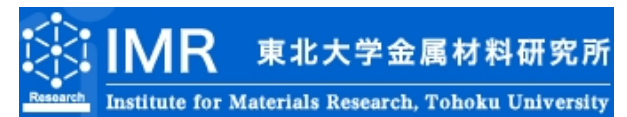
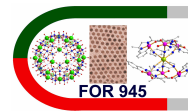


Es geht auch exakt! – Meine mathematischen Lehrjahre mit Heinz-Jürgen Schmidt

Jürgen Schnack

Fakultät für Physik – Universität Bielefeld
<http://obelix.physik.uni-bielefeld.de/~schnack/>

Festkolloquium, 27. Januar 2012



Legendenbildung*

* Bei der Legendenbildung ist es wichtig, dass man störenden Details nicht allzuviel Aufmerksamkeit zukommen lässt!

etwa um 2004 \pm 2

Ganz Germanien ist von
Evaluatoren,
Drittmitteleintreibern und
Akkreditoren besetzt!

Ganz Germanien ???

Dem Fachbereich Physik
der Universität Osnabrück
droht in Folge personeller
Auszehrung der Untergang!

**Doch eine kleine Gruppe
Ausserplanmäßiger Professoren
leistet tapfer Widerstand!**



Klaus Betzler

**Sprecher des Graduiertenkollegs und
Spiritus Rector der Betzler- und Master-Studiengänge**



Jürgen Schnack & Manfred Neumann Sprecher des Promotionsprogramms



Heinz-Jürgen Schmidt
Dekan



Wir haben so hart gearbeitet, dass wir oft nicht einmal dazu gekommen sind, den Mantel auszuziehen!



Unsere gute Laune hat darunter nicht gelitten.

Wir hatten was geschafft!

Marshall Luban about Heinz-Jürgen Schmidt



My *in-total summary* of Heinz-Jürgen is that he is one of the smartest, perhaps **THE SMARTEST**, mathematical physicist that I have met in my life!

Marshall Luban at the age of 75!

Paul Busch über Heinz-Jürgen Schmidt



Lieber Herr Schnack,
Ich kann noch hinzufügen, dass die Arbeit mit Heinz-Jürgen zurück zu den Anfängen in Grundlagen der Quantentheorie bei Ludwig geführt hat; es war beeindruckend zu sehen, wie leicht ihm das ganze Arsenal wieder zuflog oder einfach noch zur Verfügung steht.
Wenn Sie etwas technisches einfügen möchten, könnte ich vielleicht noch eine komplizierte Formel beibringen ...
Herzliche Grüße, Ihr Paul Busch

Paul Busch, York, U.K.

Heinz-Jürgen Schmidt



1. **Legendenbildung** ✓
2. Ein paar komplizierte Formeln,
soweit ich das verantworten kann
3. **Persönliche Anmerkungen**

Wenn Sie jetzt nichts verstehen,
ist's nicht so schlimm . . .

... ich versteh ja auch nichts!

Theorem 1. $(M_4, \langle | \rangle, \prec, \prec_o)$ is isomorphic to the 4-dimensional Minkowski space.

Accordingly, we will apply freely the terminology of Minkowski geometry and refer to $\langle | \rangle$ as the (Minkowski) scalar product. We use the same notation for vectors and for points in M_4 as an affine space.

The *forward* and *backward light cones* of an element $x \in M_4$ are defined as the sets

$$\mathcal{F}(x) := \{y \in M_4 : x \prec_o y\}, \quad \mathcal{B}(x) := \{y \in M_4 : x \succ_o y\}.$$

A vector $x \in M_4$ is called *lightlike* if $\langle x | x \rangle = 0$. If $\langle x | x \rangle > 0$ or < 0 , the vector x is called *timelike* or *spacelike*, respectively. Then $x \prec_o y$ is equivalent to $y - x$ being lightlike and $y_0 - x_0 \geq 0$. Elements $x, y \in M_4$ are called *spacelike related*, $x \sigma y$, if $\langle x - y | x - y \rangle < 0$.

