

Frustrated about Frustration

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An attempt will be made to sharpen the term **frustration** which appears very often in connection with antiferromagnetically coupled spin systems. Usage of this term is very sloppy, observable properties, which hold in general, obscure.

Typical statement: "It is known that frustrated spin systems exhibit spectacular phenomena: high ground state degeneracy, re-entrance, partial disorder, controversial nature of the phase transition, order by the disorder, etc."

<http://obelix.physik.uni-osnabrueck.de/~schnack/>

Hamilton Operator

Hamilton operator; AF $J < 0$, F: $J > 0$

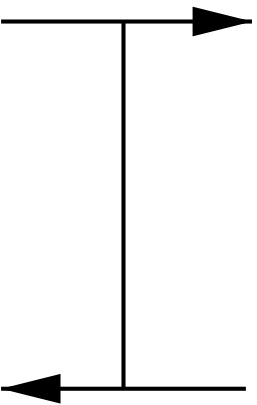
$$\begin{aligned}\tilde{H}(\gamma) &= -2J \sum_x \left\{ \tilde{s}^3(x) \tilde{s}^3(x+1) + \gamma \left[\tilde{s}^1(x) \tilde{s}^1(x+1) + \tilde{s}^2(x) \tilde{s}^2(x+1) \right] \right\} \\ &= -2J \sum_x \left\{ \tilde{s}^3(x) \tilde{s}^3(x+1) + \frac{\gamma}{2} \left[\tilde{s}^+(x) \tilde{s}^-(x+1) + \tilde{s}^-(x) \tilde{s}^+(x+1) \right] \right\} \\ \tilde{H}(1) &= -2J \sum_x \tilde{\vec{s}}(x) \cdot \tilde{\vec{s}}(x+1)\end{aligned}$$

Classical Hamilton function

$$\begin{aligned}H(\gamma) &= -2J_c \sum_x \left\{ e^3(x) e^3(x+1) + \gamma \left[e^1(x) e^1(x+1) + e^2(x) e^2(x+1) \right] \right\} \\ H(1) &= -2J_c \sum_x \vec{e}(x) \cdot \vec{e}(x+1)\end{aligned}$$

Classical Heisenberg Spin Systems I

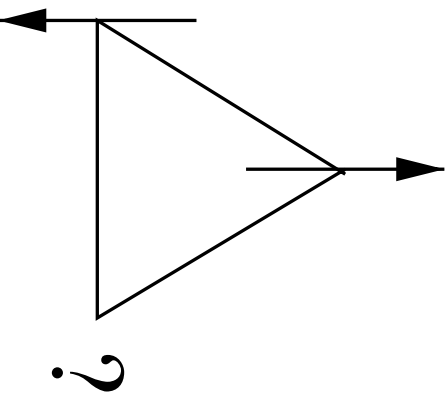
Dimer & trimer



$$E_{GS} = -2J(\vec{e}(1) \cdot \vec{e}(2) + \vec{e}(2) \cdot \vec{e}(1)) = 4J$$

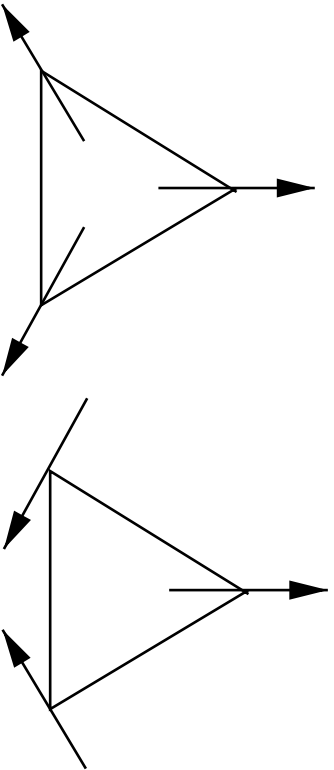
Very simple definition
of frustration:

The last spin is frustrated because it
does not know how to align.



