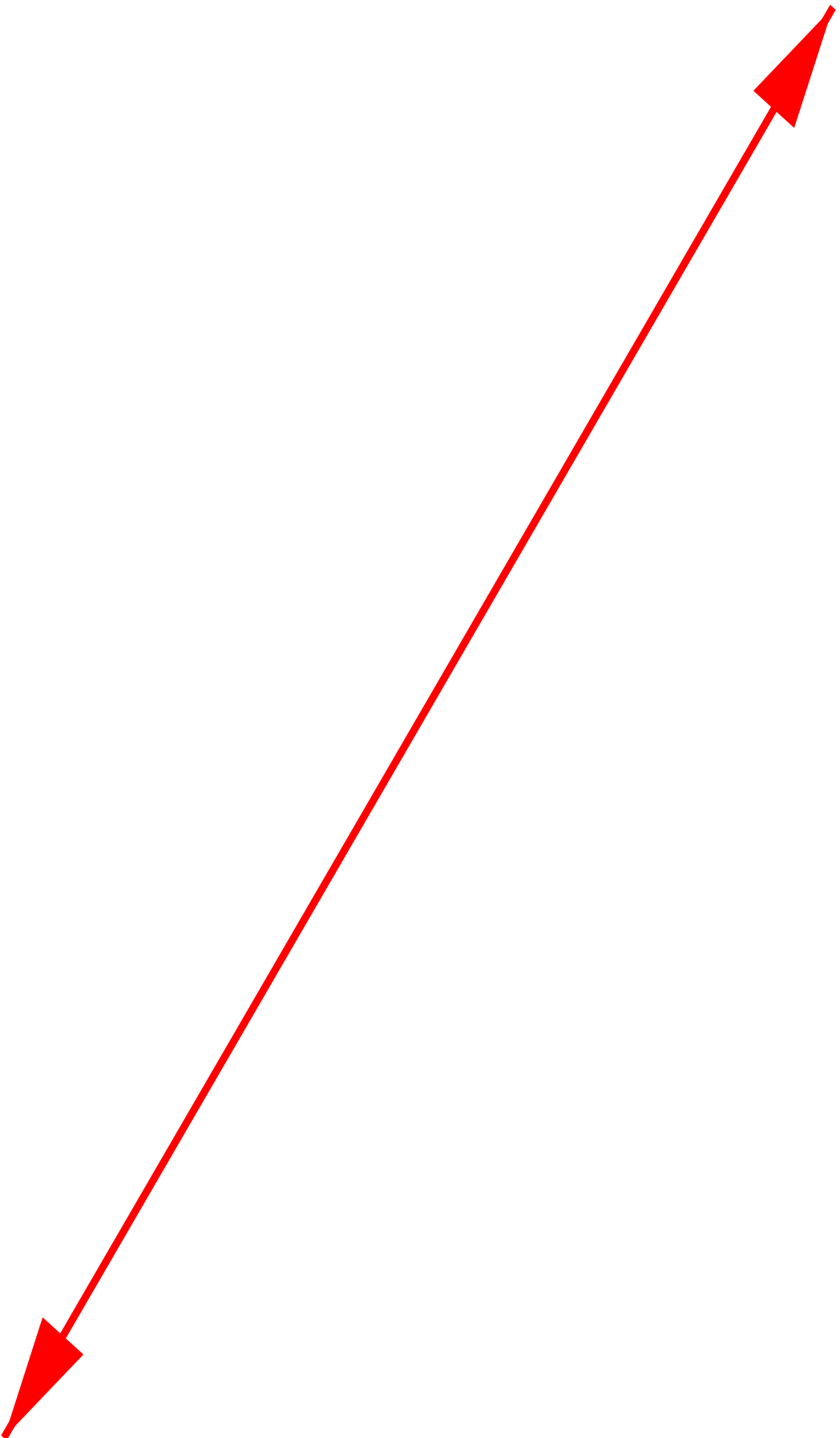


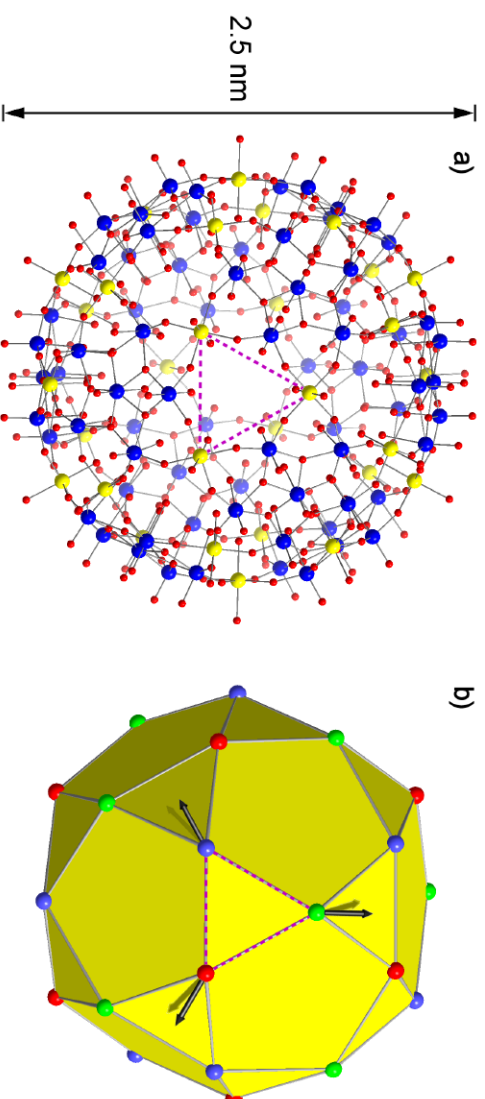
**Size**



# What does the icosidodecahedron teach us about magnetism?

Jürgen Schnack

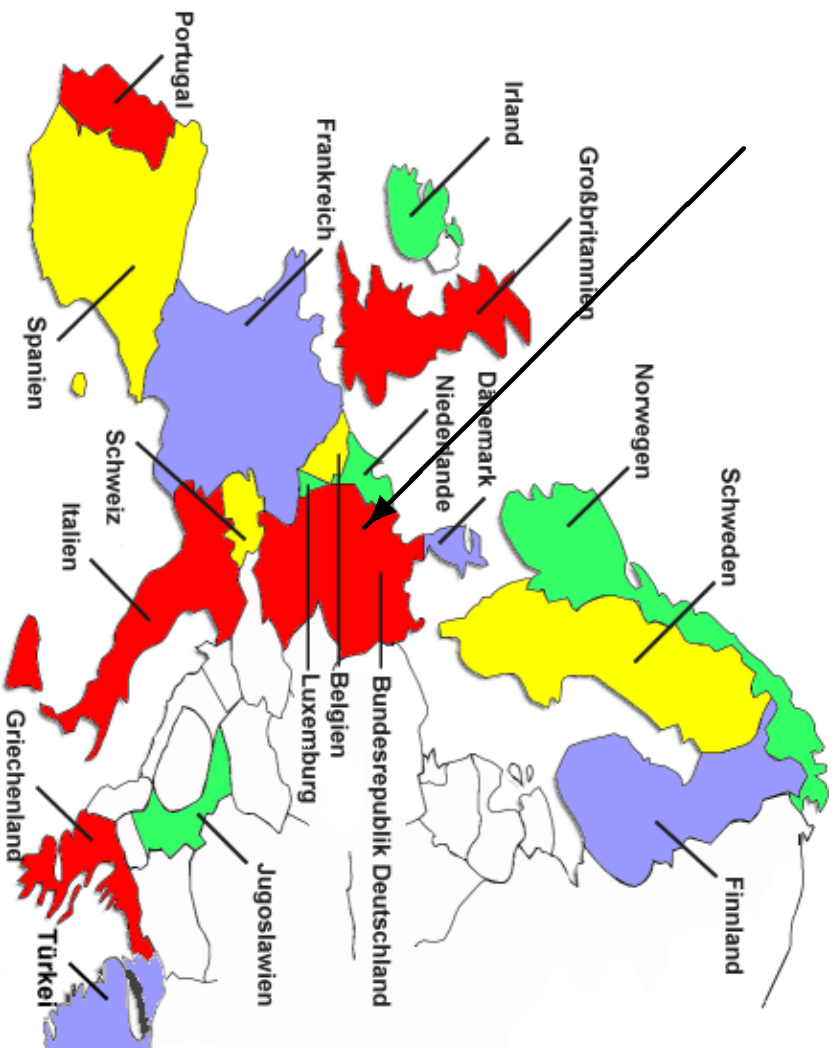
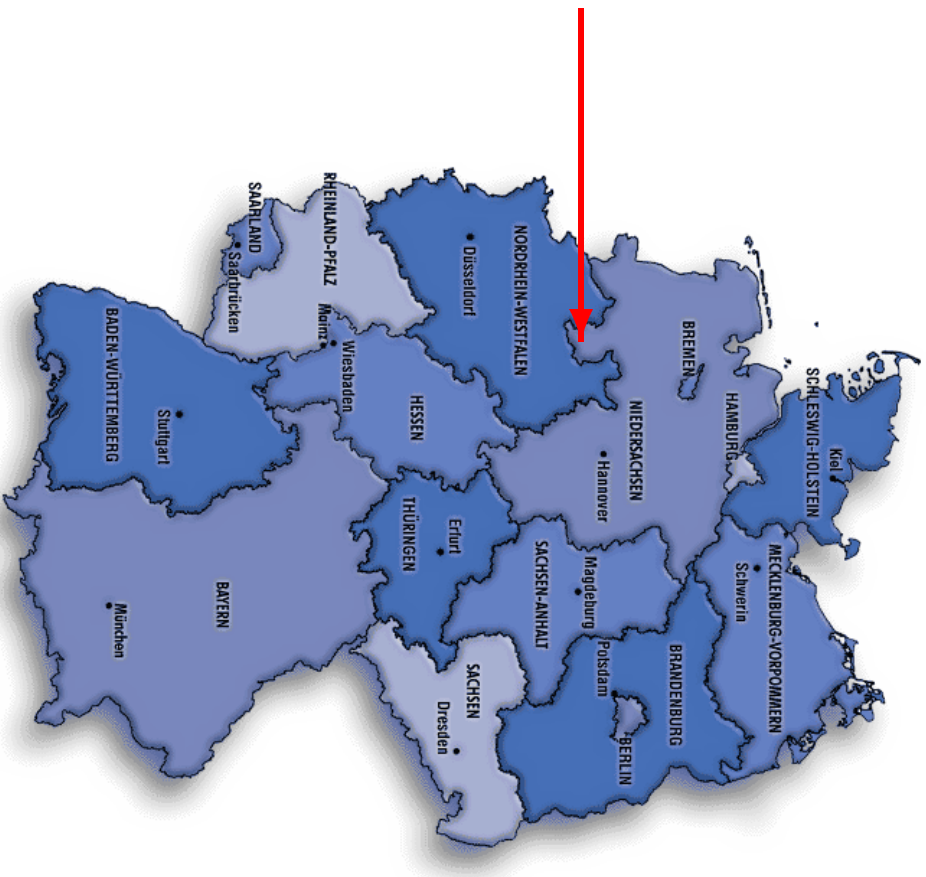
Universität Osnabrück, D-49069 Osnabrück, Germany



thanks to Paul Kögerler

<http://obelix.physik.uni-osnabrueck.de/~schnack/>

# Where is Osnabrück



# Osnabrück, Germany

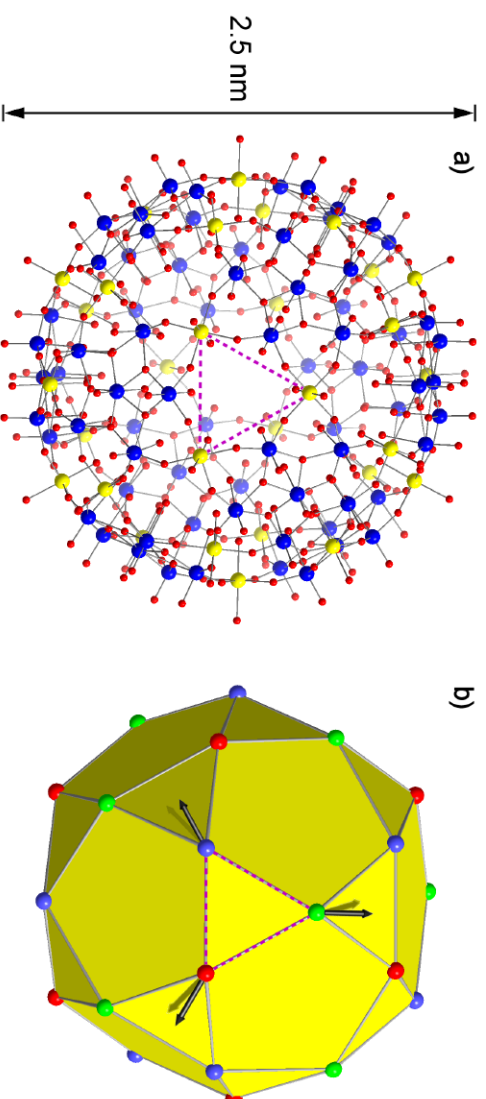
- 9 AD three Roman legions, around 20,000 men, were killed in a battle which lasted for three days. Lulled into a false sense of security by the Germanic chief Arminius, the Roman governor Publius Quinctilius Varus led his army into a trap that only a handful managed to escape alive. The loss of the Varian Legions was a massive psychological blow to the Roman Empire and, after 9 AD, the Romans gave up their plans to hold Germania and withdrew to the west bank of the Rhine.
- Osnabrück was founded by Charles the Great in 780.
- The peace treaty after the 30 years war was signed here and in Münster in 1648.
- 150000 inhabitants
- Mercedes cabrios are assembled in Osnabrück, we don't have a Barilla factory



# What does the icosidodecahedron teach us about magnetism?

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# Contents

1. Heisenberg model
2. Rotational band hypothesis
3. True lowest levels of the  $s = 1/2$  icosidodecahedron
4. Independent magnons
5. Giant magnetisation step on spin lattices
6. Back to the icosidodecahedron:  $\text{Fe}_{30}$
7. There's still a lot to be done.