The set of effects can now be written as $(\text{conv}(X))$ denotes the convex hull of a set X)

$$\mathcal{L} = [\mathbb{0}, \mathbb{1}] = \text{conv}(\mathcal{F}(\mathbb{0})) \cap \text{conv}(\mathcal{B}(\mathbb{1}))$$

\mathcal{L} is convex and compact, that is, it includes its boundary $(\mathcal{F}(\mathbb{0}), \mathcal{B}(\mathbb{1}))$.

The Minkowski scalar product $\langle e | f \rangle$ admits a simple interpretation: if e and f are effects: it is equal to the probability of joint effects e and f are measured by, say, Alice and Bob on an entangled (singlet) state $\Phi = \frac{1}{\sqrt{2}}(\psi_+ \otimes \psi_- - \psi_- \otimes \psi_+)$.

COEXISTENCE OF QUBIT EFFECTS
 PAUL BUSCH AND HEINZ-JÜRGEN SCHMIDT

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Structuralism in Physics

First published Sun Nov 24, 2002; substantive revision Wed Mar 5, 2008

Under the heading of “structuralism in physics” there are three different but closely related research programs in philosophy of science and, in particular, in philosophy of physics. These programs were initiated by the work of Joseph Sneed, Günther Ludwig, and Erhard Scheibe, respectively, since the begin of the 1970s. For the sake of simplicity we will use these names in order to refer to the three programs, without the intention of ignoring or minimizing the contributions of other scholars. (See the Bibliography.) The term “structuralism” was originally claimed by the Sneed school, see e.g., Balzer and Moulines (1996), but it also appears appropriate to subsume Ludwig's and Scheibe's programs under this title because of the striking similarities of the three approaches. The activities of the structuralists have been mainly confined to Europe, especially Germany, and, for whatever reasons, largely ignored in the Anglo-American discussion.

- 1. Common traits
- 2. The problem of theoretical terms
 - 2.1 An example
 - 2.2 Structuralistic solutions of the problem of theoretical terms
 - 2.3 The measurement problem
 - 2.4 Measurement and approximation
- 3. Problems of reduction
 - 3.1 Reduction relation between theories
 - 3.2 Reduction and incommensurability
 - 3.3 Ludwig's account
 - 3.4 Sneed's account
 - 3.5 Scheibe's account
- 4. Three structuralist programs
 - 4.1 Sneed's program
 - 4.2 Ludwig's program
 - 4.3 Scheibe's program
- 5. Summary
- Bibliography
 - Sneed's program
 - Ludwig's program
 - Scheibe's program
- Other Internet Resources
- Related Entries

