm}^{-1}$, $d_{Mn} = -1.05 \text{ cm}^{-1}$, $\vartheta_{Mn} = 15^{\circ}$, $d_{Cr} = +0.40 \text{ cm}^{-1}$.

M. Prinz, K. Kuepper, C. Taubitz, M. Raekers, B. Biswas, T. Weyhermüller, M. Uhlarz, J. Wosnitza, J. Schnack, A. V. Postnikov, C. Schröder, S. J. George, M. Neumann, P. Chaudhuri, Inorg. Chem., still struggling with the referees.

Mn₃Cr III – Angular averaging



For a good fit you need several directions, at least 10.

$\mathbf{Mn}_{6}\mathbf{Fe}\ \mathbf{I}$



- Rational design of strict C₃ symmetry of local easy axes (Thorsten Glaser): e.g. Mn₆Cr (1), Mn₆Fe (2)
- Mn₆Fe: J_1 between Mn in caps, J_2 to central Fe; Mn anisotropy modeled by local axis $\vec{e}(\vartheta, \phi)$ with $\vartheta_{Mn1} = \vartheta_{Mn2} = \vartheta_{Mn3} = 36.5^{\circ}$. Only relative $\phi = 120^{\circ}$ determined.
- Mn: s=2, g=1.98; Fe: s=1/2, g-tensor.

 T. Glaser, M. Heidemeier, T. Weyhermüller, R. D. Hoffmann, H. Rupp, P. Müller, Angew. Chem.-Int. Edit. 45, 6033 (2006).
T. Glaser, M. Heidemeier, E. Krickemeyer, H. Bögge, A. Stammler, R. Fröhlich, E. Bill, J. Schnack, Inorg. Chem. 48, 607 (2009).

Mn₆Fe II – Results



T. Glaser, M. Heidemeier, E. Krickemeyer, H. Bögge, A. Stammler, R. Fröhlich, E. Bill, J. Schnack, Inorg. Chem. **48**, 607 (2009).

$[V_4^{III}CI_6(thme)_2(bipy)_3]$

- 4 V_4^{III} ions with s = 1; approximate C_3 symmetry;
- 2 exchange interactions;
- Central V: axial anisotropy;
- Outer Vs: local anisotropy axis with azimuthal angle ϑ .

• Powder average.

Ian S. Tidmarsh, Luke J. Batchelor, Emma Scales, Rebecca H. Laye, Lorenzo Sorace, Andrea Caneschi, Jürgen Schnack and Eric J.L. McInnes, Dalton Trans. (2009) 9402



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Two equally good parameter sets.

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 V_4 III



High fields could distinguish.

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Accuracy of measurement limits modeling.

V_4 – Anisotropy tensors



Cartoon of anisotropy tensors.

Summary – Anisotropic molecules

- It is possible to determine local anisotropy axes with rather high accuracy.
- Complementary *ab initio* calculations on local D-tensors would be valuable.
- Powders have to be averaged properly.
- Single crystals would probably allow to obtain the full local D-tensor, i.e. also *E*-terms.



• Element-selective calculations possible.

Thank you very much for your attention.

Information

Molecular Magnetism Web

www.molmag.de

Highlights. Tutorials. Who is who. Conferences.