

Magnetism of the $N = 42$ kagome lattice antiferromagnet

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The Bielefeld conspiracy

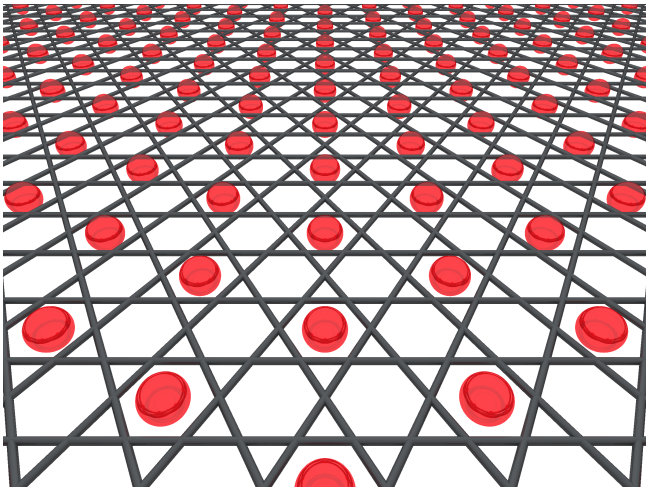


The story goes that the city of **BIELEFELD** in the German state of North Rhine-Westphalia **DOES NOT** actually **EXIST**. Rather, its existence is merely propagated by an entity known only as **THEM**, which has conspired with the authorities to create the illusion of the city's existence.

The origins of and reasons for this conspiracy are not a part of the original theory. Speculated originators jokingly include the CIA, Mossad, or aliens who use Bielefeld University as a disguise for their spaceship.

Do you know anybody from Bielefeld?

Kagome lattice antiferromagnet – the problem



- Thermodynamic functions, in particular heat capacity and susceptibility.
- Magnetization curve, in particular thermal stability of plateau at $\mathcal{M}_{\text{sat}}/3$.
- Method: Finite-temperature Lanczos.
- Comparison with tensor-network calculations.

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Model Hamiltonian

$$\tilde{H} = J \sum_{i < j} \vec{\tilde{s}}_i \cdot \vec{\tilde{s}}_j + g \mu_B B \sum_i^N \tilde{s}_i^z$$

Heisenberg
Zeeman

Finite-temperature Lanczos Method I

$$Z(T, B) = \sum_{\nu} \langle \nu | \exp \left\{ -\beta \tilde{H} \right\} | \nu \rangle$$

$$\langle \nu | \exp \left\{ -\beta \tilde{H} \right\} | \nu \rangle \approx \sum_n \langle \nu | n(\nu) \rangle \exp \left\{ -\beta \epsilon_n \right\} \langle n(\nu) | \nu \rangle$$

$$Z(T, B) \approx \frac{\dim(\mathcal{H})}{R} \sum_{\nu=1}^R \sum_{n=1}^{N_L} \exp \left\{ -\beta \epsilon_n \right\} |\langle n(\nu) | \nu \rangle|^2$$

- $|n(\nu)\rangle$ n-th Lanczos eigenvector starting from $|\nu\rangle$
- Partition function replaced by a small sum: $R = 1 \dots 10, N_L \approx 100$.