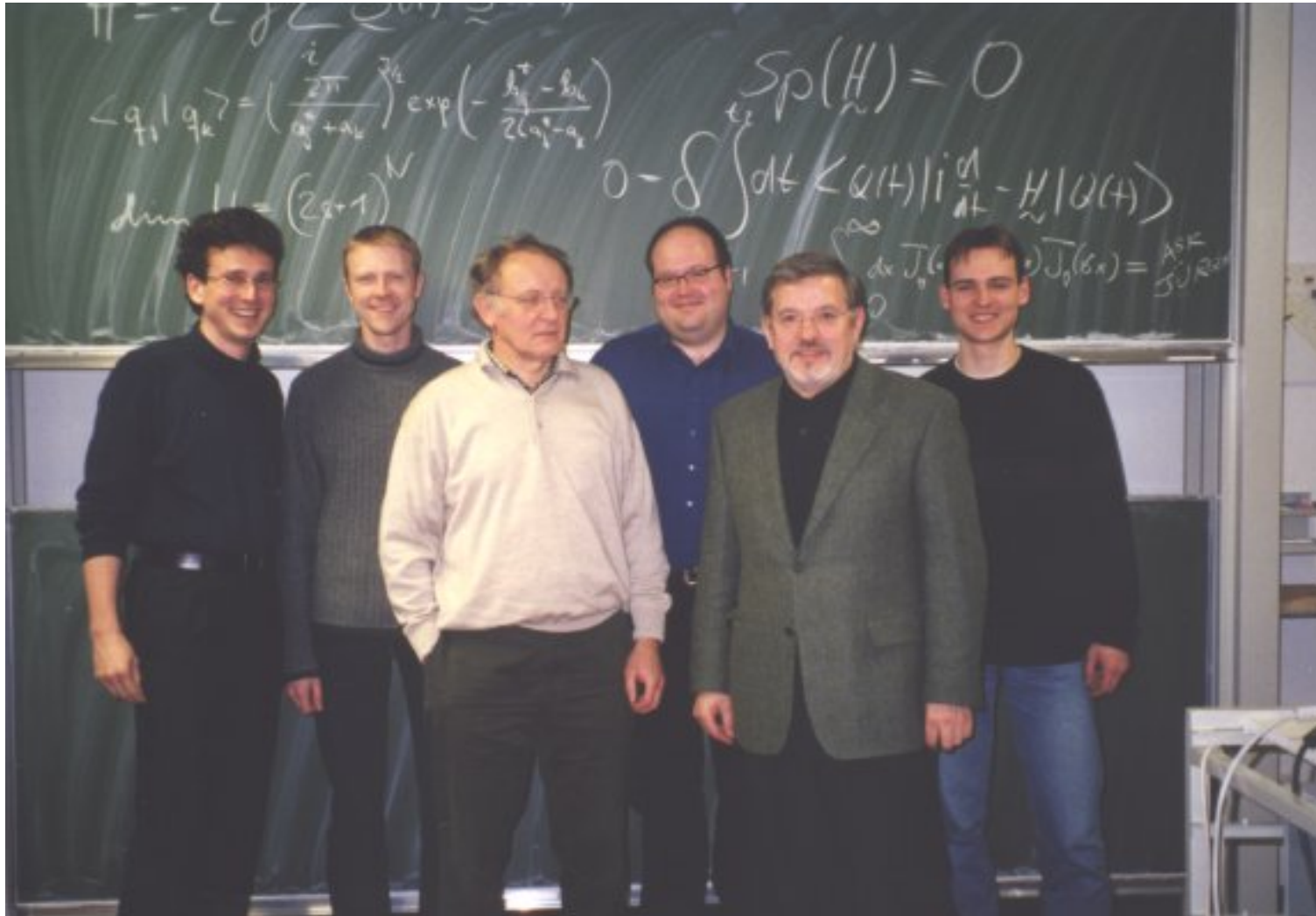
Acknowledgments

The author is indebted to John D. Norton, Edward N. Zalta, and Susanne Z. Riehemann for helpful suggestions concerning the content and the language of this entry.

