

# Symmetric finite-temperature Lanczos method for anisotropic magnetic molecules

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# Finite-temperature Lanczos method in one slide

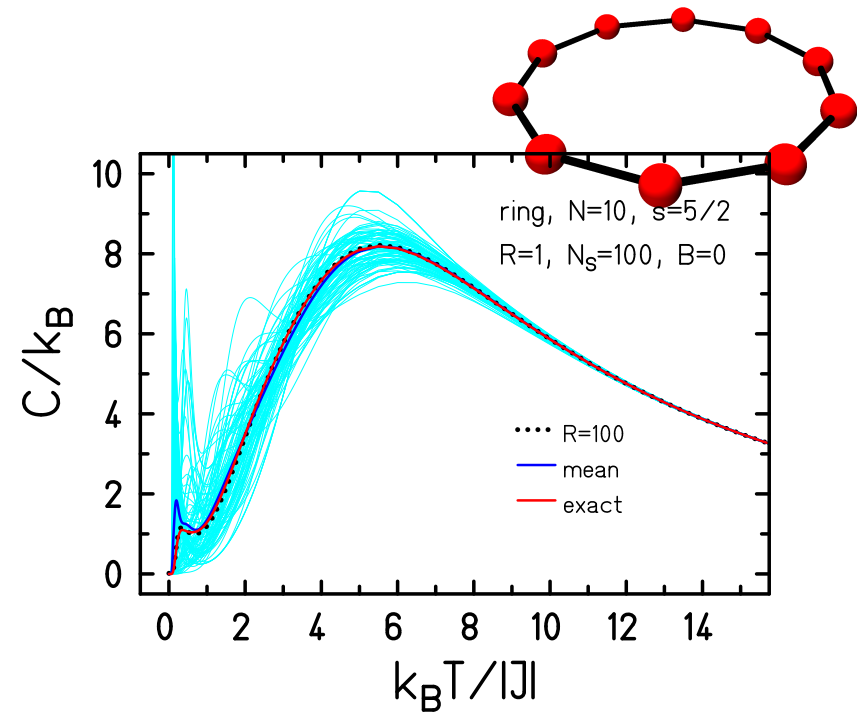
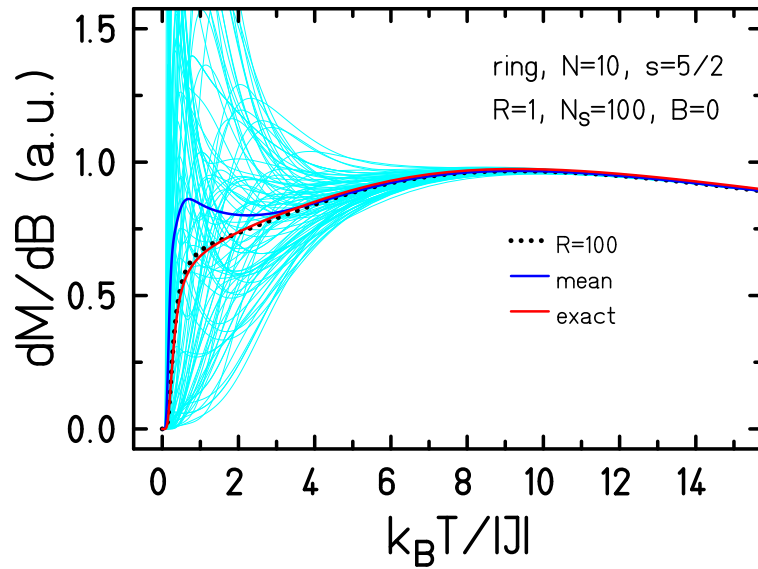
$$Z(T, B) \approx \langle r | e^{-\beta \tilde{H}} | r \rangle \approx \sum_{n=1}^{N_L} e^{-\beta \epsilon_n^{(r)}} |\langle n(r) | r \rangle|^2$$

$$O^r(T, B) \approx \frac{\langle r | Q_{\tilde{H}} e^{-\beta \tilde{H}} | r \rangle}{\langle r | e^{-\beta \tilde{H}} | r \rangle}$$

- $|r\rangle$  is a random vector. Any random vector will do:  $|r\rangle \equiv (T = \infty)$
- $e^{-\beta \tilde{H}} = \sum_{n=1}^{N_L} |n(r)\rangle e^{-\beta \epsilon_n^{(r)}} \langle n(r)|$  is the spectral representation in the Krylov space of dimension  $N_L$  grown from seed  $|r\rangle$ .