# Heisenberg model

Hamilton operator (AF:  $J > 0^a$ )

$$\tilde{H} = J \sum_{(u,v)} \vec{s}(u) \cdot \vec{s}(v) + g \mu_B B \sum_u^N s_z(u)$$

Symmetries and good quantum numbers

$$\left[ \tilde{H}, \tilde{S}^2 \right] = 0 \quad \& \quad \left[ \tilde{H}, \tilde{S}_z \right] = 0 \quad \& \quad \left[ \tilde{S}^2, \tilde{S}_z \right] = 0 \quad , \quad \tilde{S}_z = \sum_u^N s_z(u)$$

For this talk I will assume that all spins have the same length  $s$ .

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<sup>a</sup>Don't worry. This definition looks different in every publication, even in our own.

## Typical observables in the canonical ensemble

Mean energy and specific heat

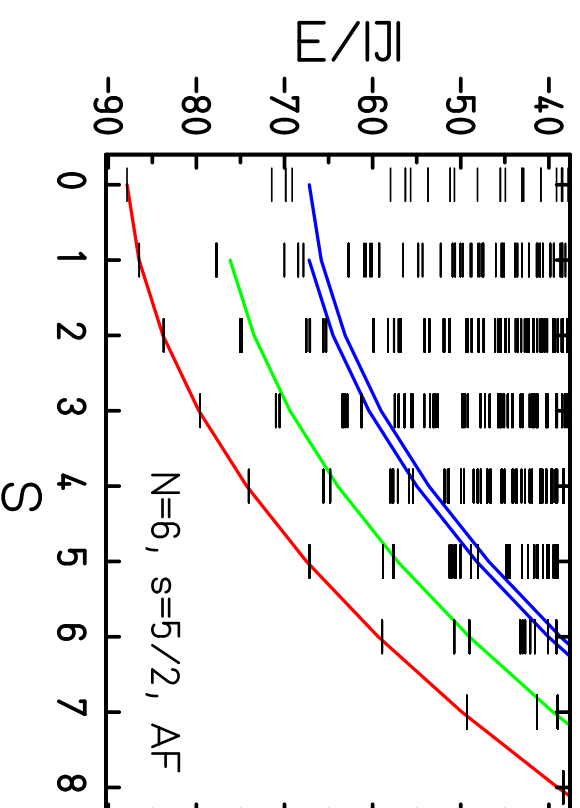
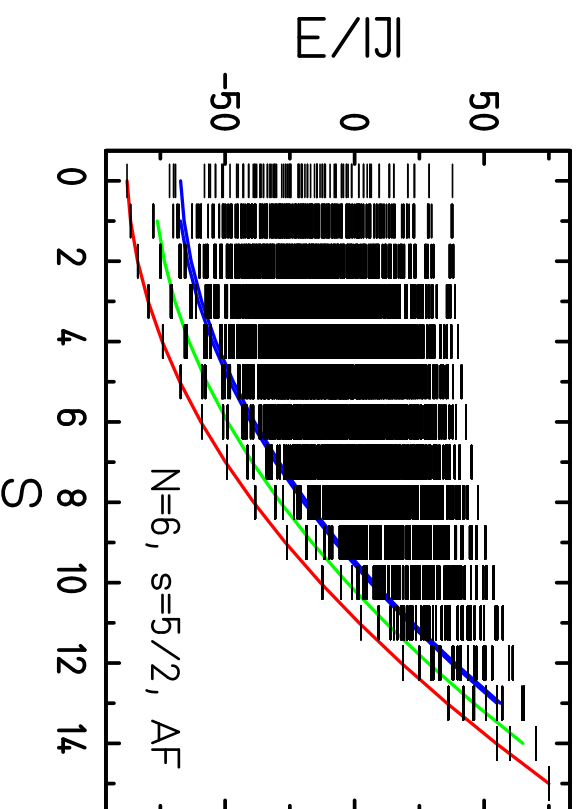
$$\begin{aligned}\langle\langle \tilde{H} \rangle\rangle &= \frac{1}{Z} \operatorname{tr} \left\{ \tilde{H} e^{-\beta \tilde{H}} \right\}, & Z &= \operatorname{tr} \left\{ e^{-\beta \tilde{H}} \right\}, & \beta &= \frac{1}{kT} \\ C &= \frac{d}{dT} \langle\langle \tilde{H} \rangle\rangle = \frac{1}{kT^2} \left( \langle\langle \tilde{H}^2 \rangle\rangle - \langle\langle \tilde{H} \rangle\rangle^2 \right)\end{aligned}$$

Magnetisation and magnetic susceptibility

$$\begin{aligned}\mathcal{M} &= g\mu_B \left( \frac{1}{Z} \operatorname{tr} \left\{ \tilde{S}_z e^{-\beta \tilde{H}} \right\} \right) \\ \chi &= \left( \frac{\partial \mathcal{M}}{\partial B} \right) = g^2 \mu_B^2 \beta \left( \langle\langle (\tilde{S}_z)^2 \rangle\rangle - \langle\langle \tilde{S}_z \rangle\rangle^2 \right)\end{aligned}$$