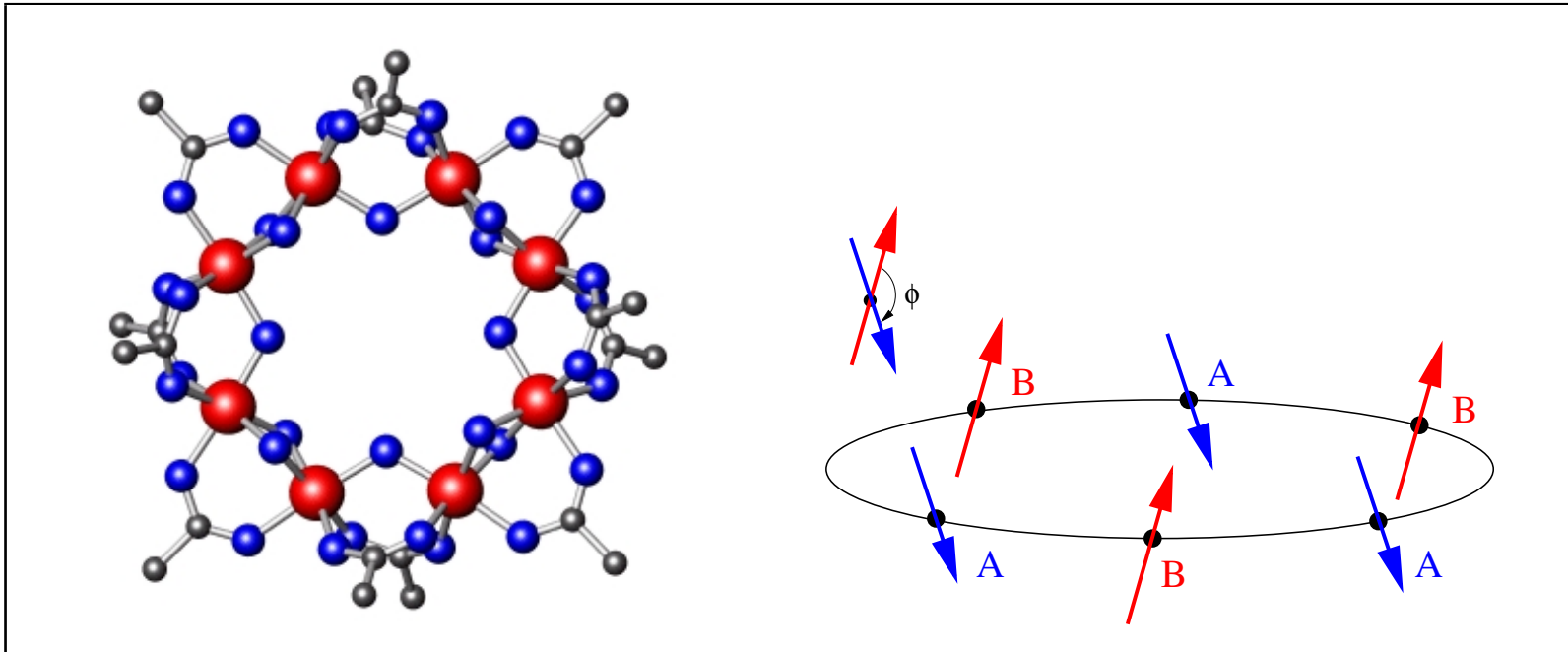
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HEINZ-JUERGEN SCHMIDT <HSCHMIDT@PHYSIK.UNI-OSNABRUECK.DE>