Classical Heisenberg Spin Systems II

Non-trivial degeneracy

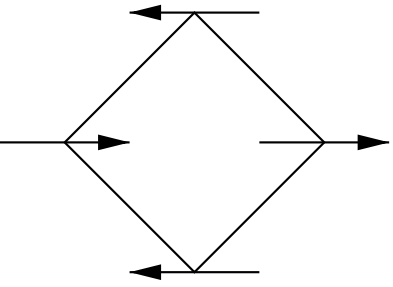


Definition I: **The ground state possesses a non-trivial degeneracy.**

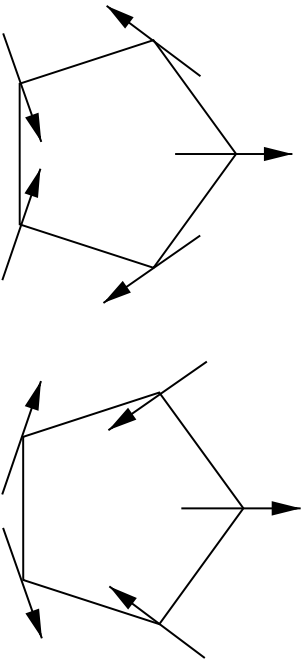
Definition II: $\langle\langle \vec{e}(x) \cdot \vec{e}(x+1) \rangle\rangle (T=0) > \langle\langle \vec{e}(1) \cdot \vec{e}(2) \rangle\rangle_D (T=0) = -1$,
i.e. anti-correlation weaker than in the dimer.

Classical Heisenberg Spin Systems III

Square & pentagon



only trivial degeneracy and $\langle\langle \vec{e}(x) \cdot \vec{e}(x+1) \rangle\rangle(T=0) = -1$



non-trivial degeneracy and
 $\langle\langle \vec{e}(x) \cdot \vec{e}(x+1) \rangle\rangle(T=0) = \cos\left(\frac{4\pi}{5}\right) \approx -0.8 > -1$

Definitions work!?

Ising Model

Hamilton operator of the q -states Potts model; $q = 2 \Rightarrow$ Ising

$$\tilde{H} = \tilde{H}(\gamma = 0) = -2J \sum_x \tilde{s}^3(x) \tilde{s}^3(x+1)$$

$$N = 2, s = \frac{1}{2}$$

ground state energy $E_{GS} = J$ twofold degenerate;

ground states: $|\frac{1}{2}, -\frac{1}{2}\rangle, |-\frac{1}{2}, \frac{1}{2}\rangle$

correlation: $\langle\langle \tilde{s}^3(x) \cdot \tilde{s}^3(x+1) \rangle\rangle (T=0)/s^2 = -1$

$$N = 3, s = \frac{1}{2}$$

ground state energy $E_{GS} = J/2$ 2 · N -fold degenerate;

ground states: $|\frac{1}{2}, \frac{1}{2}, -\frac{1}{2}\rangle, |\frac{1}{2}, -\frac{1}{2}, -\frac{1}{2}\rangle$ and cyclic shifts thereof

correlation: $\langle\langle \tilde{s}^3(x) \cdot \tilde{s}^3(x+1) \rangle\rangle (T=0)/s^2 = -1/3$

$$N = 4, s = \frac{1}{2}$$

ground state energy $E_{GS} = 2J$ twofold degenerate;

ground states: $|\frac{1}{2}, -\frac{1}{2}, \frac{1}{2}, -\frac{1}{2}\rangle, |-\frac{1}{2}, \frac{1}{2}, -\frac{1}{2}, \frac{1}{2}\rangle$

correlation: $\langle\langle \tilde{s}^3(x) \cdot \tilde{s}^3(x+1) \rangle\rangle (T=0)/s^2 = -1$

Definition I (non-trivial degeneracy) applicable! Definition II (correlation) applicable!

Heisenberg Model I

Hamilton operator

$$\tilde{H} = \tilde{H}(\gamma = 1) = -2J \sum_x \tilde{\mathfrak{s}}(x) \cdot \tilde{\mathfrak{s}}(x+1)$$

$$N = 2, s = \frac{1}{2}$$

ground state energy $E_{GS} = 3J$ non-degenerate;

ground state: $|GS\rangle = \frac{1}{\sqrt{2}} (|\frac{1}{2}, -\frac{1}{2}\rangle - |-\frac{1}{2}, \frac{1}{2}\rangle)$