(1) J. Jaklic and P. Prelovsek, Phys. Rev. B **49**, 5065 (1994).

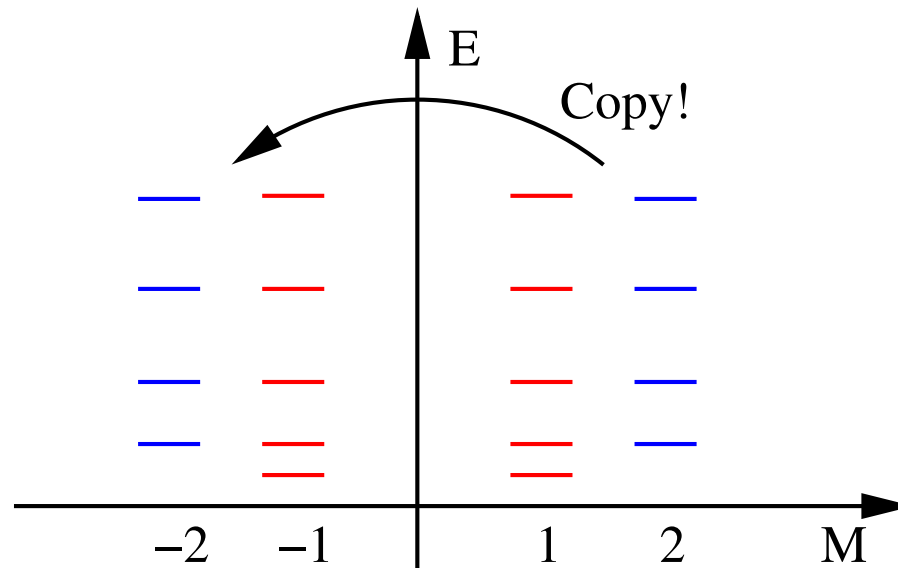
(2) J. Schnack, J. Richter, R. Steinigeweg, Phys. Rev. Research **2**, 013186 (2020).



$$O^{\text{FTLM}}(T, B) \approx \frac{\langle r_1 | Q e^{-\beta \tilde{H}} | r_1 \rangle + \langle r_2 | Q e^{-\beta \tilde{H}} | r_2 \rangle + \dots}{\langle r_1 | e^{-\beta \tilde{H}} | r_1 \rangle + \langle r_2 | e^{-\beta \tilde{H}} | r_2 \rangle + \dots} \quad (\text{correct})$$

$$O^{\text{mean}}(T, B) \approx \frac{1}{R} \left( \frac{\langle r_1 | Q e^{-\beta \tilde{H}} | r_1 \rangle}{\langle r_1 | e^{-\beta \tilde{H}} | r_1 \rangle} + \frac{\langle r_2 | Q e^{-\beta \tilde{H}} | r_2 \rangle}{\langle r_2 | e^{-\beta \tilde{H}} | r_2 \rangle} + \dots \right) \quad (\text{wrong})$$