## Rotational bands in AF Heisenberg rings



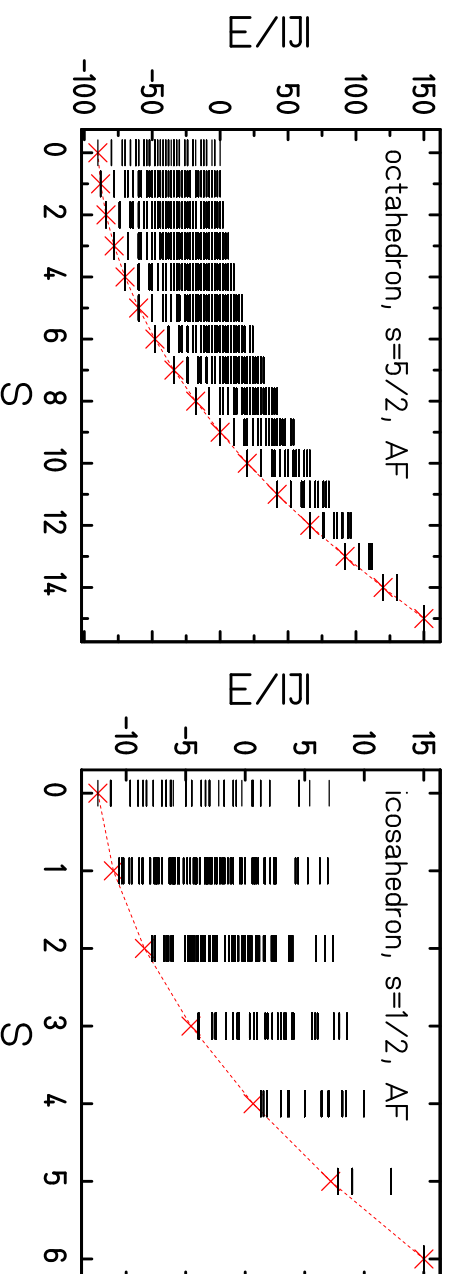
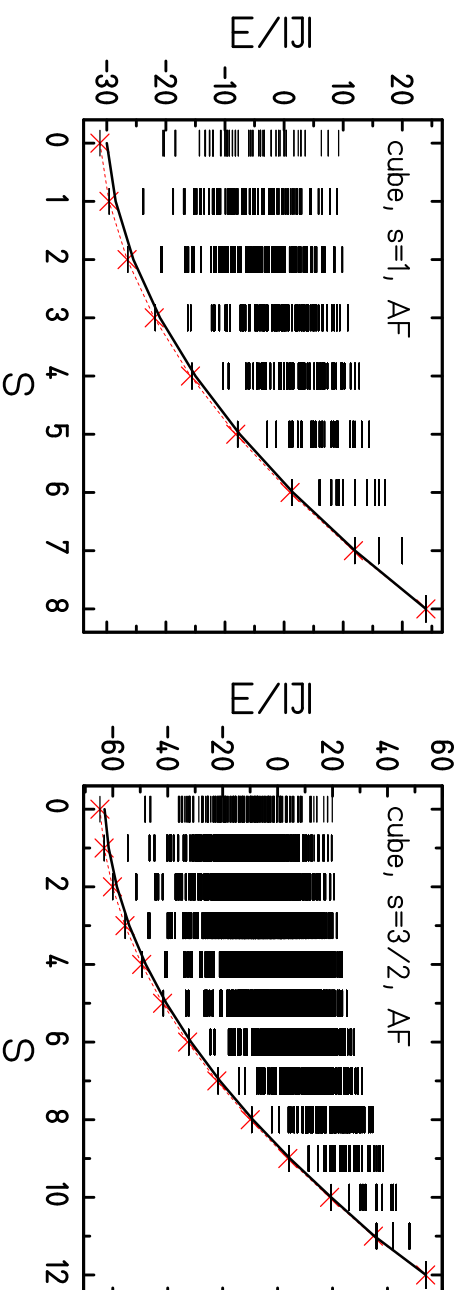
Very often the minimal energies  $E_{min}(S)$  form a rotational band, i. e. depend approximately quadratically on the total spin quantum number  $S$  (Landé interval rule).<sup>a</sup>

Sometimes the low-lying spectrum is dominated by a sequence of rotational bands.

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<sup>a</sup>Chem. Eur. J. **2**, 1379 (1996); G. L. Abbati *et al.*, Inorg. Chim. Acta **297**, 291 (2000); J. Schnack and M. Luban, Phys. Rev. B **63**, 014418 (2001).

# Rotational bands in AF Heisenberg polytopes



J. Schnack and M. Luban, Phys. Rev. B **63** (2001), 014418.

## Rotational band Hamiltonian for Fe<sub>30</sub>

$$\begin{aligned}
 \tilde{H} &= -2J \sum_{(u < v)} \vec{s}(u) \cdot \vec{s}(v) \approx -\frac{DJ}{N} \left[ \tilde{S}^2 - \sum_{j=1}^{N_{SL}} \tilde{S}_j^2 \right] = \tilde{H}_1^{\text{eff}} \\
 &\approx -J \frac{D(N, s)}{N} \left[ \tilde{S}^2 - \gamma(N, s) \left( \sum_{j=1}^{N_{SL}} \tilde{S}_j^2 \right) \right] = \tilde{H}_2^{\text{eff}}
 \end{aligned}$$

$N_{SL}$  – number of sublattices,  $\tilde{S}_j$  – sublattice spin;

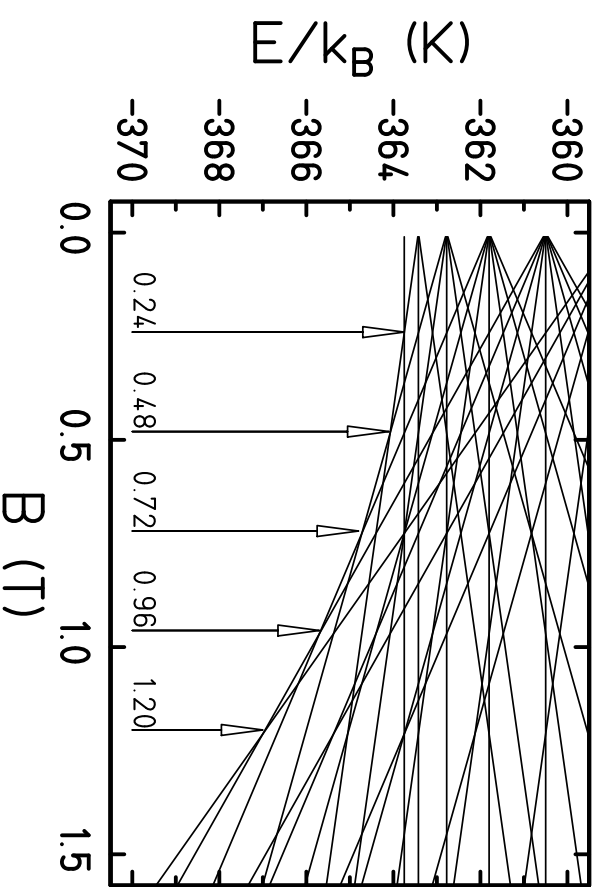
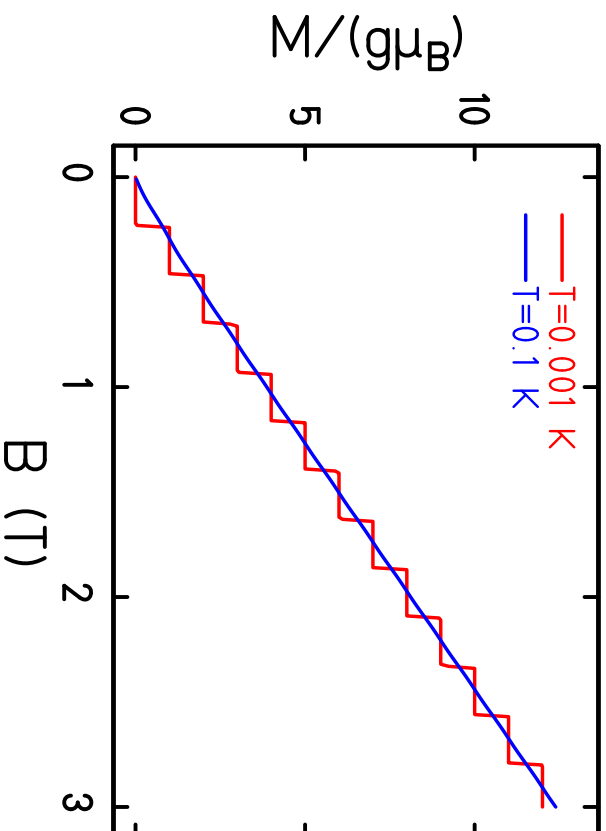
**{Mo<sub>72</sub>Fe<sub>30</sub>}**<sup>a</sup>:

- $N_{SL} = 3$ ,  $S_A, S_B, S_C = 0, 1, \dots, 25$ ,  $S = 0, 1, \dots, 75$ ;
- $D = 6$  determined from corresponding classical system or equivalently from sublattice structure;
- finite size effect:  $D(N, s) = 6.23$ ,  $\gamma(N, s) = 1.07$ .

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<sup>a</sup>J. Schnack, M. Luban, R. Modler, *Quantum rotational band model for the Heisenberg molecular magnet {Mo<sub>72</sub>Fe<sub>30</sub>}*, Europhys. Lett. (2001) submitted.

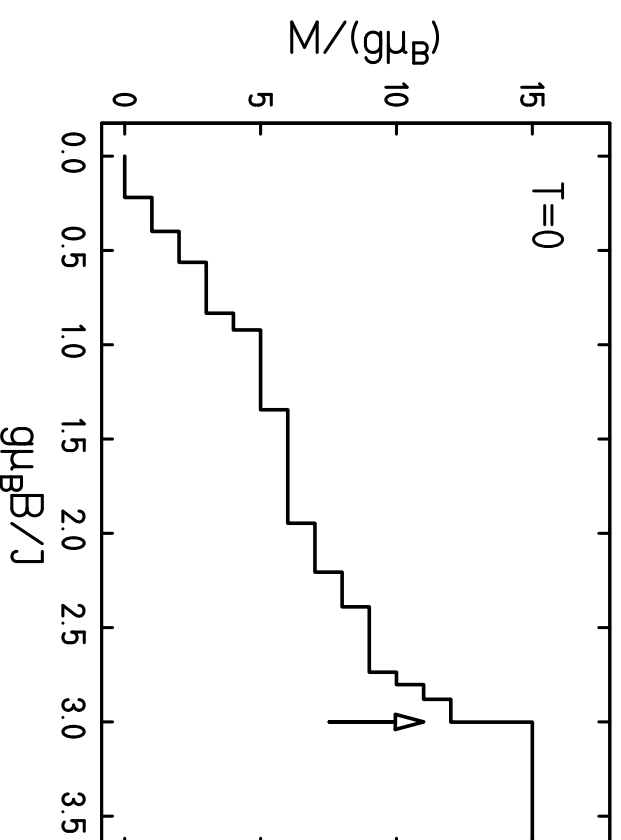
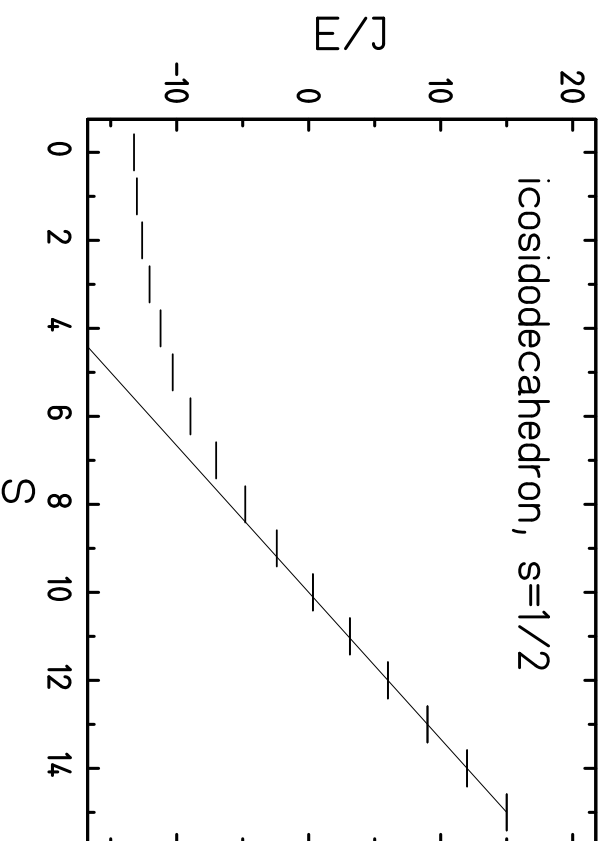
## Can one detect rotational bands?



Low-temperature magnetisation measurements as well as low-temperature NMR can detect the ground state band. Higher bands might be accessible by neutron scattering or ESR.<sup>a</sup>