correlation: $\langle\langle \tilde{\mathfrak{s}}(x) \cdot \tilde{\mathfrak{s}}(x+1) \rangle\rangle (T=0) / \sqrt{s(s+1)}^2 = -1$

$$N = 3, s = \frac{1}{2}$$

ground state energy $E_{GS} = 3J/2$ fourfold degenerate;

ground states: $|GS1\rangle = \frac{1}{\sqrt{3}} (|\frac{1}{2}, \frac{1}{2}, -\frac{1}{2}\rangle + e^{i\frac{2\pi}{3}} |-\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\rangle + e^{i\frac{4\pi}{3}} |\frac{1}{2}, -\frac{1}{2}, \frac{1}{2}\rangle)$

$|GS2\rangle = \frac{1}{\sqrt{3}} (|\frac{1}{2}, \frac{1}{2}, -\frac{1}{2}\rangle + e^{-i\frac{2\pi}{3}} |-\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\rangle + e^{-i\frac{4\pi}{3}} |\frac{1}{2}, -\frac{1}{2}, \frac{1}{2}\rangle)$

etc.

correlation: $\langle\langle \tilde{\mathfrak{s}}(x) \cdot \tilde{\mathfrak{s}}(x+1) \rangle\rangle (T=0) / \sqrt{s(s+1)}^2 = -\frac{1}{4}$

$$N = 4, s = \frac{1}{2}$$

ground state energy $E_{GS} = 4J$ non-degenerate;

ground states: $|GS\rangle = \frac{1}{\sqrt{3}} (|\frac{1}{2}, -\frac{1}{2}, \frac{1}{2}, -\frac{1}{2}\rangle + |-\frac{1}{2}, \frac{1}{2}, -\frac{1}{2}, \frac{1}{2}\rangle) +$

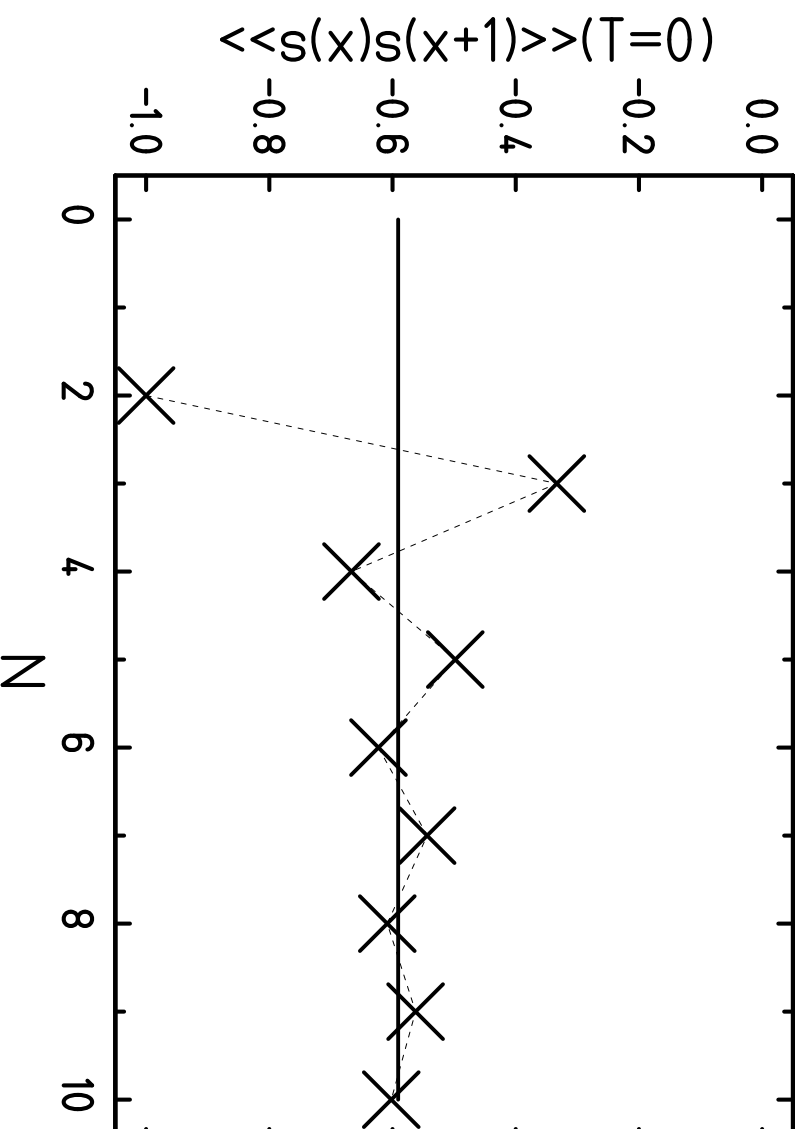
$\frac{1}{\sqrt{12}} (|\frac{1}{2}, \frac{1}{2}, -\frac{1}{2}, -\frac{1}{2}\rangle + |-\frac{1}{2}, \frac{1}{2}, \frac{1}{2}, -\frac{1}{2}\rangle) + |-\frac{1}{2}, -\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\rangle + |\frac{1}{2}, -\frac{1}{2}, -\frac{1}{2}, \frac{1}{2}\rangle)$