# Why is FTLM so good?



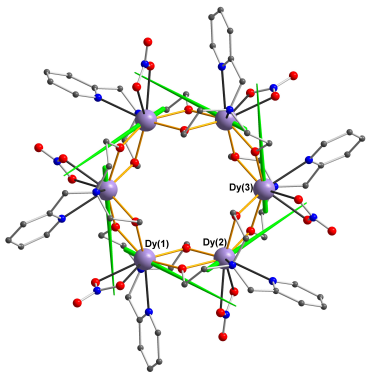
$S^z$ -symmetry is used, spectra are generated in orthogonal subspaces  $\mathcal{H}(M \geq 0)$

spectra for subspaces  $\mathcal{H}(M < 0)$  are taken as copies

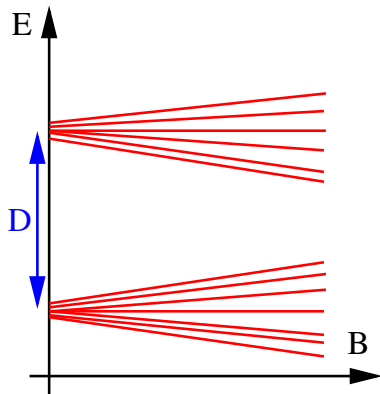
What would happen if you would generate the spectra for subspaces  $\mathcal{H}(M)$  and  $\mathcal{H}(-M)$  in independent random processes?

# Effective model for Dy<sub>6</sub>

$$\tilde{H} = \sum_{k < l} \vec{j}_k \cdot \mathbf{J}_{kl} \cdot \vec{j}_l + \sum_k \vec{j}_k \cdot \mathbf{D}_k \cdot \vec{j}_k + \mu_B \vec{B} \cdot \sum_k g_k \vec{j}_k$$