J. Jaklic and P. Prelovsek, Phys. Rev. B **49**, 5065 (1994).

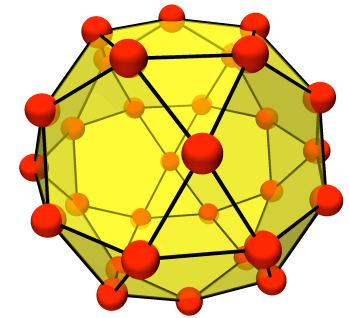
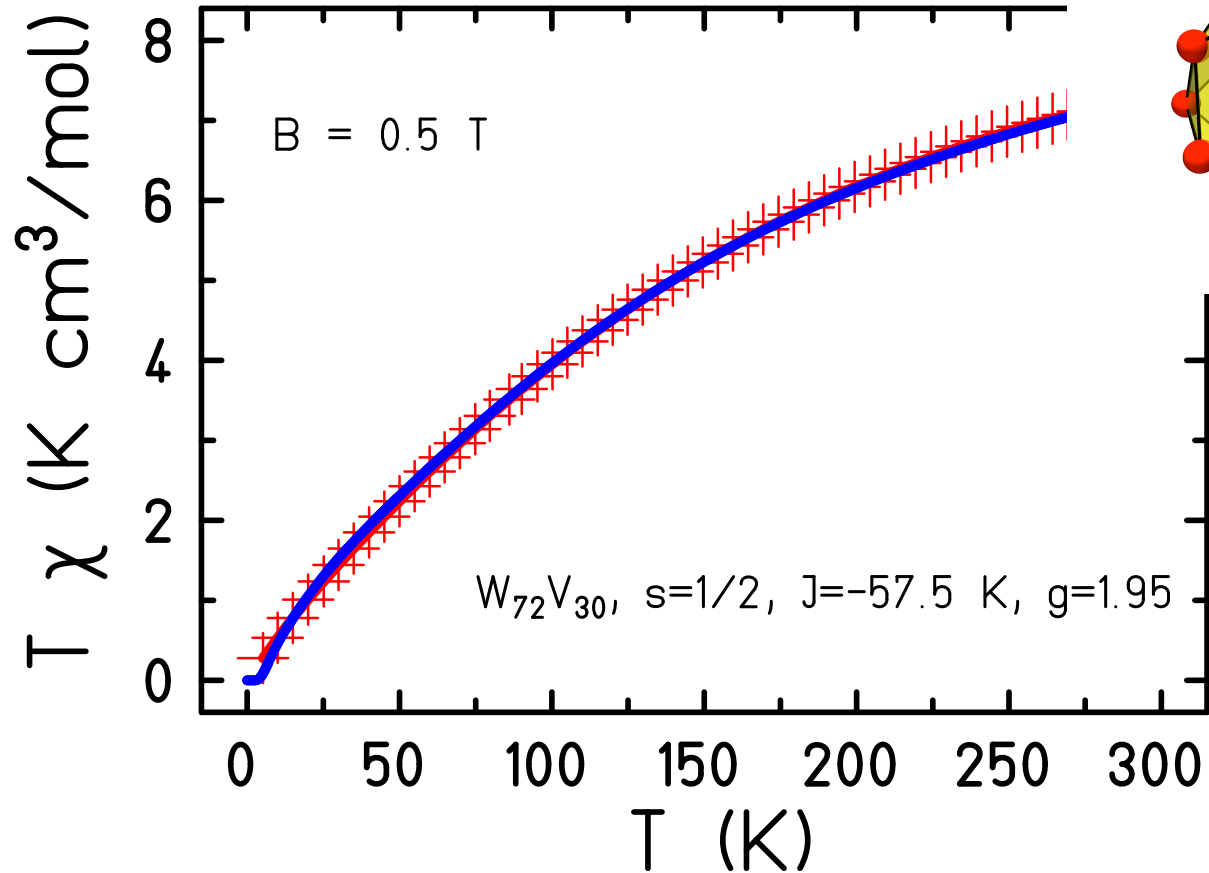
Finite-temperature Lanczos Method II

$$Z(T, B) \approx \sum_{\Gamma} \frac{\dim(\mathcal{H}(\Gamma))}{R_{\Gamma}} \sum_{\nu=1}^{R_{\Gamma}} \sum_{n=1}^{N_L} \exp\{-\beta\epsilon_n\} |\langle n(\nu, \Gamma) | \nu, \Gamma \rangle|^2$$

- Approximation better if symmetries taken into account.
- Γ denotes the used irreducible representations; often this is just the \tilde{S}^z symmetry, i.e. $\Gamma \equiv M$