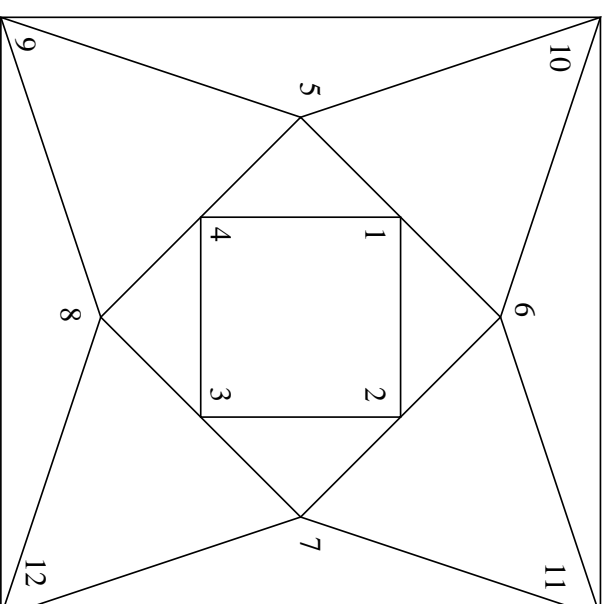
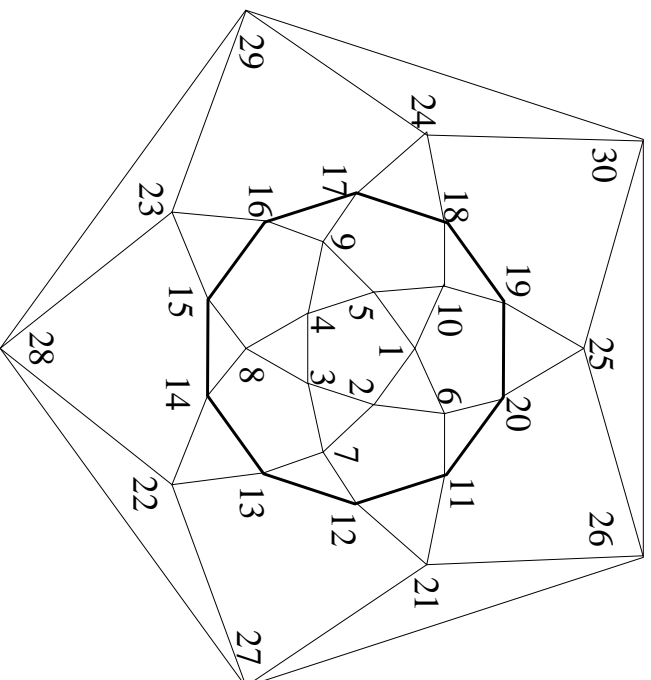
<sup>a</sup>J. Schnack, M. Luban, R. Modler, *Quantum rotational band model for the Heisenberg molecular magnet*  $\{M_{072}Fe_{30}\}$ , Europhys. Lett. (2001) submitted.

# Icosidodecahedron with $s = 1/2$

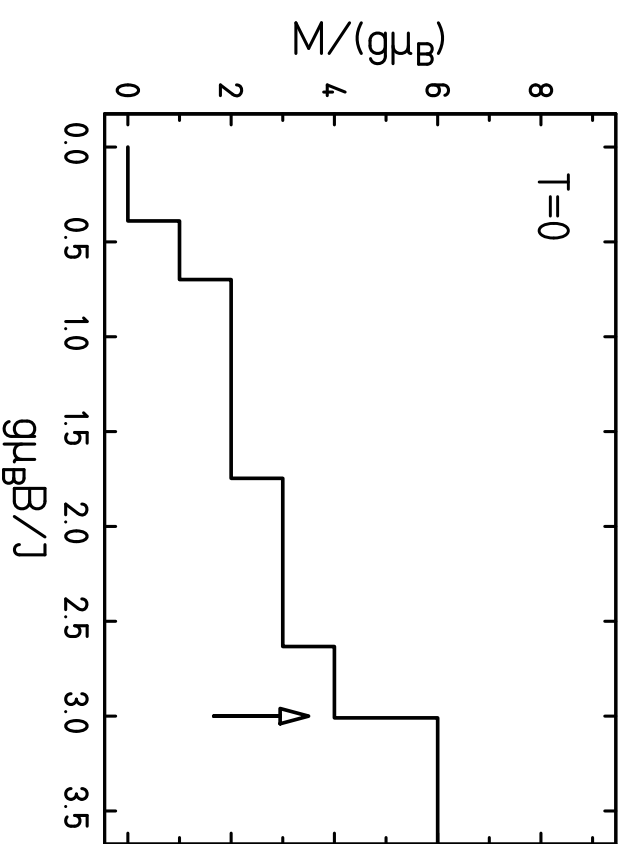
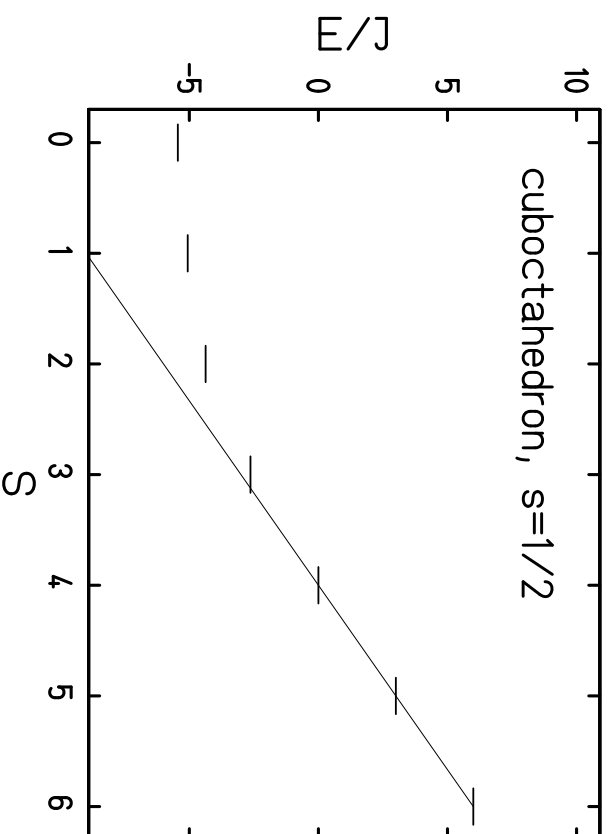


J. Schnack, H.-J. Schmidt, J. Richter, J. Schulenburg, *Independent magnon states on magnetic polytopes*, Eur. Phys. J. B (2001) submitted; cond-mat/0108432

# Structure of Icosidodecahedron and Cuboctahedron



# Cuboctahedron with $s = 1/2$



J. Schnack, H.-J. Schmidt, J. Richter, J. Schulenburg, *Independent magnon states on magnetic polytopes*, Eur. Phys. J. B (2001) submitted; cond-mat/0108432