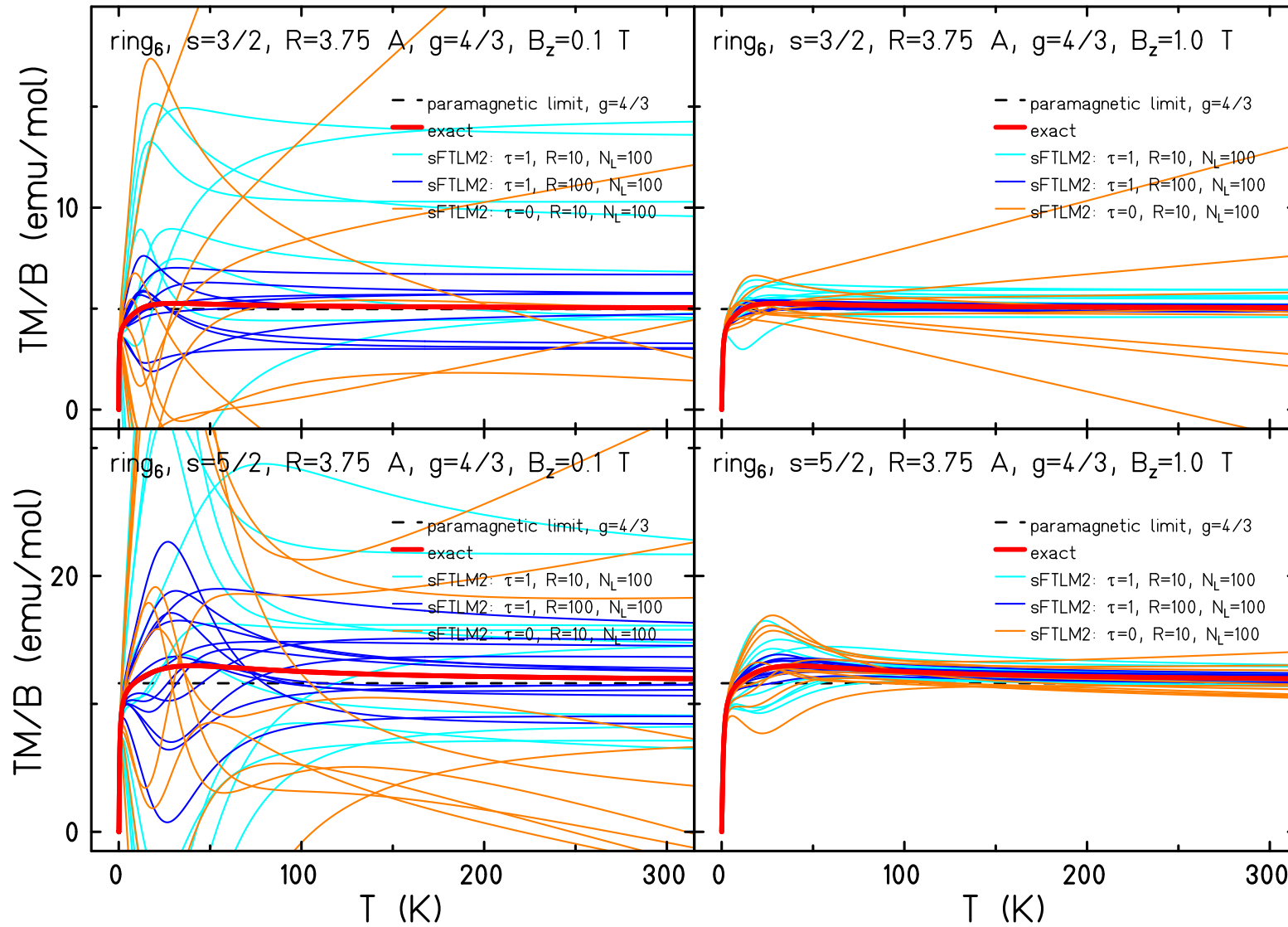


- Very strong alternating easy axes with  $D \approx -20$  K.  $J \approx -0.02$  K and (stronger) dipolar interaction.
- Hamiltonian has no symmetries! Anisotropy is dominant.
- $\dim \mathcal{H} = 16,777,216 \Rightarrow$  FTLM!

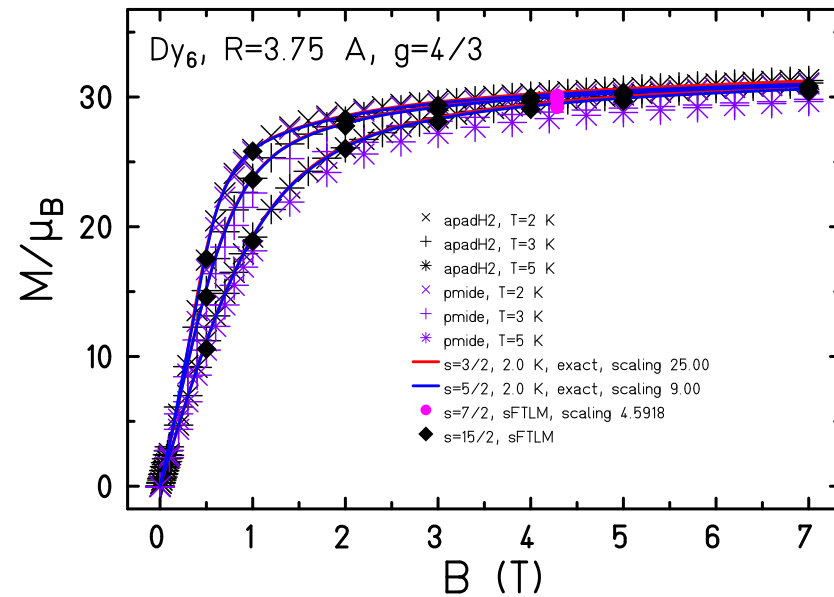
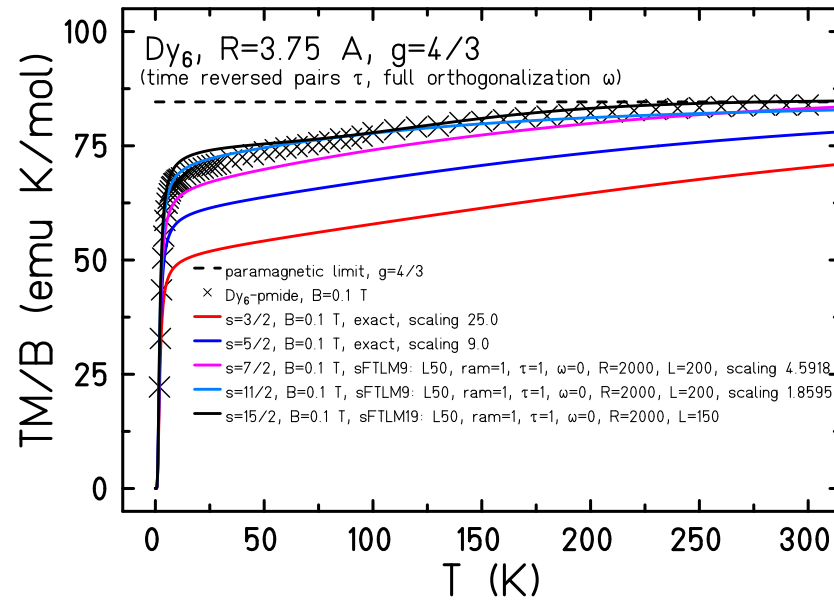