correlation: $\langle\langle \tilde{\mathfrak{s}}(x) \cdot \tilde{\mathfrak{s}}(x+1) \rangle\rangle (T=0) / \sqrt{s(s+1)}^2 = -\frac{2}{3}$

Definition I (non-trivial degeneracy) applicable! Definition II (correlation) applicable?

Heisenberg Model II

Next-neighbour correlation $s = 1/2$



The next-neighbour spin-spin correlation does not serve as a measure of frustration.

Heisenberg Model III

Hamilton operator

$$\tilde{H} = \tilde{H}(\gamma = 1) = -2J \sum_x \tilde{\mathfrak{s}}(x) \cdot \tilde{\mathfrak{s}}(x+1)$$

$N = 2, s = 1$

ground state energy $E_{GS} = 8J$ non-degenerate;

ground state: $|GS\rangle = \frac{1}{\sqrt{3}}(|1, -1\rangle + |-1, 1\rangle - |0, 0\rangle)$

correlation: $\langle\langle \tilde{\mathfrak{s}}(x) \cdot \tilde{\mathfrak{s}}(x+1) \rangle\rangle(T=0)/\sqrt{s(s+1)}^2 = -1$

$N = 3, s = 1$

ground state energy $E_{GS} = 6J$ non-degenerate;

ground states: $|GS\rangle = \frac{1}{\sqrt{6}}(|1, 0, -1\rangle + |-1, 1, 0\rangle + |0, -1, 1\rangle +$

$|-1, 0, 1\rangle + |1, -1, 0\rangle + |0, 1, -1\rangle)$

correlation: $\langle\langle \tilde{\mathfrak{s}}(x) \cdot \tilde{\mathfrak{s}}(x+1) \rangle\rangle(T=0)/\sqrt{s(s+1)}^2 = -\frac{1}{2}$

$N = 4, s = 1$

ground state energy $E_{GS} = 12J$ non-degenerate;

correlation: $\langle\langle \tilde{\mathfrak{s}}(x) \cdot \tilde{\mathfrak{s}}(x+1) \rangle\rangle(T=0)/\sqrt{s(s+1)}^2 = -\frac{3}{4}$

Definition I (non-trivial degeneracy) applicable? Definition II (correlation) applicable?

Way Out ?