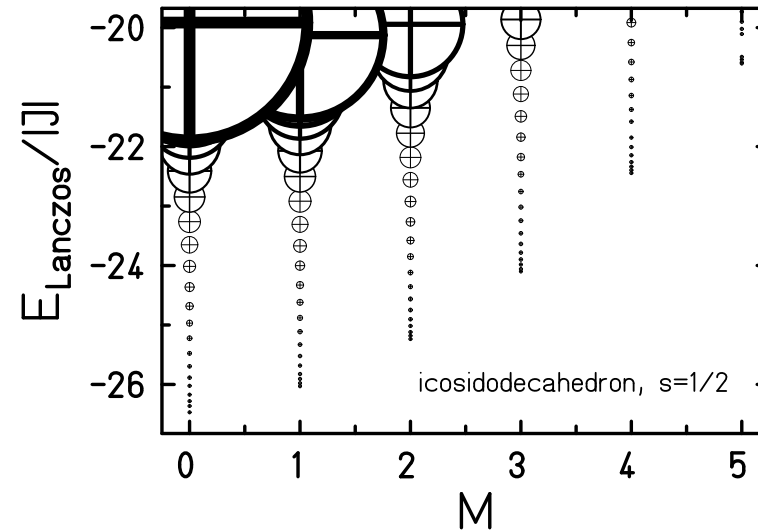
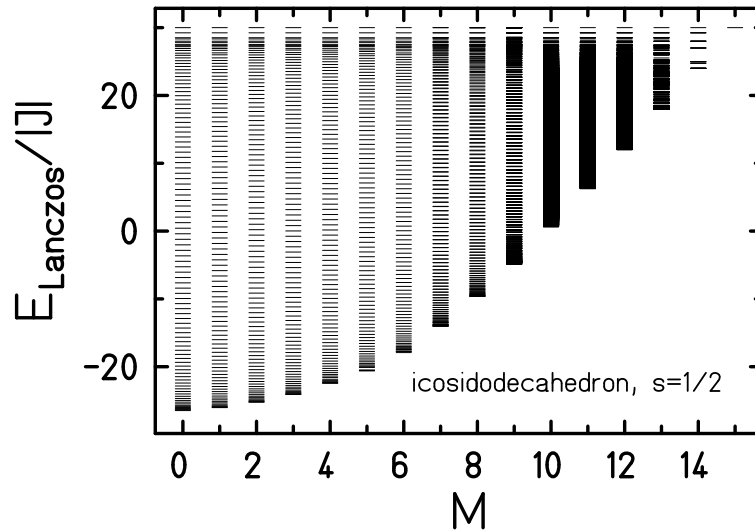
J. Schnack and O. Wendland, Eur. Phys. J. B **78** (2010) 535-541

Icosidodecahedron $s = 1/2$



Exp. data: A. M. Todea, A. Merca, H. Bögge, T. Glaser, L. Engelhardt, R. Prozorov, M. Luban, A. Müller, Chem. Commun., 3351 (2009).