## Questions

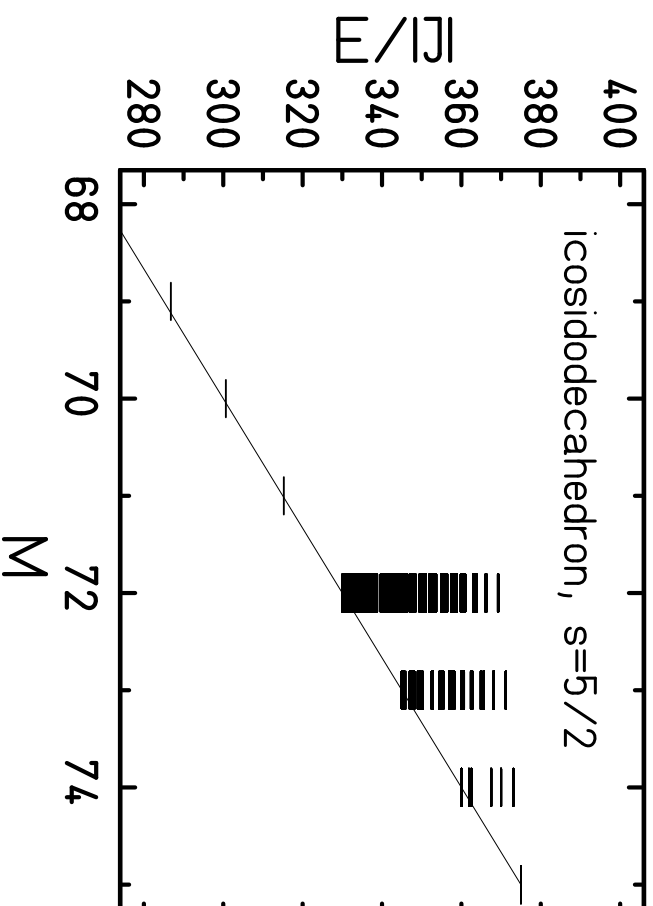
- Magnetisation curve shows plateaus and jumps!
- Where are the rotational bands?
- Accident ?
- How can one understand the jumps, i.e. the linear behaviour in  $E_{\min}(S)$ ?
- Are there other structures with such properties?
- Who saves our brilliant ideas<sup>a</sup> about  $\text{Fe}_{30}$ ?

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<sup>a</sup>J. Schnack, M. Luban, R. Modler, *Quantum rotational band model for the Heisenberg molecular magnet*  $\{M_{072}\text{Fe}_{30}\}$ , Europhys. Lett. (2001) submitted.

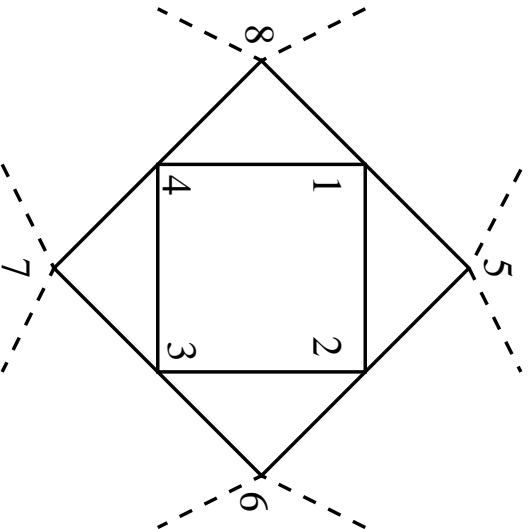


# Independent Magnons



- magnon vacuum:  $M = 75$ , one-magnon space:  $M = 74$
- one-magnon ground state highly degenerate
- localized one-magnon states can be constructed
- if spin array large enough several localized one-magnon states can be placed on the grid without interaction
- minimal energy, i.e. ground state energy in the  $n$ -magnon spaces, is linear in  $M$

## Localized Magnon – Example



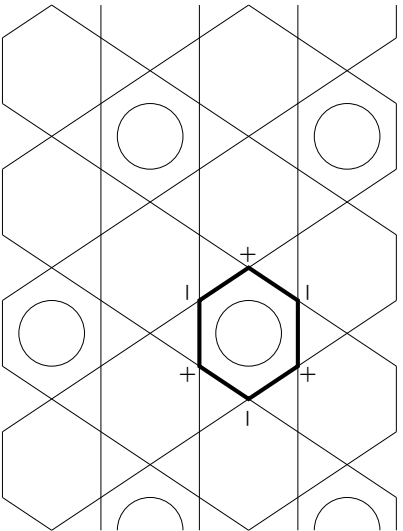
- $|\text{localized magnon}\rangle = \frac{1}{2} (|1\rangle - |2\rangle + |3\rangle - |4\rangle)$
- $|u\rangle = \tilde{s}^-(u) |\Omega\rangle$ ;  $|\Omega\rangle$  – magnon vacuum
- $\tilde{H}|1\rangle = J\{|1\rangle + 1/2(|2\rangle + |4\rangle + |5\rangle + |8\rangle)\}$
- $\tilde{H}|\text{localized magnon}\rangle \propto |\text{localized magnon}\rangle$
- triangles trap the localized magnon, amplitudes cancel at outer vertices
- proven for  $s = 1/2$ , one exchange constant  $J$ , and same number of interactions for each site<sup>a</sup>
- result also holds for  $s > 1/2$  and XXZ model ( $\Delta \geq 0$ ).

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<sup>a</sup>J. Schnack, H.-J. Schmidt, J. Richter, J. Schulenburg, *Independent magnon states on magnetic polytopes*, Eur. Phys. J. B (2001) submitted; cond-mat/0108432

# Kagomé Lattice – Independent Magnons

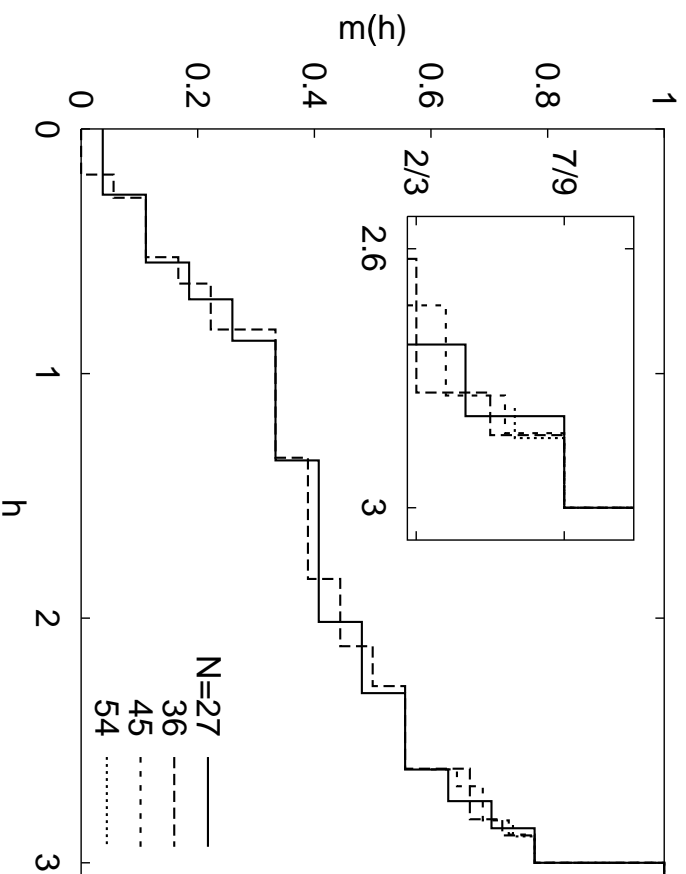
- localized one-magnon state indicated by bold lines;
- independent (non-interacting) one-magnon states can be placed on the grid (circles);
- due to the absence of attractive interaction, each state of  $n$  independent magnons is the ground state in the Hilbert subspace with  $M = Ns - n$ ;
- $\Rightarrow$  linear dependence of  $E_{\min}$  on  $M$ ;
- $\Rightarrow$  magnetisation jump;
- maximal number of independent magnons:  $N/9$ ;
- magnetisation jump is a macroscopic quantum effect!