Arbeitsgruppe im November 2001



Magnetische Moleküle

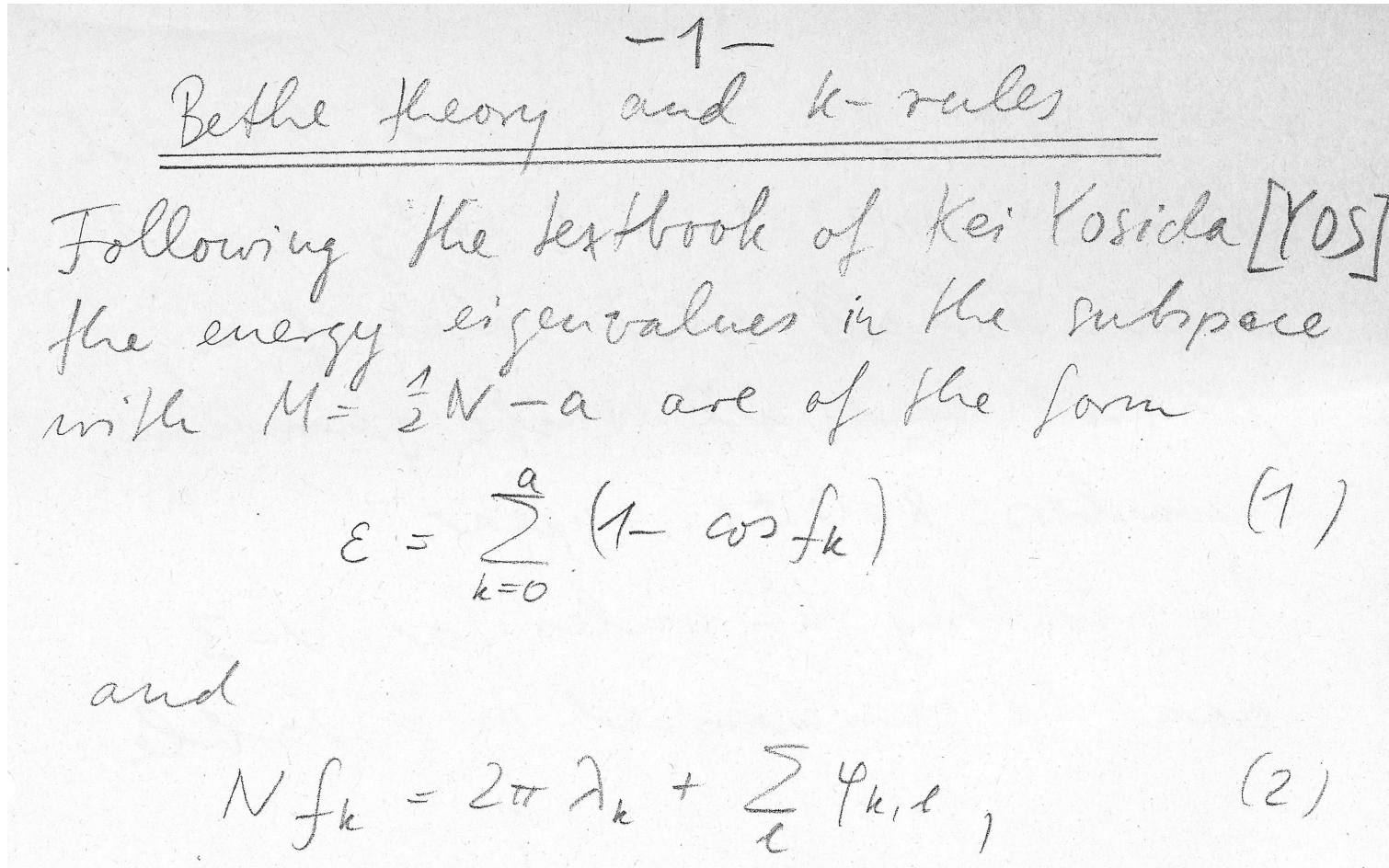


- Heisenberg-Modell (1): $\underline{H} = - \sum_{i,j} J_{ij} \underline{\tilde{s}}(i) \cdot \underline{\tilde{s}}(j);$
- Der molekulare Magnetismus hat uns im 4. Stock zusammengebracht (2).

(1) W. Heisenberg, Z. f. Phys. **49**, 619 (1928)

(2) K. Bärwinkel, H.-J. Schmidt, and J. Schnack, J. Magn. Magn. Mater. **212**, 240 (2000)