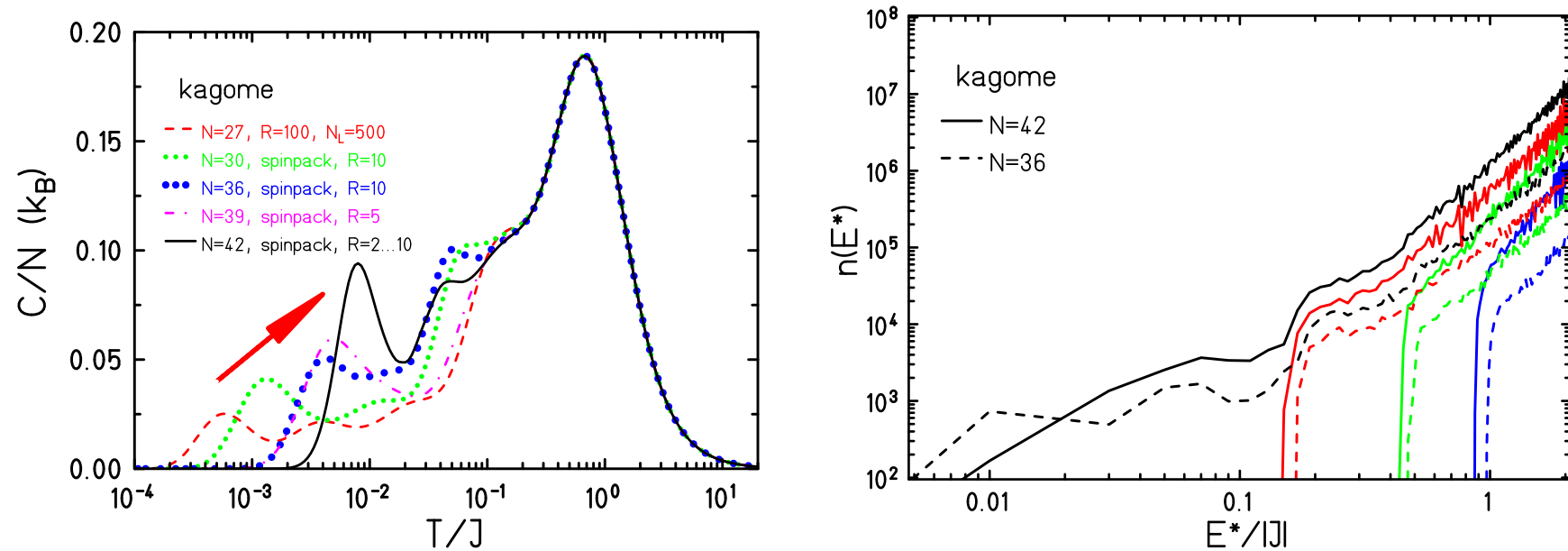
Icosidodecahedron $s = 1/2$



- The true spectrum will be much denser. This is miraculously compensated for by the weights.

$$Z(T, B) \approx \frac{\dim(\mathcal{H})}{R} \sum_{\nu=1}^R \sum_{n=1}^{N_L} \exp\{-\beta \epsilon_n\} |\langle n(\nu, \Gamma) | \nu, \Gamma \rangle|^2$$

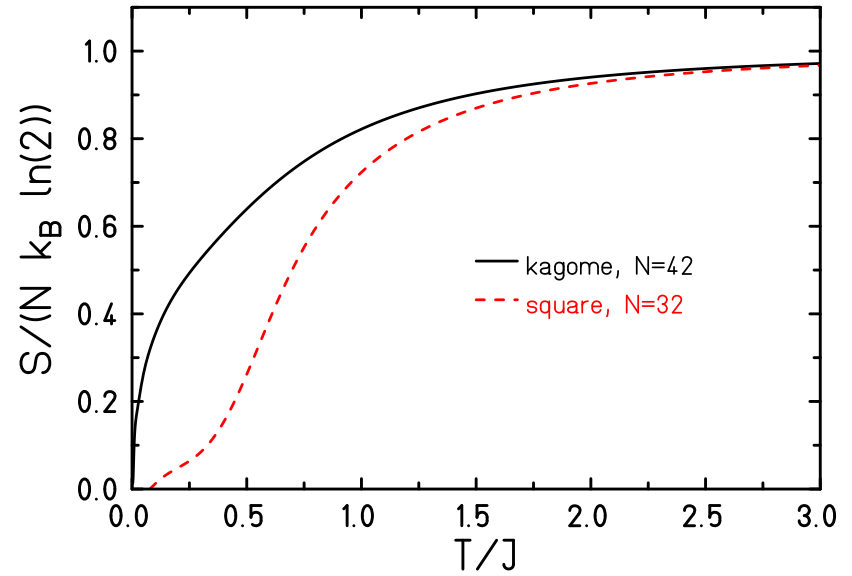
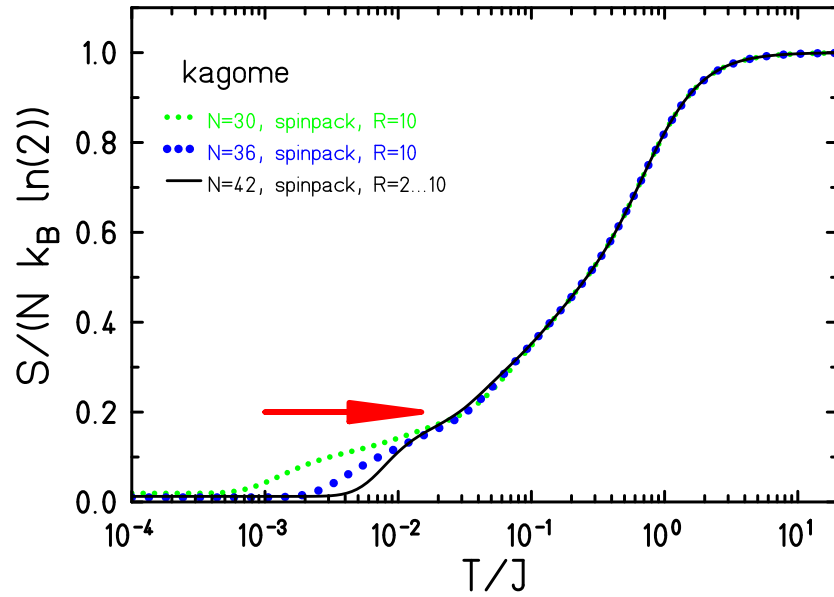
Kagome 42 – heat capacity



- Low- T peak moves to higher T with increasing N .
- Density of low-lying singlets seems to move to higher excitation energies.