Warning! Method is approximate and holds only for small enough  $B$  since spin and orbital angular momentum have got different  $g_k$ .

# Problem – FTLM converges badly for anisotropic models



# Dy<sub>6</sub> – results



1. Use pairs of time-reversed random vectors (1)

2. Use symmetric version of FTLM (2):

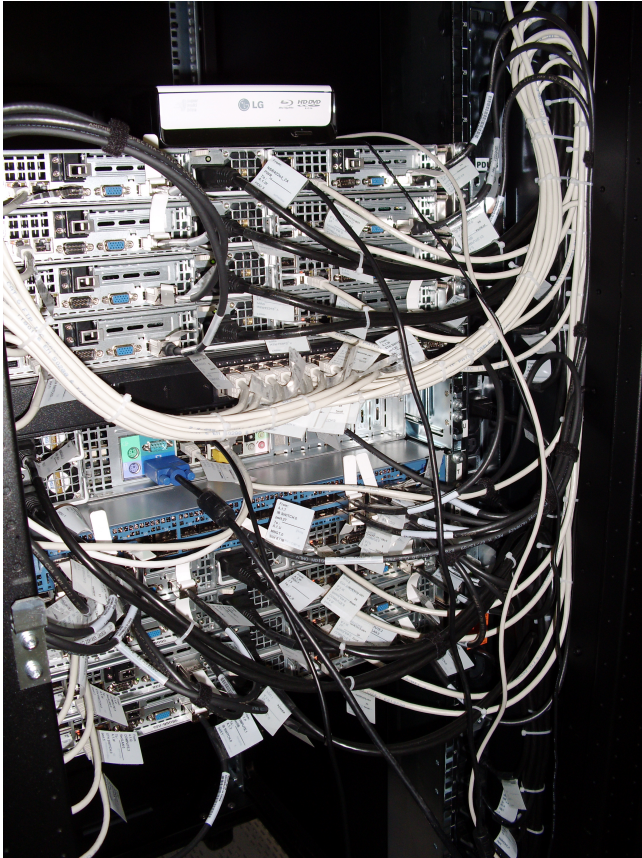
$$\text{Tr} \left( \underline{Q} e^{-\beta \underline{H}} \right) = \text{Tr} \left( e^{-\beta \underline{H}/2} \underline{Q} e^{-\beta \underline{H}/2} \right) \approx \langle r | e^{-\beta \underline{H}/2} \underline{Q} e^{-\beta \underline{H}/2} | r \rangle \neq \langle r | e^{-\beta \underline{H}} \underline{Q} | r \rangle$$

(1) O. Hanebaum, J. Schnack, Eur. Phys. J. B **87**, 194 (2014).

(2) M. Aichhorn, M. Daghofer, H. G. Evertz, and W. von der Linden, Phys. Rev. B **67**, 161103(R) (2003).

(3) D. Westerbeck, Ph.D. thesis, Bielefeld University (2025).

## Summary + To-Do-List



- Spin models are always approximate in case of strong spin-orbit interaction. Actually, they are always approximate (cf. e.g. Oliver Waldmann).
- Multicenter  $j$ -spin models allow accurate modelling and effective understanding of various aspects of  $Dy_6$ .
- How far does this description hold?
- FTLM for anisotropic models has to be further improved to converge with less random vectors.

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