Originalnotizen



Arbeitsgruppe im August 2002



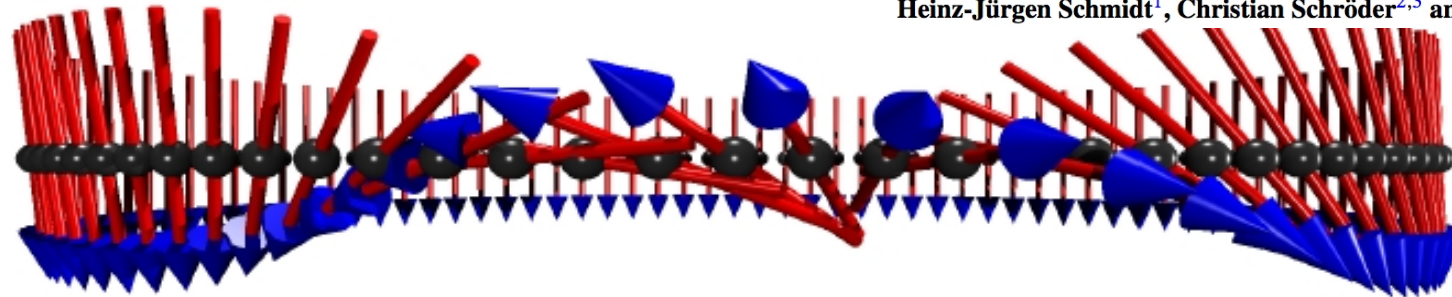
Du konntest mit jedem – 2004



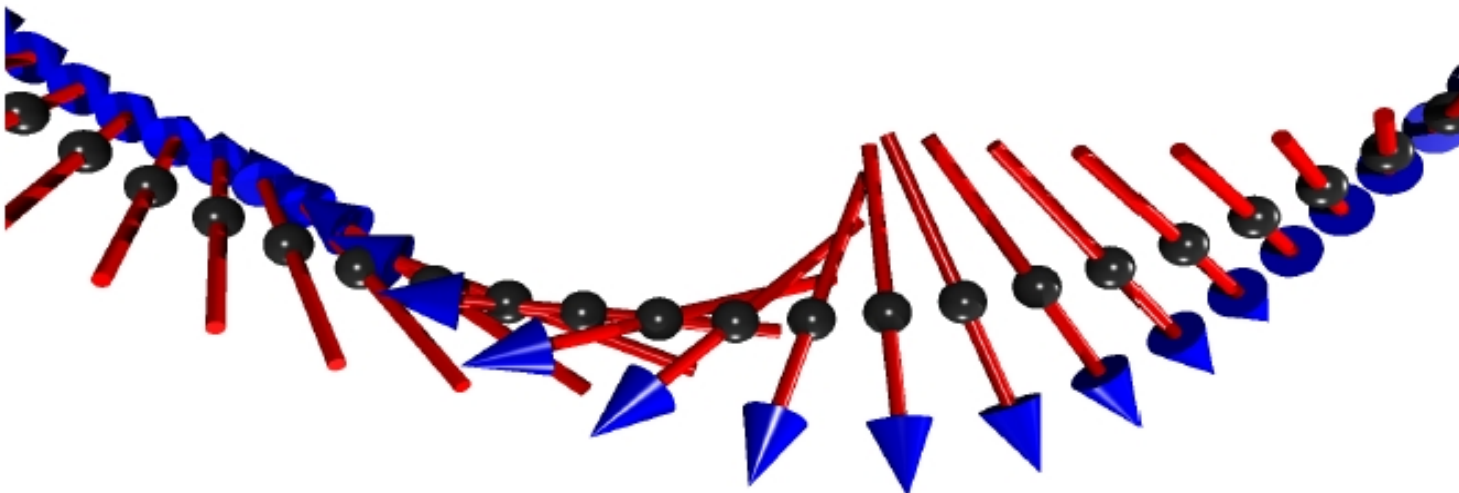
Modulated spin waves and robust quasi-solitons in classical Heisenberg rings

Heinz-Jürgen Schmidt¹, Christian Schröder^{2,3} and Marshall Luban^{3,4}

a)



b)





Thermodynamic fermion–boson symmetry in harmonic oscillator potentials

H.-J. Schmidt, J. Schnack *

Using the usual dependence on the inverse temperature β this property expresses itself as