Definition III

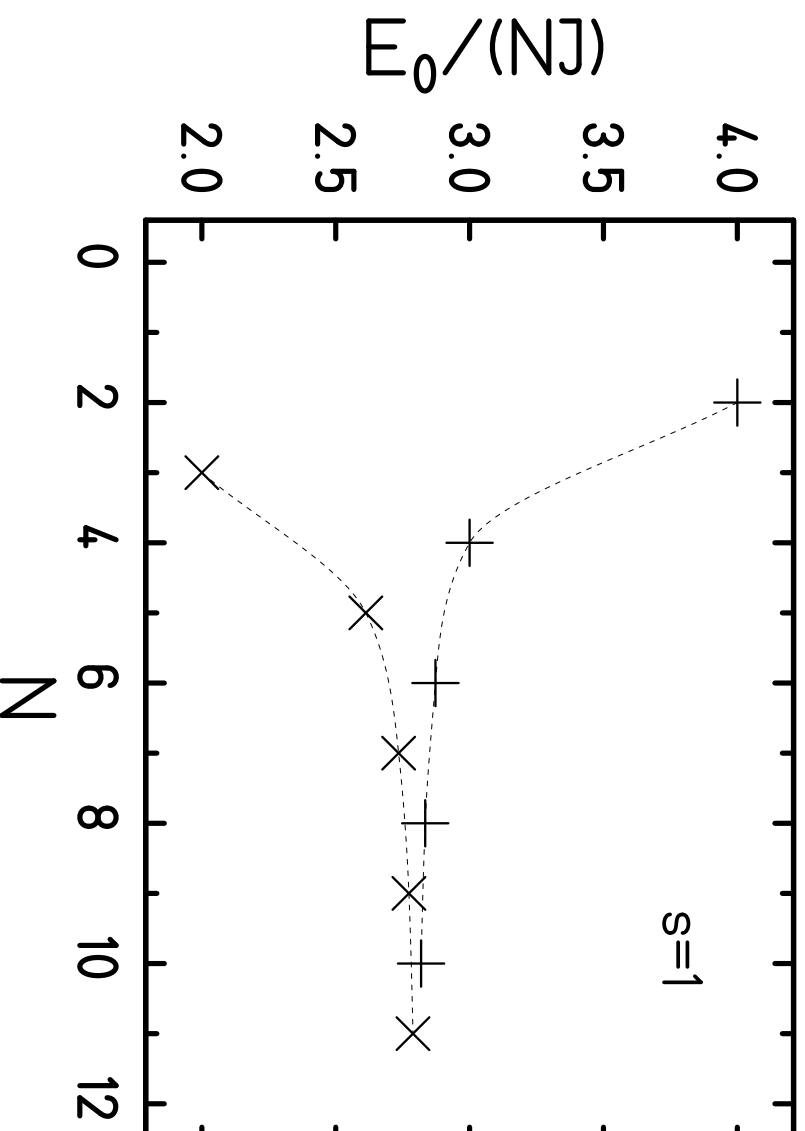
A quantum spin system is frustrated if the corresponding classical system is frustrated.

Problem

- Not very satisfactory. Reminds of the problems one has with “quantum chaos”. There one says that the spectrum should have some special properties like being a GOE.
- Are there any common observable phenomena caused by frustration?
- Or does the pessimistic statement hold, that the classical term frustration is without value in quantum mechanics except for Ising systems?
- One suggestion: weaker binding per bond compared to “neighbouring”, non-frustrated systems

Binding Energy per Bond

Heisenberg ring $s = 1$



Heisenberg ring with e.g. $N = 5$ has a weaker binding per bond than the neighbouring rings with $N = 4$ and $N = 6$. Property holds for arbitrary spin quantum number.

Measure of Frustration

O.k., we still don't know what frustration exactly means, but may be we can quantify it?^a

Satisfied and unsatisfied bonds

satisfied $\left(\langle\langle \vec{s}(x) \cdot \vec{s}(y) \rangle\rangle (T=0) \right) = +\text{sign} (J(x, y)) , J(x, y) \neq 0$

unsatisfied $\left(\langle\langle \vec{s}(x) \cdot \vec{s}(y) \rangle\rangle (T=0) \right) = -\text{sign} (J(x, y)) , J(x, y) \neq 0$

Misfit parameter

$$m = 2 \frac{E_u}{E_u + |E_s|} \qquad E_u = - \sum_{\text{unsatisfied bonds}} J(x, y) \langle\langle \vec{s}(x) \cdot \vec{s}(y) \rangle\rangle (T=0)$$
$$E_s = - \sum_{\text{satisfied bonds}} J(x, y) \langle\langle \vec{s}(x) \cdot \vec{s}(y) \rangle\rangle (T=0)$$

\Rightarrow Heisenberg spin rings always have $m = 0$! Misfit parameter works for Ising model.