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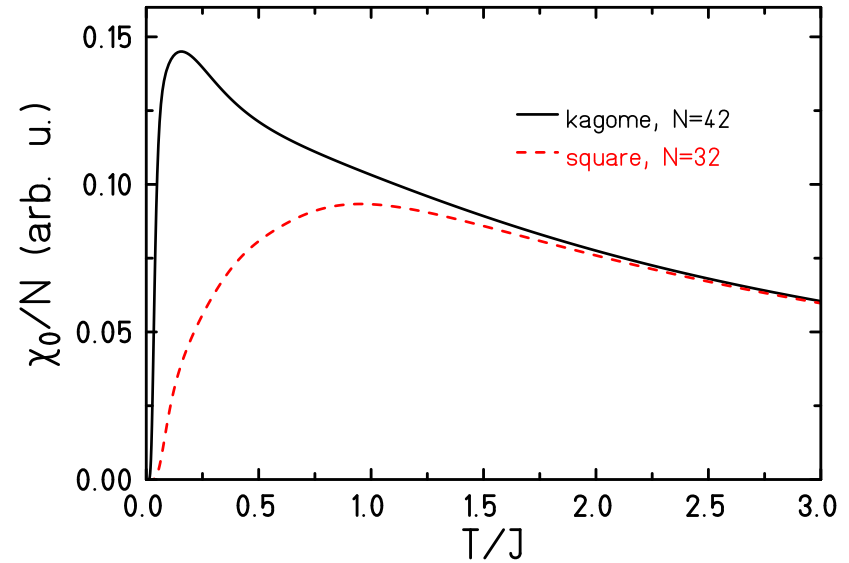
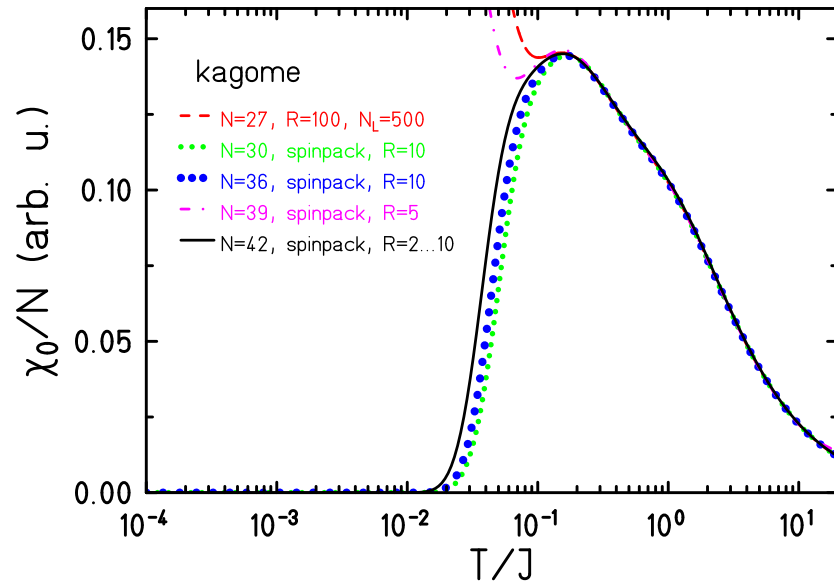
Kagome 42 – entropy



- Rise of entropy for higher T with increasing N .