J. Schulenburg, A. Honecker, J. Schnack, J. Richter, H.-J. Schmidt, *Macroscopic magnetization jumps due to independent magnons in frustrated quantum spin lattices*, Phys. Rev. Lett. (2001) submitted; cond-mat/0108498

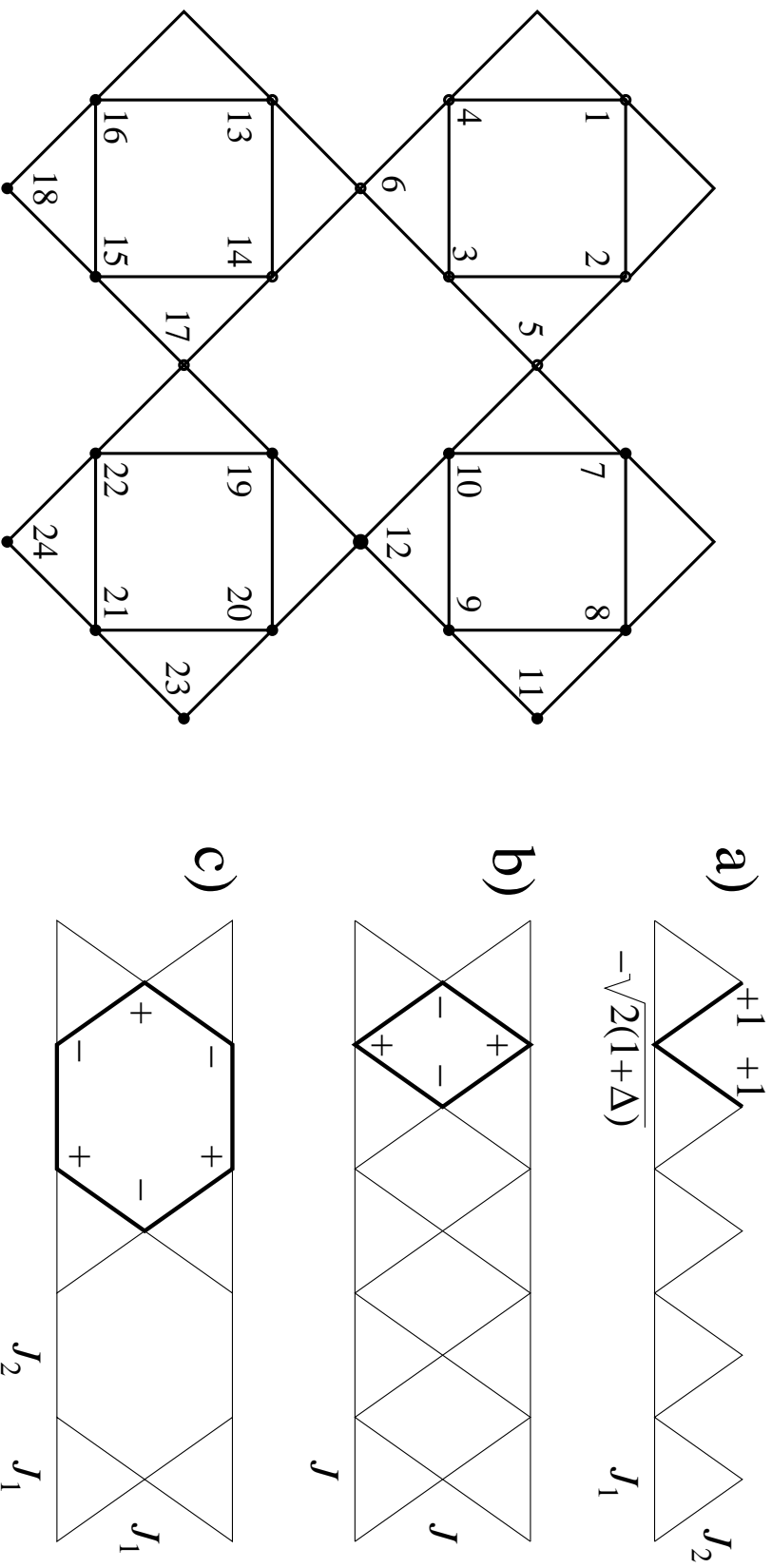
# Kagomé Lattice

## Giant Magnetisation Jump



J. Schulenburg, A. Honecker, J. Schnack, J. Richter, H.-J. Schmidt, *Macroscopic magnetization jumps due to independent magnons in frustrated quantum spin lattices*, Phys. Rev. Lett. (2001) submitted; cond-mat/0108498

# Structures With Magnetisation Jump



**Who saves our brilliant ideas  
about  $\text{Fe}_{30}$  ?**

**Marshall Luban and Heinz-Jürgen Schmidt  
&  
Matthias Exler**

There's still a lot to be done.

List of wishes for Christmas:

- Paul Kögerler – please synthesise an icosidodecahedron with  $s = 1/2$ .  
What about  $\text{Cu}_{30}$  or  $\text{V}_{30}$ ?
- Marshall Luban and Heinz-Jürgen Schmidt – please publish your proof.
- Robert Modler – Can't we go to 50 mK with  $\text{Fe}_{30}$ ?
- understand specific heat of  $\text{Fe}_{30} \Rightarrow$  need better low-energy spectrum;
- Neutron scattering with  $\text{Fe}_{30}$ ?
- find material with giant magnetisation jump

Meanwhile, what I have to do





## Acknowledgement

- Marshall Luban, Robert Modler, Paul Kögerler – Ames Lab
- Klaus Bärwinkel, Heinz-Jürgen Schmidt, Detlef Mentrup, Matthias Exler – University of Osnabrück
- Johannes Richter, Jörg Schulenburg – University of Magdeburg
- Andreas Honecker – University of Braunschweig
- National Science Foundation and the Deutscher Akademischer Austauschdienst for supporting a mutual exchange program