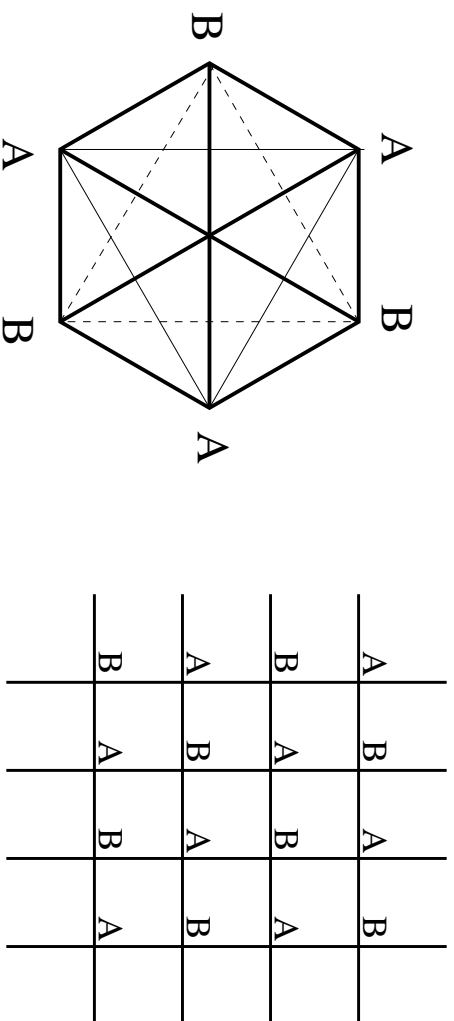
^aJ. Richter, S. Kobe, J. Phys. C: Solid State Phys. **15** (1982) 2193

Definition by Geometry of Interaction I

“Bipartiteness”

Definition IV: **A non-bipartite system is called frustrated.**
 (see e.g. Johannes Richtes, Magdeburg)

Bipartite: If the system can be decomposed into subsystems A and B such that the coupling constants fulfill $J(x_A, y_B) \leq g^2$, $J(x_A, y_A) \geq g^2$, $J(x_B, y_B) \geq g^2$, the system is called bipartite.



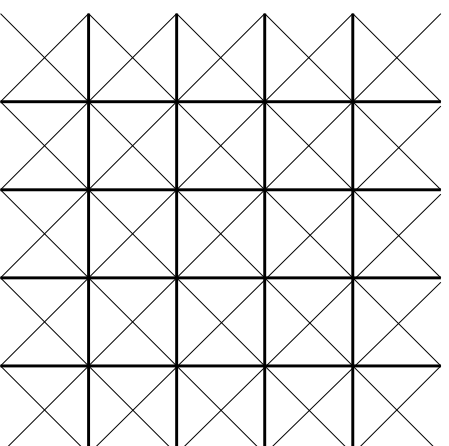
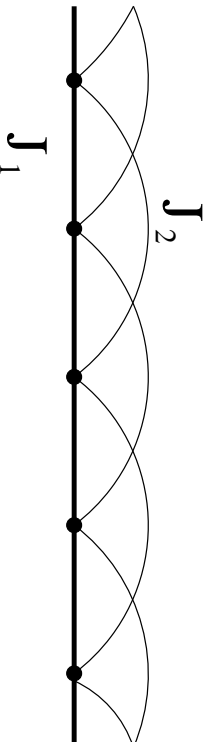
Definition by Geometry of Interaction II

Examples

$$\tilde{H} = -2J_1 \sum_{nn} \vec{s}(x) \cdot \vec{s}(y) - 2J_2 \sum_{mnn} \vec{s}(x) \cdot \vec{s}(y)$$

Consider e.g. rings with an even number of sites or infinite lattices, i.e. $J_1 < 0$ & $J_2 = 0$ should allow a bipartite lattice.

- $J_1 < 0$ & $J_2 > 0$ results in a bipartite lattice, the system is not frustrated.
- $J_1 < 0$ & $J_2 < 0$ results in a non-bipartite lattice, the system is frustrated and the coupling strength J_2 is sometimes itself called frustration.



Summary

Classical Heisenberg and Ising models

- term frustration well defined;
- phenomena like non-trivial ground state degeneracy, weakening of binding energy, spin-spin correlation not -1 etc.;

Quantum Heisenberg models

- term frustration may be defined like in definition IV: non-bipartite \equiv frustrated;
- but resulting general phenomena unclear: ground state energy sometimes degenerate, spin-spin correlation almost always not -1 , weakening of binding compared to “neighbouring” non-frustrated systems - but what is “neighbouring”?