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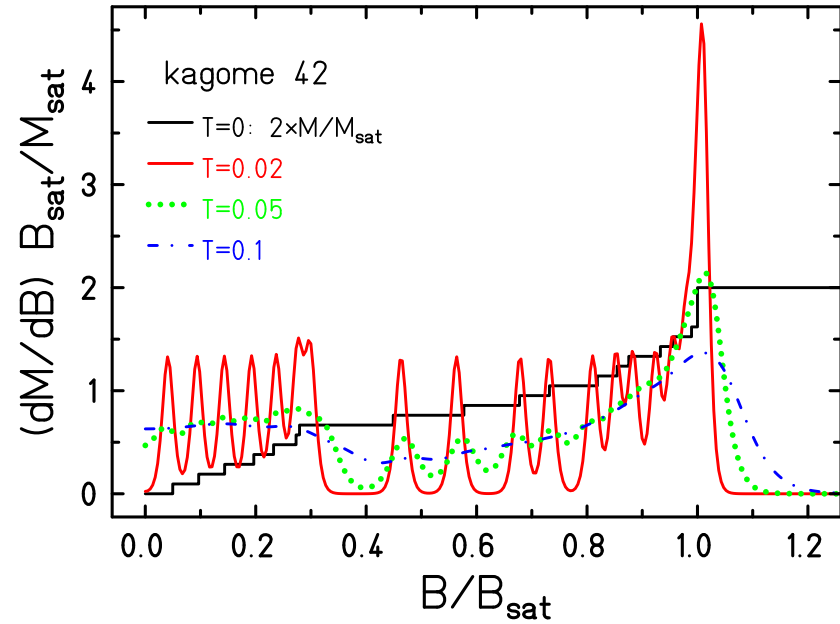
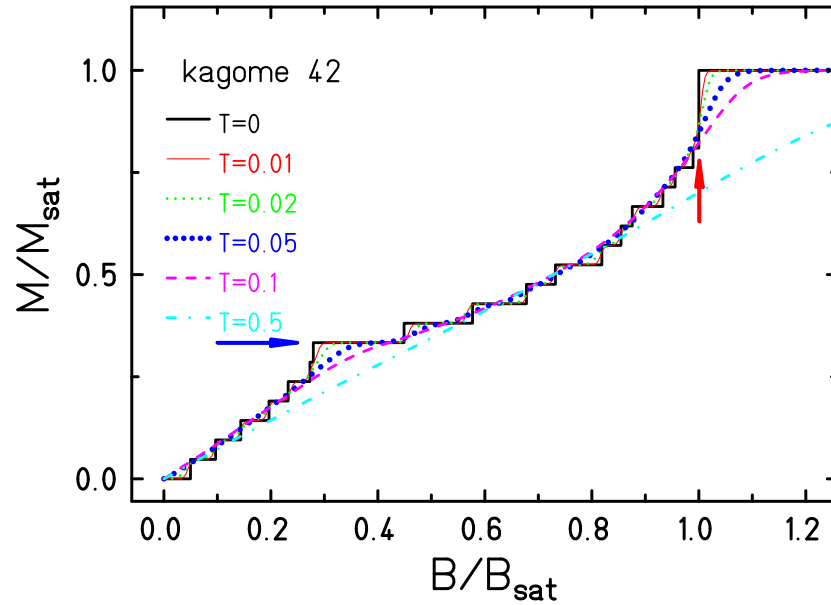
Kagome 42 – susceptibility



- Singlet-triplet gap shrinks very slowly with increasing N .

(1) A. Laeuchli, J. Sudan, and R. Moessner, arXiv:1611.06990.
 (2) J. Schnack, J. Schulenburg, J. Richter, Phys. Rev. B **98** (2018) 094423