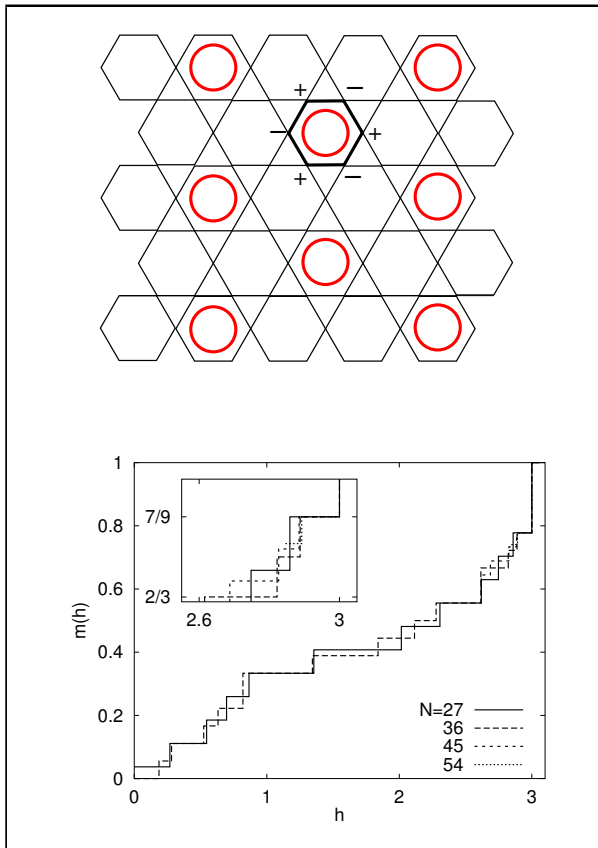
$$Z_N^+(\beta) = (-1)^N Z_N^-(-\beta), \tag{11}$$

where the partition function with the negative argument has to be understood as the analytic continuation into the region of $y = \exp(-\beta\hbar\omega) > 1$. In thermodynamic mean values like mean energy or specific heat this symmetry also shows up

$$E_N^+(\beta) = -E_N^-(-\beta), \quad C_N^+(\beta) = C_N^-(-\beta). \tag{12}$$

A straightforward application of the above result is to calculate fermionic partition functions and mean values by evaluating the respective bosonic ones at negative temperatures and thereby to avoid the sign problem.

Giant magnetization jumps in frustrated antiferromagnets III Kagome Lattice



- Non-interacting one-magnon states can be found on various lattices, e. g. kagome or pyrochlore
- Each state of n independent magnons is a state in the Hilbert subspace of dimension 2^n
Kagome: max. number of independent magnons is 12
- Linear dependence of magnetization on magnetic field \Rightarrow (T) **anomalous magnetization effect!**
- Jump in magnetization at $h = 2/3, 7/9, 3$
- A rare states!

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22 APRIL 2002

Macroscopic Magnetization Jumps due to Independent Magnons in Frustrated Quantum Spin Lattices

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 (Received 29 August 2001; published 8 April 2002)

J. Schnack, H.-J. Schmidt, J. Richter, J. Schulenburg, Eur. Phys. J. B **24**, 475 (2001)

J. Schulenburg, A. Honecker, J. Schnack, J. Richter, H.-J. Schmidt, Phys. Rev. Lett. **88**, 167207 (2002)

Persönliche Anmerkungen

Engagierte Lehre



voller Einsatz – selbstkritisch – kommunikativ

Workshop 2005



Ehrendoktorwürde für Marshall Luban 2006





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van der Waals revisited

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4. Summary

The summary of our elaboration is that for a large region of temperatures, the functions $\langle\langle F_{\pm} \rangle\rangle = a$ and $\langle\langle G_{\pm} \rangle\rangle = b$ may be nearly considered as constants, i.e. not dependant on temperature, see Fig. 1. Moreover, the respective kind of statistics (Bose–Einstein or Fermi–Dirac) does not matter. This constitutes a fermion-boson symmetry in non-ideal quantum gases. For ideal quantum gases such symmetries have been known for about ten years [8,9].

The slight residual dependence of a and b on temperature (cf. Fig. 1) reflects the fact that these are parameters not of an exact but of a model theory.

Acknowledgments

A long-standing collaboration with our colleague and friend Heinz–Jürgen Schmidt is gratefully acknowledged. This article is dedicated to him on the occasion of his 60th birthday.

Lieber Heinz-Jürgen,
in Zeiten, in denen jeder*
Daten fittet bis der Arzt kommt,

* Falls Sie es nicht selbst merken,
hier ist auch etwas Selbstkritik verborgen!

... hast Du mich fragen und
denken gelehrt!

Vielen Dank für 9 schöne Jahre
1997 – 2006

Ich bin sicher, Dein Ruhestand
wird zwar ruhig, aber intellektuell
voll ausgefüllt sein.

Alles Gute!