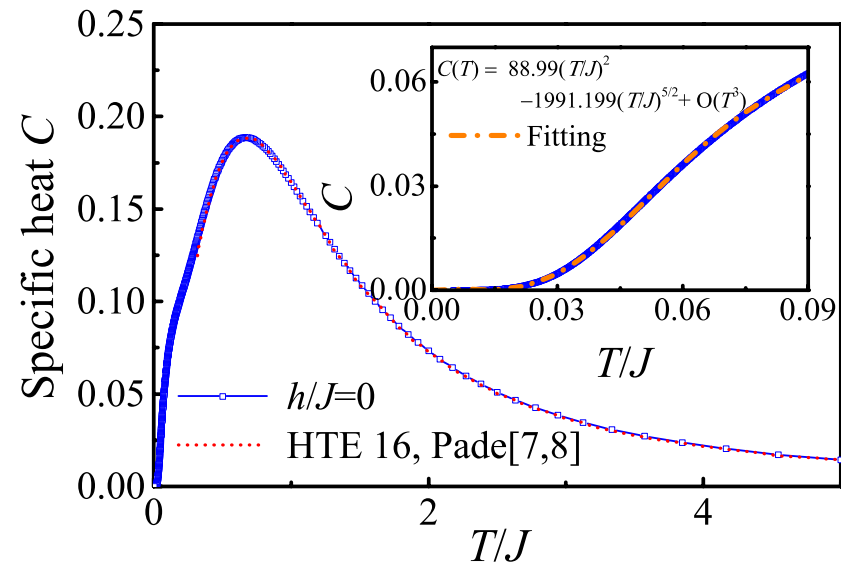
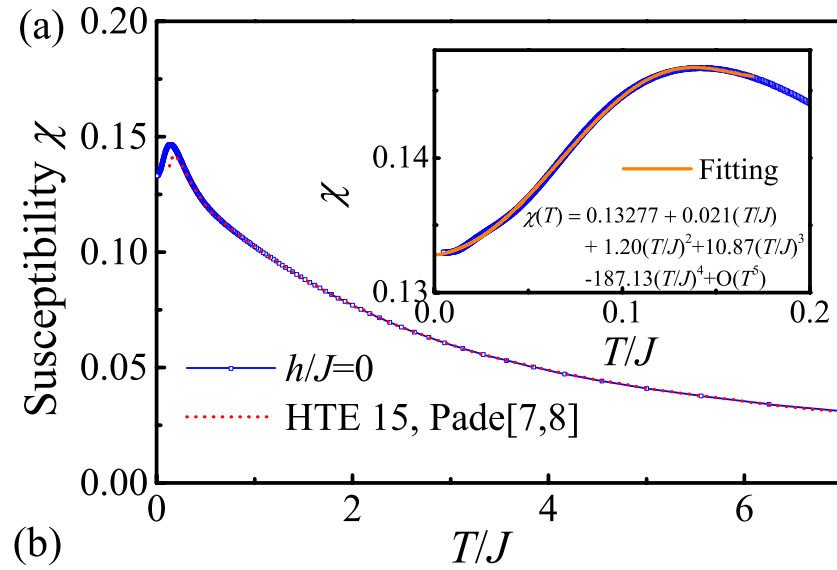
Kagome 42 – magnetization



- Plateaus and jump; asymmetric melting of the plateau at $\mathcal{M}_{\text{sat}}/3$.

(1) S. Capponi, O. Derzhko, A. Honecker, A. M. Laeuchli, J. Richter, Phys. Rev. B **88**, 2 144416 (2013).
 (2) J. Schulenburg, A. Honecker, J. Schnack, J. Richter, H.-J. Schmidt, Phys. Rev. Lett. **88**, 167207 (2002).
 (3) H. Nakano and T. Sakai, J. Phys. Soc. Jpn. **83**, 104710 (2014).

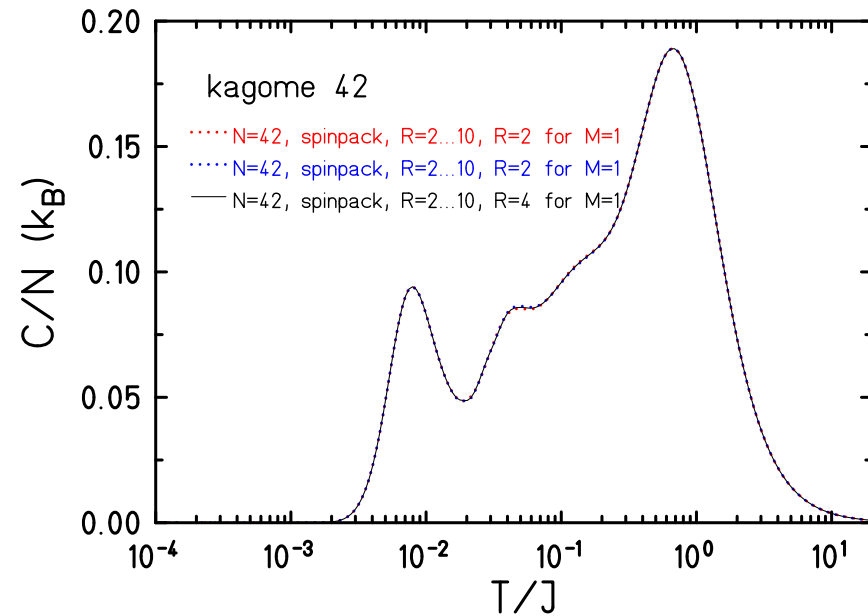
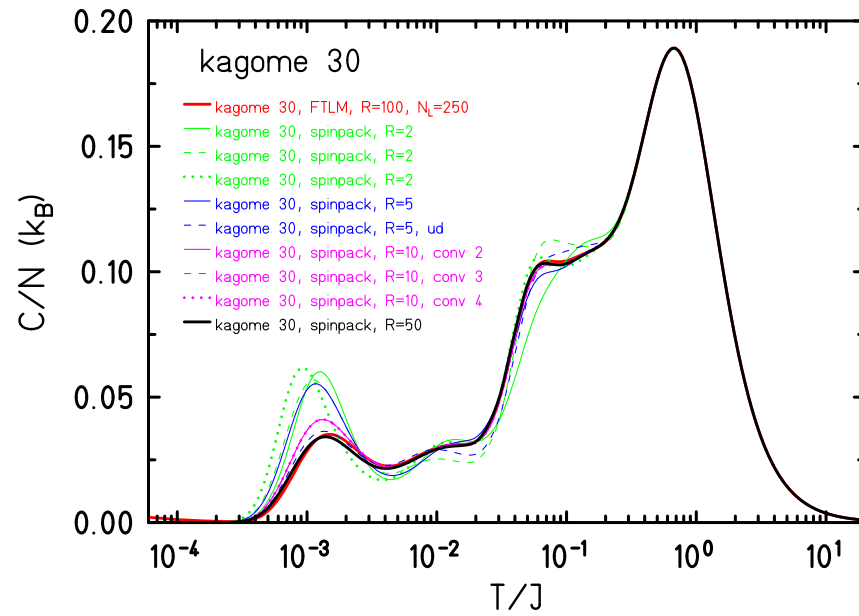
Kagome – tensor network calculations



- Tensor network calculations for the infinite system (1).

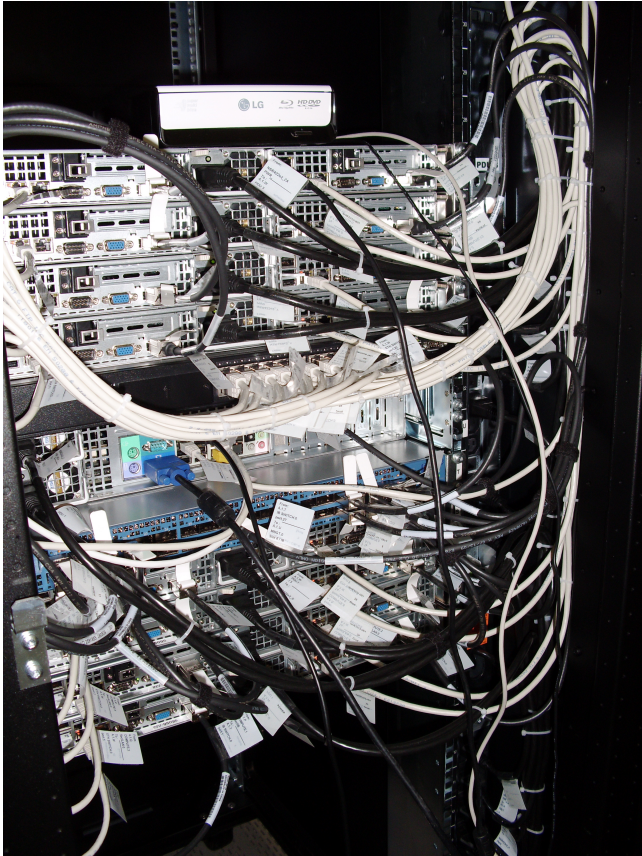
(1) Xi Chen, Shi-Ju Ran, Tao Liu, Cheng Peng, Yi-Zhen Huang, Gang Su, Science Bulletin **63**, 1545 (2018).

Error estimates of FTLM for the kagome calculations



FTLM more accurate for larger dimensions of Hilbert space.

Summary



- Largest FTLM calculation for a spin system so far (5 Mio. core hours).
- Unexpected N -dependence of low- T peak of heat capacity.
- B -dependence of density of states leads to asymmetric melting of plateaus.

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Thank you very much for your
attention.

The end.

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